

# Infrared spectra of C<sub>2</sub>H<sub>2</sub> under jet-cooled and *para*-H<sub>2</sub> matrix conditions

Ying-Chi Lee <sup>a</sup>, V. Venkatesan <sup>a</sup>, Yuan-Pern Lee <sup>a,b,\*</sup>, P. Macko <sup>c,1</sup>, K. Didiache <sup>c,2</sup>,  
M. Herman <sup>c,\*</sup>

<sup>a</sup> Department of Applied Chemistry, Institute of Molecular Science, National Chiao Tung University, Hsinchu 30010, Taiwan

<sup>b</sup> Institute of Atomic and Molecular Sciences, Academia Sinica, Taipei 10617, Taiwan

<sup>c</sup> Laboratoire de Chimie quantique et Photophysique, CP160/09, Université libre de Bruxelles, Ave. Roosevelt, 50, B-1050, Brussels, Belgium

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## Abstract

In spectra of jet-cooled C<sub>2</sub>H<sub>2</sub> recorded with an FTIR spectrometer, the  $\nu_5$ ,  $\nu_4 + \nu_5$ ,  $\nu_3$  and  $\nu_2 + \nu_4 + \nu_5$  bands all exhibit an intensity distribution corresponding to  $\sim 6$  K for rotation, with no evidence of nuclear spin conversion. Spectra of C<sub>2</sub>H<sub>2</sub> isolated in solid *p*-H<sub>2</sub> show no evidence of rotation of C<sub>2</sub>H<sub>2</sub>. The strong interaction between  $\nu_3$  and  $\nu_2 + \nu_4 + \nu_5$  in the gas phase is diminished in solid *p*-H<sub>2</sub>. Lines associated with dimer, trimer and tetramer of C<sub>2</sub>H<sub>2</sub> are identified. Spectral features characteristic of solid state acetylene are observed under jet-cooled conditions.

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## 1. Introduction

Acetylene (C<sub>2</sub>H<sub>2</sub>) is an important species in planetary atmospheres and interstellar space. Its gas-phase spectra have been extensively investigated [1,2]. Clusters and crystalline C<sub>2</sub>H<sub>2</sub> are expected also to be present in a planetary atmosphere [3]; their spectra are reasonably well characterized [4].

Infrared absorption spectra of C<sub>2</sub>H<sub>2</sub> dispersed in various matrices have been reported [5–7]. Although several small non-diatomic molecules such as CH<sub>4</sub>, NH<sub>3</sub>, H<sub>2</sub>O are known to rotate in matrices [8,9], no evidence of rotation of C<sub>2</sub>H<sub>2</sub> in low-temperature matrices has been reported.

Because of the ‘softness’ associated with the properties of *p*-H<sub>2</sub> as a quantum solid, guest molecules are expected to rotate more readily in solid *p*-H<sub>2</sub> than in other matrices. The rotational parameters of species isolated in *p*-H<sub>2</sub> are typically  $\sim 90\%$  of those for the gas phase [10,11]. Internal rotation of CH<sub>3</sub>OH is reported to occur in solid *p*-H<sub>2</sub>, but not in solid Ne or Ar [12]. Because of the special properties associated with *p*-H<sub>2</sub>, it would be interesting to discover whether C<sub>2</sub>H<sub>2</sub> can rotate in this quantum solid. A direct comparison of spectra of jet-cooled C<sub>2</sub>H<sub>2</sub> and matrix-isolated C<sub>2</sub>H<sub>2</sub> not only helps in the assignments but also reveals spectral differences between C<sub>2</sub>H<sub>2</sub> in the gas phase and in matrices.

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## 2. Experiments

Spectra of jet-cooled C<sub>2</sub>H<sub>2</sub> were measured under high resolution in Belgium. The jet-expansion system named FANTASIO is described elsewhere [13] and only the major features are presented here. The circular nozzle has a diameter 500  $\mu$ m. Spectra are recorded with a Fourier-transform infrared spectrometer (FTIR) (Bruker IFS120HR) at

\* Corresponding authors. Address: Department of Applied Chemistry, Institute of Molecular Science, National Chiao Tung University, Hsinchu 30010, Taiwan (Y.-P. Lee). Fax: +886 3 5713491.

E-mail addresses: [yplee@mail.nctu.edu.tw](mailto:yplee@mail.nctu.edu.tw) (Y.-P. Lee), [mherman@ulb.ac.be](mailto:mherman@ulb.ac.be) (M. Herman).

<sup>1</sup> FNRS and ARC postdoctoral researcher; Permanent address: Faculty of Mathematics, Physics and Informatics, Comenius University, Mlynská dolina, 84248 Bratislava, Slovakia.

<sup>2</sup> FRIA Researcher.

0.0043  $\text{cm}^{-1}$  resolution (defined as  $0.9/\delta$ , in which  $\delta$  is the maximum optical path difference). Multipass optics around the expansion leads to an increase of signal to noise ratio  $\sim 5$  times over single pass. The  $v_5$  and  $v_4 + v_5$  bands were recorded with a HgCdTe detector, and flow conditions were  $0.08 \text{ L min}^{-1}$  for  $\text{C}_2\text{H}_2$  and  $4.2 \text{ L min}^{-1}$  for Ar, with injection and residual pressures of 700 and  $1.6 \times 10^{-2}$  torr, respectively. For the  $v_3$  region, an InSb detector was used; flow conditions were  $0.05 \text{ L min}^{-1}$  for  $\text{C}_2\text{H}_2$ , and  $2.8 \text{ L min}^{-1}$  for Ar, with injection and stagnation pressures of 470 and  $9 \times 10^{-3}$  torr, respectively. The contribution of the residual gas in the cell, having higher temperature, could be eliminated from the spectrum by a subtraction procedure based on measurements of molecular density from a quadrupole mass filter with a retractable probe coupled to the chamber.

The matrix experiments were performed in Taiwan [12]. A nickel-plated copper plate, maintained at 3.6 K with a closed-cycle refrigerator, served both as a cold substrate for the matrix sample and as a mirror to reflect the incident infrared (IR) beam to the detector. IR absorption spectra were recorded with a FTIR (Bomem, DA8) with a HgCdTe detector. A gaseous mixture of  $\text{C}_2\text{H}_2/p\text{-H}_2$  (1/5000 to 1/22000) was deposited over a period of 1–3 h. Typically, 200 scans at a resolution of  $0.05 \text{ cm}^{-1}$  were recorded. In some experiments the matrix sample was maintained  $\sim 4.5 \text{ K}$  for 0.5–1.0 h for annealing, but all IR measurements were performed at 3.6 K.

$\text{C}_2\text{H}_2$  (99.6%) was degassed at 77 K before use.  $\text{H}_2$  (99.9999%) was used after passage through a trap at 77 K before its conversion to  $p\text{-H}_2$ . The efficiency of conversion is controlled by the temperature of the catalyst; the concentration of *o*- $\text{H}_2$  is  $\sim 100$  ppm at a conversion temperature of 15 K.

### 3. Results and discussion

#### 3.1. Spectra of jet-cooled $\text{C}_2\text{H}_2$

The region  $700\text{--}1400 \text{ cm}^{-1}$  contains the  $v_5$  ( $729.163 \text{ cm}^{-1}$ ) and  $v_4 + v_5$  ( $1328.081 \text{ cm}^{-1}$ ) bands. The  $v_5$  band, corresponding to excitation of the *cis*-bending motion, is of  $\Pi_u \leftarrow \Sigma_g^+$  type, thus with P, R and prominent Q branches (Fig. 1a). The weaker  $v_4 + v_5$  band, of  $\Sigma_u^+ \leftarrow \Sigma_g^+$  type with P and R branches (Fig. 2a), corresponds to simultaneous excitation of *cis*- and *trans*-bending motions [2].

The region  $3270\text{--}3310 \text{ cm}^{-1}$  contains the  $v_3$  and  $v_2 + v_4 + v_5$  ( $3281.899$  and  $3294.839 \text{ cm}^{-1}$ ) bands, both are of  $\Sigma_u^+ \leftarrow \Sigma_g^+$  type with P and R branches (Fig. 3a) and are already reported under jet-cooled conditions using FTIR [14]. A strong anharmonic resonance connecting the two upper vibrational states is responsible for the observation of  $v_2 + v_4 + v_5$  that hence acquires intensity similar to that of the  $v_3$  [15].

The rotational temperature was determined to be 6.1 and 7.0 K, respectively, from relative intensities of rovibrational lines in  $v_3$  and  $v_4 + v_5$ . The latter value therefore also

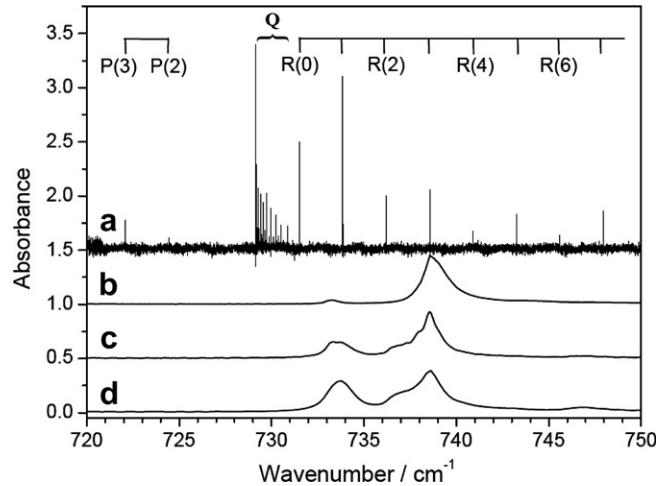


Fig. 1. IR absorption spectra of  $\text{C}_2\text{H}_2$  in the region  $720\text{--}750 \text{ cm}^{-1}$ . (a) Under jet-cooled conditions (see Section 2), with rotational assignments indicated; (b) annealed  $\text{C}_2\text{H}_2/p\text{-H}_2$  (1/22000) matrix; (c)  $\text{C}_2\text{H}_2/p\text{-H}_2$  (1/10000) matrix; (d)  $\text{C}_2\text{H}_2/p\text{-H}_2$  (1/5000) matrix.

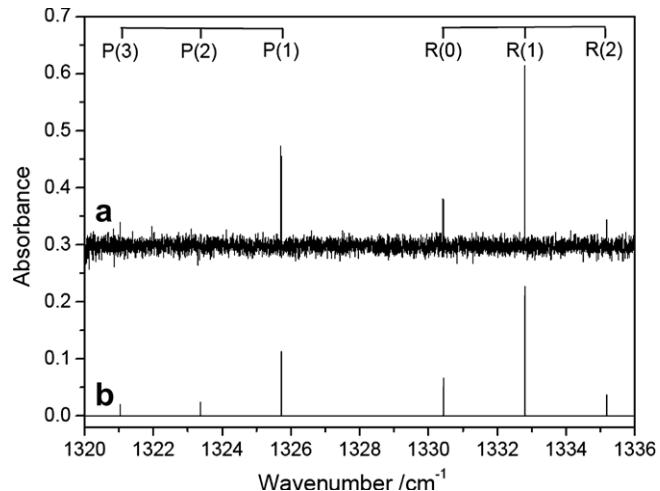


Fig. 2. IR absorption spectra of  $\text{C}_2\text{H}_2$  under jet-cooled conditions in the region  $1320\text{--}1340 \text{ cm}^{-1}$ . (a) Observed, with the P(2) line within the experimental noise limit and P(3) barely identifiable; (b) simulated for  $T_{\text{rot}} = 7 \text{ K}$  with no nuclear-spin conversion; rotational assignments are indicated.

characterizes  $v_5$ , which was recorded simultaneously with  $v_4 + v_5$  but suffers from saturation. The intensity conforms to nuclear-spin statistics 1:3 in favor of odd J-lines, as confirmed by spectral simulations. An example is presented in Fig. 2 for  $v_4 + v_5$ . It thus indicates the absence of nuclear-spin conversion during cooling processes to reach 6 K; these results are similar to those reported previously for  $\text{C}_2\text{H}_2$  at 31 K in the literature [14].

#### 3.2. Spectra of $\text{C}_2\text{H}_2$ in solid $p\text{-H}_2$

IR spectra of samples of  $\text{C}_2\text{H}_2/p\text{-H}_2$  (1/5000 to 1/22000) at 3.6 K exhibit intense features near  $738$  and  $3279 \text{ cm}^{-1}$ , as shown in Traces B–D of Figs. 1 and 3, respectively, for var-

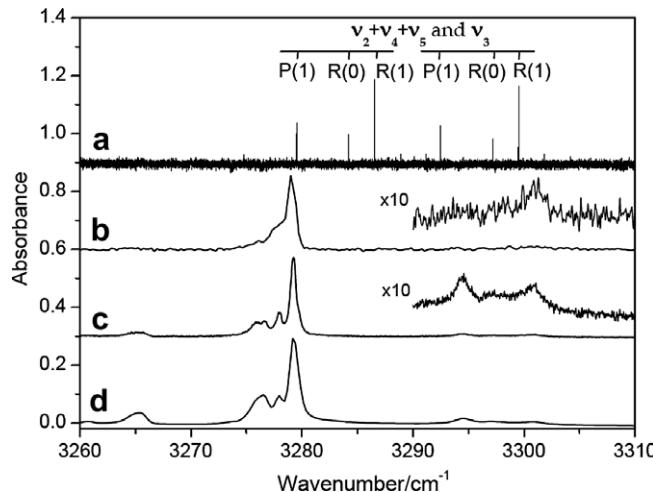


Fig. 3. IR absorption spectra of  $\text{C}_2\text{H}_2$  in the region 3250–3310  $\text{cm}^{-1}$ . (a) Under jet-cooled conditions (see Section 2), with rotational assignments indicated for the two observed bands ( $v_3$  and  $v_2 + v_4 + v_5$ ); (b) annealed  $\text{C}_2\text{H}_2/p\text{-H}_2$  (1/22000) matrix; (c)  $\text{C}_2\text{H}_2/p\text{-H}_2$  (1/10000) matrix; (d)  $\text{C}_2\text{H}_2/p\text{-H}_2$  (1/5000) matrix.

ious molar ratios. At such a low temperature, if rotation were feasible and nuclear spin conserved, one would expect to observe absorption lines originating only from ground-state levels  $J'' = 0$  (*para*) and 1 (*ortho*) of  $\text{C}_2\text{H}_2$  in solid *p*- $\text{H}_2$ . If rotation of  $\text{C}_2\text{H}_2$  does not occur in the *p*- $\text{H}_2$  matrix, only a single line corresponding to the purely vibrational transition is predicted to appear.

### 3.2.1. The $v_5$ band

An experiment with an annealed sample of  $\text{C}_2\text{H}_2$  in highly pure *p*- $\text{H}_2$  (*o*- $\text{H}_2$  less than 0.23% and a slight  $\text{H}_2\text{O}$  impurity) yielded a spectrum showing side bands much diminished relative to the main band at 738.5  $\text{cm}^{-1}$  (Fig. 1b). The observed spectral pattern is inconsistent with that expected for  $\text{Q}(1)$ ,  $\text{R}(0)$  and  $\text{R}(1)$ , with no nuclear-spin conversion (similar to Fig. 1a). These observed features of  $\text{C}_2\text{H}_2$  in *p*- $\text{H}_2$  are therefore unlikely to be attributable to rotational lines of  $v_5$ . The feature at 738.5  $\text{cm}^{-1}$  is thus assigned to a purely vibrational – thus rotationless – transition of  $v_5$ .

It should be pointed out that at higher concentrations two lines at 733.7 and 738.5  $\text{cm}^{-1}$  were observed (Fig. 1c and d); their separation of 4.8  $\text{cm}^{-1}$  is near that, 4.7  $\text{cm}^{-1}$ , between  $\text{Q}(1)$  and  $\text{R}(1)$  lines in the jet-cooled spectrum (Fig. 1a) and might lead to some confusion. Similarly, two lines were observed when substantial *o*- $\text{H}_2$  impurity was present, as shown in Fig. 4; the intensity of the feature near 733.7  $\text{cm}^{-1}$  increased relative to that near 738.5  $\text{cm}^{-1}$  as the concentration of *o*- $\text{H}_2$  increased from 0.28% (trace C) to 0.51% (trace B) and to 1.21% (trace A), and the maximum also shifted from 733.8  $\text{cm}^{-1}$  to 735.0  $\text{cm}^{-1}$ . The feature near 733.7  $\text{cm}^{-1}$  can therefore be ascribed, at least in part, to  $\text{C}_2\text{H}_2$  with nearby *o*- $\text{H}_2$ . The observation in Fig. 4 that  $v_5$  (*cis*-bending) but not  $v_3$  (C–H stretching) is affected by the presence of *o*- $\text{H}_2$  indicates

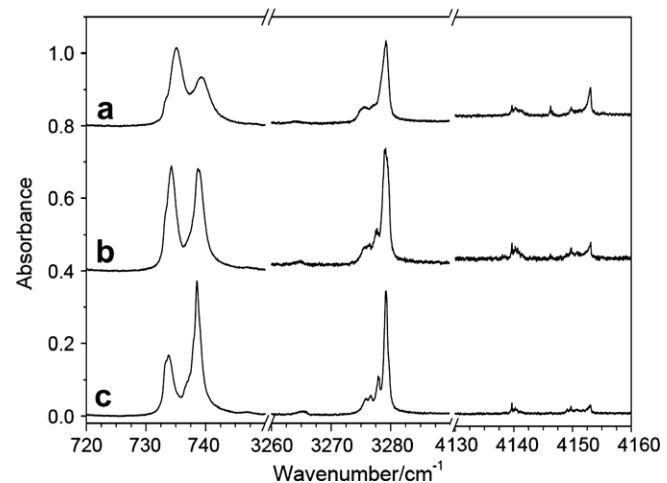


Fig. 4. IR absorption spectra of  $\text{C}_2\text{H}_2$  in the region 720–750, 3260–3290 and 4130–4160  $\text{cm}^{-1}$  for solid  $\text{C}_2\text{H}_2/p\text{-H}_2$  (1/10000); impurity levels of *o*- $\text{H}_2$  are (a) 1.21%, (b) 0.51% and (c) 0.28%. The intensity of the band at 4153  $\text{cm}^{-1}$ , corresponding to the  $\text{Q}_1(0)$  line of  $\text{H}_2$ , indicates the relative concentration of *o*- $\text{H}_2$  in various samples.

that *o*- $\text{H}_2$  might be attached at the side of  $\text{C}_2\text{H}_2$ . A doublet structure was also observed in  $v_5$  for T-shaped complexes of  $\text{C}_2\text{H}_2$  with  $\text{HX}$  ( $\text{X} = \text{F}, \text{Cl}, \text{Br}$ ) in solid  $\text{Ar}$  [16], and for  $\text{C}_2\text{H}_2$  in solid  $\text{N}_2$  [7].

### 3.2.2. The $v_3$ band

The matrix spectrum in highly pure *p*- $\text{H}_2$  shows an intense line at 3279.2  $\text{cm}^{-1}$  and some unresolved weak features in the  $v_3$  region (Fig. 3b); the spectral pattern also indicates that there is no rotation for  $\text{C}_2\text{H}_2$ . At higher concentrations, three subsidiary features at 3278.0, 3276.6 and 3275.9  $\text{cm}^{-1}$  increase in intensity (Fig. 3c and d). These features are unlikely to be associated with rotational structure of  $v_3$ . Furthermore, they are unlikely to arise from any other vibrational transition of the monomer. An exhaustive survey of possible vibrational energy states in  $^{12}\text{C}_2\text{H}_2$  indicates that no accessible state other than  $v_2 + v_4 + v_5$  exists within 100  $\text{cm}^{-1}$  of  $v_3$  [1]. In the gas phase, the two bands are separated by  $\sim 13 \text{ cm}^{-1}$  and have nearly identical intensities (Fig. 3a) [15]. In the matrix, the observed weak features are separated from  $v_3$  by less than 4  $\text{cm}^{-1}$ . Assigning one of them to  $v_2 + v_4 + v_5$  would imply a significantly reduced intramolecular anharmonic coupling between  $v_2 + v_4 + v_5$  and  $v_3$  in the *p*- $\text{H}_2$  matrix.

The weak feature at 3300.9  $\text{cm}^{-1}$  (Fig. 3b) might correspond to the  $v_2 + v_4 + v_5$  assignment. A 2 by 2 interaction matrix model reproducing the anharmonic coupling between  $v_3$  and  $v_2 + v_4 + v_5$  using the parameters in Ref. [15] shows that, with a splitting  $\sim 20 \text{ cm}^{-1}$ , the intensity ratio for bands  $v_2 + v_4 + v_5$  to  $v_3$  is expected to decrease from  $\sim 1$  in the gas phase to  $\sim 0.056$  if the interaction matrix element is assumed to be unchanged from the gas to the matrix. An observed intensity ratio of  $0.04 \pm 0.01$  for these two features is consistent with this prediction. The intensity of the  $v_2 + v_4 + v_5$  band relative to the  $v_3$

band was observed to be reduced to  $\sim 0.5$  in an Ar matrix with energy splitting  $\sim 14 \text{ cm}^{-1}$  [5], and to  $< 0.03$  in an  $\text{N}_2$  matrix with energy splitting  $\sim 28.4 \text{ cm}^{-1}$  [7].

### 3.2.3. The $v_4 + v_5$ band

Two broad features at  $1331.6$  and  $1340.1 \text{ cm}^{-1}$  were observed in the  $v_4 + v_5$  region. Similar to that observed for  $v_5$ , the intensity of the former increases relative to that of the latter as the concentration of *o*- $\text{H}_2$  increases from 0.28% to 1.21%. Hence, we assign the feature at  $1340.1 \text{ cm}^{-1}$  to  $\text{C}_2\text{H}_2$  and the feature at  $1331.6 \text{ cm}^{-1}$  to  $\text{C}_2\text{H}_2$  with nearby *o*- $\text{H}_2$ .

**Table 1** compares observed wavenumbers of the various spectral features assigned to monomer absorption in various environments.

### 3.3. Absorption bands of clusters

Three intense lines at  $3272$ ,  $3266$  and  $3261 \text{ cm}^{-1}$  were reported in the literature for clusters of  $\text{C}_2\text{H}_2$  produced under jet-cooled conditions and are assigned to the T-shaped dimer, trimer and tetramer of  $\text{C}_2\text{H}_2$ ; these lines are separated from the monomer band by  $17$ ,  $23$  and  $28 \text{ cm}^{-1}$ , respectively [4]. They were also observed using FANTASIO (Fig. 5, flow conditions  $\text{C}_2\text{H}_2 = 0.35 \text{ L min}^{-1}$  and  $\text{Ar} = 7 \text{ L min}^{-1}$ ) but are not specifically analyzed yet. Upon dilution of  $\text{C}_2\text{H}_2$  in solid *p*- $\text{H}_2$ , we observed that lines at  $3265.8$ ,  $3260.9$  and  $3254.8 \text{ cm}^{-1}$  have reduced intensities relative to that of the main feature at  $3279.2 \text{ cm}^{-1}$  (Fig. 3b–d); the separations of  $13.6$ ,  $18.5$  and  $24.6 \text{ cm}^{-1}$  are slightly smaller but parallel to those observed in the gas phase. We tentatively assign these lines to absorption of dimer (perhaps also the line at  $3275.9 \text{ cm}^{-1}$ ), trimer and tetramer of  $\text{C}_2\text{H}_2$ , respectively. Analogously, we assign the observed line at  $746.5 \text{ cm}^{-1}$  (and perhaps a line overlapping the feature at  $733.7 \text{ cm}^{-1}$ ) to a dimer of  $\text{C}_2\text{H}_2$ . A band due to solid state acetylene near  $3235 \text{ cm}^{-1}$ , similar to the one reported in a pulsed-jet experiment [17] was also observed here in the jet (Fig. 5), but not in the matrix, at molar ratios of acetylene/*p*- $\text{H}_2 \leq 1/400$ . At this concentration, direct deposition of a flowing (not jet-cooled) mixture of acetylene in *p*- $\text{H}_2$  cannot form solid acetylene on the cold support.

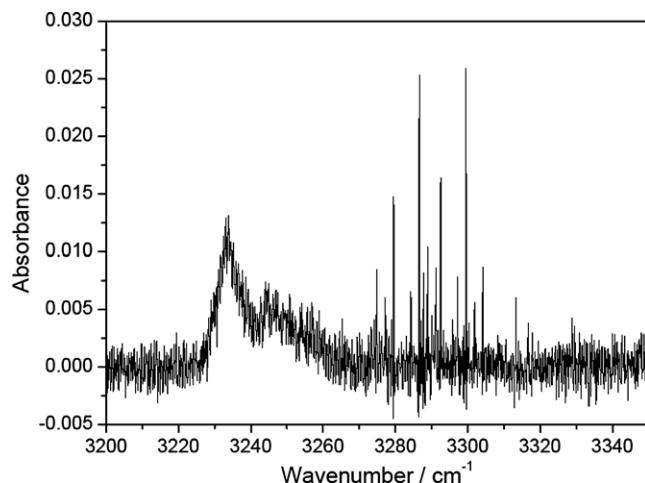


Fig. 5. IR absorption spectra of solid-state  $\text{C}_2\text{H}_2$  under jet-cooled conditions (see Section 2) in the region  $3200$ – $3350 \text{ cm}^{-1}$ .

## 4. Conclusion

Infrared absorption spectra of  $\text{C}_2\text{H}_2$  under jet-cooled conditions and in solid *p*- $\text{H}_2$  were compared. No evidence for rotation of  $\text{C}_2\text{H}_2$  in solid *p*- $\text{H}_2$  has been detected. The interaction between the  $v_3$  and  $v_2 + v_4 + v_5$  states of  $\text{C}_2\text{H}_2$ , known to be strong in the gas phase, is much weaker, if not absent, in solid *p*- $\text{H}_2$ . Additional features associated with bending modes of  $\text{C}_2\text{H}_2$  were ascribed to  $\text{C}_2\text{H}_2$  with nearby *o*- $\text{H}_2$ . Bands due to dimer, trimer and tetramer of  $\text{C}_2\text{H}_2$  were tentatively assigned for  $\text{C}_2\text{H}_2$  in solid *p*- $\text{H}_2$ , whereas evidence for solid acetylene was obtained in the expansion.

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Table 1  
Origin/ $\text{cm}^{-1}$  of  $\text{C}_2\text{H}_2$  bands observed under jet conditions and in matrices

Band	Jet	$p$ - $\text{H}_2$ matrix	Ar matrix [5]	Kr matrix [6]	$\text{N}_2$ matrix [7]
$v_5$	729.163	738.5	736.8	732	742.0, 747.4
$v_4 + v_5$	1328.081	1340.1	1334.5	1325.5	–
$v_3^a$	3281.899	3279.2	3288.9	3280	3282.6
$v_2 + v_4 + v_5^a$	3294.839	3300.9	3302.9	3293	3311.0
	(~1.0)	(~0.04)	(~0.5)	(~1.0)	(<0.03)

<sup>a</sup> The  $v_3$  and  $v_2 + v_4 + v_5$  states are connected by anharmonic resonance. The intensities of the  $v_2 + v_4 + v_5$  band relative to that of the  $v_3$  band are listed in parentheses.

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