Ultrahigh supermode noise suppressing ratio of a semiconductor optical amplifier filtered harmonically mode-locked Erbium-doped fiber laser

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Abstract: The supermode noise suppressing ratio (SMSR) and the phase noise of a harmonically mode-locked Erbium-doped fiber laser (HML-EDFL) with an intra-cavity semiconductor optical amplifier (SOA) and an optical band-pass filter (OBPF) are improved and compared with a state-of-the-art Fabry-Perot laser diode (FPLD) injection-mode-locked EDFL. By driving the intra-cavity SOA based high-pass filter at unitary gain condition, the SMSR of the HML-EDFL is enhanced to 82 dB at the cost of degrading phase noise, increasing jitter, and broadened pulsewidth. The adding of OBPF further improves the SMSR, pulsewidth, phase noise, and jitter of the SOA-filtered HML-EDFL to 90 dB, 42 ps, -112 dBc/Hz, and 0.7 ps, respectively. The ultrahigh SMSR of the SOA-filtered HML-EDFL can compete with that of the FPLD injection-mode-locked EDFL without sacrificing its pulsewidth and jitter performances.

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1. Introduction

Typically, the harmonically mode-locking (HML) and injection-mode-locking (IML) are two intriguing techniques used for generating ultrashort and high repetition-rate optical pulse-train in Erbium-doped fiber laser (EDFL) systems [1, 2] which are mandatory for applications in high-speed optical time division multiplexed transmission system [3] and electro-optical sampling systems [4]. In principle, the HML-EDFL can be implemented by either using an integrated electro-optic/electro-absorption modulator (EOM/EAM) [1, 2] or using a semiconductor optical amplifier (SOA) based gain modulator [5, 6]. Alternatively, Li *et al.* used a feedback injected FPLD as a mode-locker to implement an IML-EDFL [7], which can simply be realized as a mutually injection-locked link between the gain-switched FPLD and the EDFL. Particularly, the FPLD-IML-EDFL link was reported to exhibit extremely low single-sided-band (SSB) phase noises, small timing jitter, and supermode-noise-free performances [8]. In contrast, most of the HML-EDFLs are sensitive to environmental perturbations such as temperature fluctuations and mechanical vibrations due to their

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relatively long cavity length associated with numerous longitudinal modes. Therefore, the HML-EDFL pulses usually suffer from the SSB phase noise, pulse-to-pulse timing jitter and the supermode noise, which seriously degrade the temporal resolution or the bit error rate for the HML-EDFL used in versatile applications [9]. To overcome, a regeneratively HML-DFL was demonstrated to suppress the supermode noise and the SSB phase noise (as well as timing jitter) [10-12]. Later on, Duan *et al.* reported that the supermode noise and intensity noises in pulses could be greatly suppressed by adding an SOA as a high-pass filter [13]. In addition, Schell *et al.* reported the reduction of SSB phase noise and timing jitter in a gain-switched FPLD with self-feedback injection [14].

Although the maximum supermode noise-suppressing ratio (SMSR) of up to 33 dB was reported for the SOA-filtered HML-EDFL [13], its SSB phase noise characteristic has yet never been investigated. In principle, the disadvantage of the SOA-filtered HML-EDFL is that both the SSB phase noise and timing jitter are degraded due to the gain of SOA. In this paper, the SSB phase noise, pulse-to-pulse timing jitter, and supermode noise characteristics of the HML-EDFL with intra-cavity SOA filter are theoretically investigated and experimentally compared with a start-of-the-art FPLD-IML-EDFL. We demonstrate that the SMSR of the SOA-filtered HML-EDFL can be significantly improved without sacrificing its pulsewidth, SSB phase noise, and timing performances, by concurrently adding an optical band pass filter (OBPF) and adjusting the driving current of the SOA. Theoretical simulations are also performed to elucidate that the overall performances of the SOA-and-OBPF-filtered HML-EDFL can be comparable with those of a FPLD-IML-EDFL.

2. Experiment

Figure 1(a) illustrates a typical HML-EDFL with an intra-cavity EOM based mode-locker, which consists of an Erbium-doped fiber amplifier (EDFA), a pair of Faraday optical isolators, a polarization controller (PC), a LiNbO₃ Mach-Zehnder interferometer based EOM, and an optical coupler (OC) with 10 % output coupling ratio.



Fig. 1. The schematic diagrams of the (a) HML-EDFL and (b) SOA-and-OBPF-filtered HML-EDFL. PC: polarization controller; OC: optical coupler; EDFA: Erbium-doped fiber amplifier; MZM: Mach-Zehnder intensity modulator; SOA: semiconductor optical amplifier; OBPF: optical bandpass filter.

The EOM DC-biased at its half-wave voltage ($V_{\pi} \approx 8$ V) is driven by a microwave synthesizer (Rhode & Schwartz, SML01) with the output power and the SSB phase noise of 22 dBm and -130 dBc/Hz (measured at offset frequency of 100 kHz from carrier), respectively, at frequency of 1 GHz. The PC is properly adjusted to optimize the polarization of the circulating pulses for the EOM, and two Faraday optical isolators ensure the unidirectional propagation of light in the EDFL cavity. As a result, the stable HML-EDFL pulse-train can be obtained when the modulating frequency is detuned to coincide with one harmonic of longitudinal modes in the EDFL cavity. The length of the EDFL ring cavity is 32 m, providing a longitudinal mode spacing of 6.25 MHz. On the other hand, the schematic diagram of an SOA-filtered HML-EDFL is shown in Fig. 1(b).

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pigtailed SOA exhibits a small-signal gain of 25 dB and a saturated output power of -1.5 dBm when operating at unitary gain (or nearly transparent) condition. In this case, the bias current, operating temperature, 3-dB linewidth and central wavelength of the SOA are set as 57 mA, 24 °C, 30.4 nm, and 1530 nm, respectively. A PC is used to adjust the polarization of light for the SOA. To suppress the supermode noise and SSB phase noise, an OBPF (JDS Uniphase, TB1500B) with 3dB bandwidth 1.38 nm is inserted into the HML-EDFL.



Fig. 2 The schematic diagram of the EDFL mutually injection-mode-locked with a gainswitched FPLD. Comb: Electrical pulse generator; FPLD: Fabry-Perot laser diode.

On the other hand, the FPLD-IML-EDFL is performed by seeding the EDFL with a gainswitched FPLD, as shown in Fig. 2. It consists of a close-loop EDFA with 10 % output coupling ratio, a comb-generator-driven FPLD, an OBPF, and an optical circulator. The wavelength, threshold current, and longitudinal mode spacing of the free-running FPLD operated at 25 °C are about 1550 nm, 8 mA, and 1.2 nm, respectively. A comb generator driven by 27-dBm microwave signal at 1 GHz is employed to provide an electrical pulse-train for gain-switching the FPLD DC-biased at 3.4 mA. The central wavelength of the OBPF is adjusted to match that of the FPLD, which avoids the feedback injection of the amplified spontaneous emission (ASE) of the EDFL into the FPLD. In particular, the gain-switched FPLD pulse amplified by the EDFL is fedback into the FPLD itself, achieving a mutual IML between the FPLD and EDFL link. Such a configuration effectively suppresses the ASE as well as the SSB phase noise of the EDFL. The feedback injection from EDFL is also used to facilitate single longitudinal mode lasing of the FPLD with improved SMSR. Note that the feedback wavelength of the amplified FPLD pulses must coincide with the central longitudinal mode of the FPLD at 1550 nm in order to obtain the lowest SSB phase noise (timing jitter) and the highest SMSR.

The pulsewidth and the output power are monitored using a digital sampling oscilloscope (Agilent, 86100A+86109B), an optical spectrum analyzer with 0.01-nm resolution (Advantest, Q8384), and an optical power meter (ILX Lightwave, OMM-6810B+6727B), respectively. Assuming that the optical pulse-train exhibits small and stationary phase deviations (*i.e.* small timing jitter) without any secondary phase deviating sidebands, a spectral domain technique is employed to characterize the SSB phase noise and pulse-to-pulse timing jitter characteristics of the EDFL pulses [10, 15-19]. Such a technique neglects the influences of pulse shape and pulsewidth fluctuations, which assumes no correlation between amplitude and phase fluctuations. The SSB phase noise is quantified by measuring the noise spectral power density with an electrical spectrum analyzer (ESA, Agilent HP8565E) at a resolution bandwidth of 1 Hz in connection with a photodetector (New Focus 1014). The timing jitter and the SMSR in a given bandwidth are then calculated and analyzed by using phase-noiseutility software of the RF spectrum. By subtracting the SSB phase noise spectrum at a higher harmonic frequency (for example, the 10^{th} harmonics, n = 10) with that at the fundamental frequency (n = 1) of the EDFL pulses, the rms timing jitter in a bandwidth extending from f_L to f_H is given by [20]:

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$$\sigma(f) = \frac{1}{2\pi f_0} \left\{ 2 \int_{f_L}^{f_H} \left[\left[\left[10^{L_n(f)/10} - 10^{L_1(f)/10} \right] \right] / \left(n^2 - 1\right) \right] df \right\}^{\frac{1}{2}}$$
(1)

where f_L and f_H are integration boundaries, $L_I(f)$ and $L_n(f)$ are phase noise power spectral densities of fundamental and nth harmonics signals, respectively. The n denotes the harmonic number and f_0 is the repetition frequency of the laser pulse.



Fig. 3. The peak power (solid up-triangle) and pulsewidth (solid square) as a function of detuning frequency for (a) the HML-EDFL and (b) the FPLD-IML-EDFL systems.

3. Results and discussion

3.1 The HML-EDFL and the FPLD-IML-EDFL

The optimized HML-EDFL pulse is determined from the evolution of pulse shape, especially the observation on the maximum peak power and the shortest pulsewidth versus detuning frequency, as shown in Fig. 3(a). It is clearly observed that the peak power is attenuated and the pulse is broadened as the modulating frequency detuned. The 3-dB detuning bandwidth (defined as the difference between the optimized mode-locking frequency and the frequency corresponding to the doubling in pulsewidth) of the HML-EDFL is only 7 kHz. The shortest pulsewidth of 36 ps and the maximum output peak power of 12.3 mW of the HML-EDFL are obtained, as shown in Fig. 3(a).



Fig. 4. (a) Supermode noise spectrum (left) and pulse shape (right) of the HML-EDFL (upper part) and the HML-EDFL with intra-cavity SOA and OBPF (lower part). (b) The supermode noise spectra measured at resolution bandwidth of (1) 1 Hz; (2) 10 Hz; and (3) 100 kHz.

The supermode noise results from the beating between longitudinal modes and ASE of the HML-EDFL, which inevitably causes large intensity fluctuations in the HML-EDFL pulsetrain. Typically, the HML-EDFL exhibits a relatively small SMSR of 45 dB (see Fig. 4), owing to the extremely long upper-level lifetime of excited erbium ions in EDFL (~10 ms) and the relatively low mode-beating frequencies of the HML-EDFL with a long cavity length. Note that the SMSR measurement was carefully performed in our experiments, since which is dependent on the background noise of the electrical spectrum analyzer under an extremely low supermode noise condition. For example, the measured SMSR could be distorted when

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the electrical spectrum analyzer is set at larger resolution bandwidth (RBW), as shown in Fig. 4 (b). By reducing the RBW from 100 kHz to 1 Hz, the SMSR is enlarged from 45 dB to 90 dB as the background noise of the spectrum analyzer decreases from -85 dBm to -126 dBm. As a result, the SMSR can only be precisely obtained by setting the smallest RBW of the electrical spectrum analyzer. In conventional HML-EDFL, the SSB phase noises and timing jitter at offset frequency of 100 kHz are -114 dBc/Hz and 0.6 ps, respectively, as shown in Fig. 7. The low-frequency spurs at 20 Hz, 55 Hz, and 60 Hz are mainly due to the fans, ambient vibrations, wall current, and their harmonics, respectively [21].



Fig. 5. (a) The pulse shapes and (b) the supermode noise spectrum (RBW = 1 Hz) of (1) the free-running FPLD, and (2) the FPLD-IML-EDFL.

In a FPLD-IML-EDFL, the gain-switched FPLD pulse experiences a resonant amplification process in the EDFL, and reaches its highest peak power and shortest pulsewidth as the modulating frequency of the FPLD coincides with the harmonic frequency of the longitudinal mode in EDFL. The output power and the pulsewidth of the FPLD-IML-EDFL versus the frequency detuning are shown in Fig. 3(b). The optimized pulse and the peak power of 22 ps and 40.7 mW, respectively, are shown in Fig. 5(a). When the repetition frequency detunes, the pulses circulated in the EDFL arrives the gain medium with a slight temporal deviation, which eventually leads to a broadened pulse shape contributed by the superposition of circulated pulses. A multiple pulse splitting effect happens as the frequency detuned beyond a bandwidth of 17.8 kHz, as shown in Fig. 3(b).



Fig. 6. (a) The SSB phase noise spectra and (b) the timing jitter of (1) the free-running FPLD and (2) the FPLD-IML-EDFL.

The SSB phase noise and the corresponding timing jitter of a free-running FPLD pulse are -114 dBc/Hz at 100 kHz offset frequency and 0.32 ps between a frequency interval of 10 Hz and 100 kHz, respectively, as illustrated in Fig. 6. In comparison, the FPLD-IML-EDFL has lower phase noise of -121 dBc/Hz and timing jitter of 0.25 ps than those of a free-running FPLD due to the mutually IML configuration (see Fig. 6). The spontaneous emission noise of the FPLD can be greatly reduced under such a FPLD-IML-EDFL link, and the energies of adjacent pulses are equalized due to the intensity-dependent gain of the FPLD [8]. Such a

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FPLD based filtering effect is attributed to its relatively fast carrier recovery rate (from 0.5 to 1 ns) and gain saturation effect, providing the enhancement on SMSR from 87 to 91 dB, as shown in Fig. 5(b). In comparison, the SMSR, SSB phase noise, and timing jitter of the FPLD-IML-EDFL are far better than those of the HML-EDFL.



Fig. 7 (a) The SSB phase noise in HML-EDFL (1) without SOA, (2) with SOA and (3) with SOA and OBPF; (b) the timing jitter in HML-EDFL (1) without SOA, (2) with SOA and (3) with SOA and OBPF.

3.2 The HML-EDFL with SOA

Due to a relatively fast carrier recovery rate and gain saturation effect of the SOA [22, 23], some SOA based high-pass filtering schemes proposed for suppressing low-frequency noise in EDFL have previously been demonstrated by Xu *et al.* and Yao *et al.* [22, 24]. Theoretically, the supermode noise beating effect can be quantitatively discussed by defining the gain (R) of output power fluctuation to that of input power fluctuation (defined as $\Delta P/P$ within a frequency interval Δv) in a laser. This can be derived from the rate equation of the small-signal gain under a perturbation on average power [24, 25]. With a simplified gain fluctuation of $g = g_0 + \Delta g$, Peng *et al.* [25] described the equation as

$$\frac{d\Delta g}{dt} + \Delta g \left(\frac{1}{\tau_c} + \frac{p_{in} \cdot e^{g_0}}{E_{sat}} \right) = -\frac{\Delta P_{in}}{E_{sat}} \left(e^{g_0} - 1 \right), \qquad (2)$$

where P_{in} denotes the average power of the stable mode-locked pulses, ΔP_{in} denotes the power fluctuation due to the beating between the supermodes, E_{sat} is the saturated energy of gain medium, and τ_c is the carrier lifetime. Yao *et al.* [24] have concluded the analytical expression of R by solving Eq. (2)

$$R = \frac{\Delta P_{out}/P_{out}}{\Delta P_{in}/P_{in}} = \left[1 - (1 - e^{-g_0}) \times \frac{e^{g_0} P_{avg}}{E_{sat}} \times \frac{\frac{1}{\tau_c}}{\left[\frac{1}{\tau_c}\right]^2 + \Delta \omega^2} \right]^2 + \left[(1 - e^{-g_0}) \times \frac{e^{g_0} P_{avg}}{E_{sat}} \times \frac{\Delta \omega}{\left[\frac{1}{\tau_c}\right]^2 + \Delta \omega^2} \right]^2, \quad (3)$$

where $\tau_c = [\tau_c^{-1} + e^g P_{avg}/E_{sal}]^{-1}$ denotes the effective carrier lifetime, and P_{avg} is the average power of the laser. The first term in the square root describes the power fluctuation of the EDFA or SOA, while the second term denotes the supermode beating fluctuation by longitudinal modes. Typically, the small mode-spacing of conventional EDFL ($\Delta v < 10$ MHz) causes a numerous cavity modes within the mode-locking spectrum (with 3-dB bandwidth >100 GHz), which beats each other in the EDFL cavity to produce supermodes. The *R* of EDFA approaches 1 since $\tau_c \ge 10$ ms and $\Delta v << (2\pi\tau_c)^{-1}$, which indicates a large power fluctuation and a strong supermode beating are usually involved in the output HML-EDFL pulses. Such a supermode_noise_induced power fluctuation can be solved with an intra-cavity SOA in deeply saturating condition ($P_{in}e^g_0\tau_c/E_{sat} >> 1$). The ultrashort τ_c of SOA (≤ 1 ns) makes the *R* as small as e^{rg} . Therefore, the SMSR of the HML-EDFL can be greatly enhanced

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Fig. 8. The SMSR and the SSB phase noise of (a) the HML-EDFL with SOA and (b) the HML-EDFL with SOA and OBPF at different SOA currents.

By taking the advantage of the SOA based high-pass filtering effect, the SMSR performance of a HML-EDFL with an intra-cavity SOA is characterized. The supermode noise located in low-frequency region can thus be suppressed by adding an SOA [23], while the SMSR is significantly improved from 45 dB to 62.4 dB even the SOA is operated at below threshold (absorption) condition. However, the SSB phase noise of the EDFL are degraded from -114 dBc/Hz to -96 dBc/Hz. As the SOA is changed from absorption to unitary gain regime, the SMSR of the HML-EDFL can further be improved to 82 dB or larger. The minimum SSB phase noise and the minimum timing jitter of -104.2 dBc/Hz and 1.4 ps, respectively, can be obtained by driving the SOA at nearly transparent current (~66 mA), as shown in Fig. 8(a). The optimized driving current of the SOA based high-pass filter at its transparent condition is clearly interpreted, while the SMSR can be greatly improved by 40 dB as compared to that of the same HML-EDFL without SOA. Nonetheless, the nearly transparent SOA still causes the EDFL pulsewidth broadening from 36 ps (without the SOA) to 61 ps (with the SOA) and leads to degradation on the SSB phase noise by 10 dB (as compared to the HML-EDFL without SOA). It is evident that the SMSR almost saturates (due to the output power saturation) but the SSB phase noise increases rapidly at higher gain condition of the SOA, which results in the distorted and fluctuated HML-EDFL pulses. In principle, the total noise power of an SOA is contributed by signal-spontaneous beating noise, spontaneous-spontaneous beating noise, amplified signal shot noise, and spontaneous emission shot noise [26-28]. The noise in the SOA is dominated by the spontaneousspontaneous beating noise when the input signal level is extremely small. When the SOA is operated at larger gain region with higher input signal levels, the signal-spontaneous beating noise and the spontaneous-spontaneous beating noise become more pronounced than other noise terms.

3.3 The HML-EDFL with SOA and OBPF

The SSB phase noise of a HML-EDFL with an intra-cavity SOA filter is strongly correlated with the stimulated or spontaneous emission induced carrier density fluctuation [26, 29, 30]. The spontaneous-emission-induced SSB phase noise is written as $S_{\phi l}(f) = hV(G-1)n_{sp}/GP_{in}$, where *h* is Planck's constant, *v* is optical frequency, *G* is amplifier gain, n_{sp} is the spontaneous emission factor, and P_{in} is input optical power. The stimulated_emission_induced indirect SSB phase noise is described as $S_{\phi 2}(f) = [(2\pi K\Gamma/\lambda A)^2 4G(G-1)n_{sp}\tau_c^2 P_{in}]/hV[(2\pi f\tau_c)^2+1]$ with $K = \Delta n_r(z, t)/\Delta n_e(z, t)$, where Δn_r is real refractive index change, Δn_e is fluctuation of carrier density, Γ is optical confinement factor, λ is wavelength, *A* is amplifier cross section, and τ_e is carrier lifetime. The theoretical simulation of $S_{total}(f)=S_{\phi l}(f)+S_{\phi 2}(f)$ interprets that the minimum phase-noise operation of an SOA can be achieved by setting its driving condition in the unitary gain (or transparent) regime, as depicted in Fig. 9.

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Fig. 9. The simulated (solid line) and measured SSB phase noises as a function of the SOA gain for the SOA-filtered EDFL without (solid square) and with OBPF (up triangle).

The spontaneous emission dominates the phase noise of the EDFL-SOA link as the SOA operates at under transparent gain condition, whereas the SSB phase noise is mainly contributed by the stimulated emission of the SOA operated at over transparent region. Note that the signal-spontaneous beating noise is inevitable, however, the spontaneous-spontaneous beating noise, which arises from the beating of the ASE components themselves over a wide gain spectrum, can be reduced by loading a narrow-band OBPF matched to the signal frequency. Without an OBPF or an etalon filter, the reduction of signal-spontaneous beating noise is inevitable for SOA operated at gain regime. Subsequently, the pulsewidth broadening and SSB phase-noise degradation problems are concurrently solved by adding an intra-cavity OBPF and operating the SOA at unitary gain regime. These improve the SSB phase noise performance from -104.2 to -112 dBc/Hz and enhance the SMSR from 86 to 90 dB by suppressing the ASE in the HML-EDFL, as shown in Fig. 8(b). In addition, such a configuration also leads to a shortened pulsewidth and a reduced timing jitter of 42 ps and 0.7 ps, respectively, as shown in Figs. 4 and 7(b). Time-domain analyses shown in Fig. 4 have revealed that the insertion of the SOA and the OBPF greatly suppresses the intensity fluctuations. The SMSR of 90 dB in HML-EDFL with SOA and OBPF is far better than that of 45 dB in a typical HML-EDFL. The parametric comparisons on different EDFL systems are summarized in Table 1.

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Configurations	HML-EDFL	HML-EDFL with SOA	HML-EDFL with SOA and OBPF	FPLD-IML-EDFL
Pulsewidth (ps)	36	61	42	22
Peak Power (mW)	12.3	14	11.1	40.7
SSB Phase Noise (dBc/Hz)	-114	-104.2	-112	-121.2
Timing Jitter (ps)	0.6	1.4	0.7	0.25
SMSR (dB)	45	86	90	91

Table 1. Comparisons on performances of versatile EDFL systems

In comparison, the SMSR of the HML-EDFL with intra-cavity SOA and OBPF is better than those systems using only OBPF [31-33] by 10 dB and is equivalent to that of a FPLD-IML-EDFL system. The minimum SSB phase noise of the SOA-and-OBPF-filtered HML-EDFL found at the transparent current condition of the SOA is comparable with that of conventional HML-EDFL. By adding both SOA and OBPF, the reduction in supermode noise and spontaneous-beating-induced SSB phase noise in a HML-EDFL has been confirmed.

4. Conclusions

We have greatly improved the supermode noise-suppressing ratio (SMSR) of a harmonically

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mode-locked Erbium-doped fiber laser (HML-EDFL) by adding intra-cavity semiconductor optical amplifier (SOA) and optical band-pass filter (OBPF). A typical HML-EDFL system exhibits the SSB phase noise, corresponding timing jitter, and supermode noise suppressing ratio (SMSR) of -114 dBc/Hz (at 100 kHz offset frequency), 0.6 ps (integrating from 10 Hz to 100 kHz), and 45 dB, respectively. In comparison, a gain-switched Fabry-Perot laser diode (FPLD) mutually injection-mode-locked EDFL (FPLD-IML-EDFL) exhibits the start-of-theart performances, such as the SMSR, the SSB phase noise and timing jitter of 91dB, -121.2 dBc/Hz, and 0.25 ps, respectively. The extremely high SMSR is due to the relatively fast carrier recovery rate and the gain saturation effect of the mutual injection-locked FPLD, which acts as a supermode noise suppressor. With an intra-cavity SOA based high-pass filter driven at unitary gain condition, we primarily demonstrate that the SMSR of the HML-EDFL increased to 86 dB becomes comparable with that of the FPLD-IML-EDFL, however, at the cost of degrading SSB phase noise (-104.2 dBc/Hz) and increasing timing jitter (1.4 ps). Theoretical and experimental results conclude that the optimized driving current of the SOA based high-pass filter for the EDFL is at transparent current, while the SSB phase noise is suppressed without sacrificing the SMSR of the EDFL. The SMSR saturates at higher driving current of SOA, whereas the pulsewidth, SSB phase noise and timing jitter of the HML-EDFL seriously degrade. By adding an OBPF into the cavity of the HML-EDFL with SOA, the SMSR and SSB phase noise can further be improved to 90 dB and -112 dBc/Hz, respectively. The pulsewidth and the timing jitter of the HML-EDFL with both SOA and OBPF are reduced to 42 ps and 0.7 ps, respectively. The OBPF maintains the SMSR performance of the HML-EDFL with SOA without sacrificing its pulsewidth, SSB phase noise and timing jitter performances. Such a high SMSR response has already been comparable with that of a stateof-the-art FPLD-IML-EDFL system.

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