

滿足 $\text{Alg } T = \{T\}'$ 之線性變換
 On Contractions Satisfying $\text{Alg } T = \{T\}'$

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ABSTRACT — For a bounded linear operator T on a Hilbert space let $\{T\}'$, $\{T\}''$ and $\text{Alg } T$ denote the commutant, the double commutant and the weakly closed algebra generated by T and 1 , respectively. Assume that T is a completely non-unitary contraction with a scalar-valued characteristic function $\psi(\lambda)$. In this note we show that the condition that $|\psi(e^{it})|=1$ on a set of positive Lebesgue measure implies that $\text{Alg } T = \{T\}'$. Moreover, if $\psi(\lambda)$ is assumed to be outer, then these two conditions are equivalent. Our main result generalizes the well-known fact that the compressions of the shift satisfy $\text{Alg } T = \{T\}'$.

For an arbitrary operator T , it is easily seen that $\text{Alg } T \subseteq \{T\}'' \subseteq \{T\}'$ holds. Now we assume that T is the compression of the shift on the space $H^2 \ominus \psi H^2$, that is, T is defined by

$$(Tf)(\lambda) = P[\lambda f(\lambda)] \text{ for } |\lambda| < 1 \text{ and } f \in H^2 \ominus \psi H^2,$$

where ψ is a scalar-valued inner function and p denotes the (orthogonal) projection onto the space $H^2 \ominus \psi H^2$. Then it was shown by Sarason [2] that $\text{Alg } T = \{T\}'$. (In fact, he showed more than this. He proved that every operator in $\{T\}'$ is of the form $u(T)$, for some $u \in H^\infty$.) Note that the compression of the shift is a completely non-unitary (c.n.u.) contraction whose characteristic function ψ is scalar-valued and satisfies $|\psi(e^{it})|=1$ a.e. In this note a sufficient condition that a c.n.u. contraction having a scalar-valued characteristic function satisfy $\text{Alg } T = \{T\}'$ is given. Indeed we show that for such contractions the condition that $|\psi(e^{it})|=1$ on a set of positive Lebesgue measure implies that $\text{Alg } T = \{T\}'$ holds. Hence Sarason's result follows as a special case. We also show that if ψ is an outer function then the condition is also necessary and both are equivalent to the condition that every invariant subspace for T is hyperinvariant.

Note that for an arbitrary operator whether the latter condition is equivalent to $\text{Alg } T = \{T\}'$ is still unknown.

We should also point out that Sz-Nagy and Foias [5] proved the following result:

A c.n.u. contraction T with the scalar-valued characteristic function ψ satisfies $\{T\}'' = \{T\}'$ if and only if $\psi(\lambda) \neq 0$. Hence in a sense our results strengthen theirs.

In the proofs of our theorems we will extensively use the functional model for c.n.u. contractions. The readers are referred to reference [4] for the basic definitions and terminology. Throughout this note results from [4] will be used without specific mentioning.

Our main result is the following:

THEOREM 1. Let T be a c.n.u. contraction having a scalar-valued characteristic function ψ . Assume $|\psi(e^{it})| = 1$ on a set of positive Lebesgue measure. Then T satisfies $\text{Alg } T = \{T\}'$.

Consider the functional model of T , that is, consider T acting on the space $H = [H^2 \oplus \Delta L^2] \ominus \{\psi W \oplus \Delta W : W \in H^2\}$ by

$$T(f \oplus g) = P(e^{it}f \oplus e^{it}g) \quad \text{for } f \oplus g \in H,$$

where $\Delta = (1 - |\psi|^2)^{\frac{1}{2}}$ and P denotes the (orthogonal) projection onto H . Let S be an operator in $\{T\}'$. Then S is of the form

$$S = P \begin{pmatrix} A & 0 \\ B & C \end{pmatrix},$$

where $A \in H^\infty$ and $B, C \in L^\infty$ satisfy $B\psi + C\Delta = \Delta A$ a.e. (cf. [5]). For an arbitrary operator T on H , let $\text{Lat } T$ denote the lattice of subspaces invariant under T , and $T^{(n)}$ denote the operator $\underbrace{T \oplus \dots \oplus T}_n$ acting on $\underbrace{H \oplus \dots \oplus H}_n$. To show that $S \in \text{Alg } T$, it suffices to show that $\text{Lat } T^{(n)} \subseteq \text{Lat } S^{(n)}$ for any $n \geq 1$ (cf. [1], Theorem 7.1). Note that the characteristic function of $T^{(n)}$ is given by the n by n matrix-valued function

$$\psi = \begin{pmatrix} \psi & & 0 \\ & \ddots & \\ 0 & & \psi \end{pmatrix}.$$

Let $K \in \text{Lat } T^{(n)}$ and $\phi = \phi_2 \phi_1$ be the corresponding regular factorization. We first show the following

LEMMA 2. ϕ_1 and ϕ_2 are n by n matrix-valued functions.

PROOF: Assume that ϕ_1 and ϕ_2 are, respectively, m by n and n by m matrix-valued functions. Consider

$$\delta(e^{it}) = \dim \overline{\Delta(e^{it}) \mathbb{C}^n}$$

$$\delta_1(e^{it}) = \dim \overline{\Delta_1(e^{it}) \mathbb{C}^n}$$

and

$$\delta_2(e^{it}) = \dim \overline{\Delta_2(e^{it}) \mathbb{C}^m},$$

where \mathbb{C} denotes the complex plane, $\Delta(e^{it}) = (1 - \overline{\phi}(e^{it})^* \phi(e^{it}))^{\frac{1}{2}}$ and $\Delta_j(e^{it}) = (1 - \overline{\phi_j}(e^{it})^* \phi_j(e^{it}))^{\frac{1}{2}}$, $j=1, 2$. That $\phi = \phi_2 \phi_1$ is a regular factorization is equivalent to the condition that

$$\delta(e^{it}) = \delta_1(e^{it}) + \delta_2(e^{it}) \quad \text{a.e.} \quad (1)$$

Since $|\psi(e^{it})|=1$ on a set of positive Lebesgue measure, say, α , it follows that $\Delta(e^{it})=0$ on α . Hence $\delta(e^{it})=0$ on α . If $m > n$ then $\phi_2(e^{it})$ cannot be isometric from \mathbb{C}^m to \mathbb{C}^n . Thus $\delta_2(e^{it}) > 0$ a.e., which contradicts (1). Hence we must have $m \leq n$. Now we want to show that actually $m=n$. Assume the contrary, that is $m < n$. Let $\phi_1 = (\xi_{ij})$ and $\phi_2 = (\psi_{ij})$. Carrying out the matrix multiplication in $\phi = \phi_2 \phi_1$, we obtain n^2 equations

$$\sum_{k=1}^m \psi_{ik} \xi_{kj} = \begin{cases} \psi & \text{if } i=j \\ 0 & \text{otherwise.} \end{cases} \quad (2)$$

Note that the determinant of any m by m submatrix of ϕ_1 is the zero function. For otherwise using the Cramer's rule we can show that $\psi_{i1} = \dots = \psi_{im} = 0$ for some $i=1, \dots, n$, which, by (2), implies $\psi=0$, contradicting our assumption that $|\psi(e^{it})|=1$ on a set of positive Lebesgue measure. (A submatrix of ϕ_1 is obtained by crossing out some rows and columns from ϕ_1 .) On the other hand, consider the m by m matrix-valued functions

$$\phi_1' = \begin{pmatrix} \psi & & 0 \\ & \ddots & \\ 0 & & \psi \end{pmatrix}, \quad \phi_2' = \begin{pmatrix} \xi_{11} & \dots & \xi_{1m} \\ \vdots & & \vdots \\ \xi_{m1} & \dots & \xi_{mm} \end{pmatrix}$$

and $\phi' = \begin{pmatrix} \psi_{11} & \dots & \psi_{1m} \\ \vdots & & \vdots \\ \psi_{m1} & \dots & \psi_{mm} \end{pmatrix}$

Since $\phi' = \phi'_2 \phi'_1$, we have $\psi^m = \det \phi' = (\det \phi'_2)(\det \phi'_1) = 0$. This implies $\psi = 0$, again a contradiction. Thus we must have $m = n$, as asserted.

Proof of Theorem 1. To complete the proof we have to show that $K \in \text{Lat } S^{(n)}$. Note that in the functional model,

$$K = \{ \phi_2 u \oplus Z^{-1}(\Delta_2 u \oplus v) : u \in H^2(\mathbb{C}^n), v \in \overline{\Delta_1 L^2(\mathbb{C}^n)} \} \\ \oplus \{ \phi w \oplus \Delta w : w \in H^2(\mathbb{C}^n) \},$$

where Z denotes the unitary operator from $\overline{\Delta L^2(\mathbb{C}^n)}$ to $\overline{\Delta_2 L^2(\mathbb{C}^n)} \oplus \overline{\Delta_1 L^2(\mathbb{C}^n)}$ defined by

$$Z(\Delta v) = \Delta_2 \phi_1 v \oplus \Delta_1 v, \quad v \in L^2(\mathbb{C}^n)$$

Let $\phi_2 u \oplus t$ be an element in K , where $u = (u_i) \in H^2(\mathbb{C}^n)$ and $t = (t_i) \in \overline{\Delta L^2(\mathbb{C}^n)}$ satisfying $zt = \Delta_2 u \oplus v$, for some $v = (v_i) \in \overline{\Delta_1 L^2(\mathbb{C}^n)}$. We want to show that $S^{(n)}(\phi_2 u \oplus t) \in K$. Note that $\phi_2 u \oplus t = \sum_{j=1}^n \psi_{ij} u_j \oplus t_i$ and $\sum_{j=1}^n \psi_{ij} u_j \oplus t_i \in [H^2 \oplus \overline{\Delta L^2}] \oplus \{ \psi w \oplus W : w \in H^2 \}$. Hence

$$S^{(n)}(\phi_2 u \oplus t) = \sum_{i=1}^n \oplus [P \begin{pmatrix} A & 0 \\ B & C \end{pmatrix} \begin{pmatrix} \sum_{j=1}^n \psi_{ij} u_j \\ t_i \end{pmatrix}]$$

$$= \sum_{i=1}^n \oplus \left[P \begin{pmatrix} A \sum_{j=1}^n \psi_{ij} u_j \\ B \sum_{j=1}^n \psi_{ij} u_j + C t_i \end{pmatrix} \right]$$

$$= \sum_{i=1}^n \oplus \begin{pmatrix} A \sum_{j=1}^n \psi_{ij} u_j - \psi w_i \\ B \sum_{j=1}^n \psi_{ij} u_j + C t_i - \Delta w_i \end{pmatrix} \tag{3}$$

for some $W_i \in H^2$, $i=1,2,\dots,n$. Solving the following n equations for $\psi_{i1}, \dots, \psi_{in}$

$$\sum_{k=1}^n \psi_{ik} \xi_{kj} = \begin{cases} \psi & , \text{ if } j=i \\ 0 & , \text{ otherwise} \end{cases} \quad j=1,2,\dots,n,$$

by means of the Cramer's rule, we obtain

$$(\det \phi_1) \psi_{ik} = \psi \eta_{ik} \quad i,k=1,2,\dots,n,$$

where η_{ik} is the determinant, multiplied by $(-1)^{i+k}$, of the matrix obtained from ϕ_1 by deleting its i -th row and k -th column. It follows that

$$\begin{aligned} (\det \phi_1) B_j \sum_{i=1}^n \psi_{ij} u_j &= B_j \sum_{i=1}^n (\det \phi_1) \psi_{ij} u_j = B_j \sum_{i=1}^n \psi \eta_{ij} u_j \\ &= \sum_{i=1}^n \Delta(A-C) \eta_{ij} u_j = \Delta(A-C) \sum_{i=1}^n \eta_{ij} u_j. \end{aligned}$$

Note that $\sum_{j=1}^n \oplus [(\det \phi_1) B_j \sum_{i=1}^n \psi_{ij} u_j]$ is an element of $\overline{\Delta L^2(C^n)}$. Thus we have

$$\begin{aligned} Z \left(\sum_{i=1}^n \oplus [(\det \phi_1) B_j \sum_{i=1}^n \psi_{ij} u_j] \right) &= Z \left(\sum_{i=1}^n \oplus [\Delta(A-C) \sum_{j=1}^n \eta_{ij} u_j] \right) \\ &= \Delta_2 \phi_1^{(A-C)} \begin{pmatrix} \sum_{j=1}^n \eta_{1j} u_j \\ \vdots \\ \sum_{j=1}^n \eta_{nj} u_j \end{pmatrix} \oplus \Delta_1^{(A-C)} \begin{pmatrix} \sum_{j=1}^n \eta_{1j} u_j \\ \vdots \\ \sum_{j=1}^n \eta_{nj} u_j \end{pmatrix} \\ &= \Delta_2^{(A-C)} \begin{pmatrix} \sum_{k=1}^n \xi_{1k} & \sum_{j=1}^n \eta_{kj} u_j \\ \vdots & \vdots \\ \sum_{k=1}^n \xi_{nk} & \sum_{j=1}^n \eta_{kj} u_j \end{pmatrix} \oplus \Delta_1^{(A-C)} \begin{pmatrix} \sum_{j=1}^n \eta_{1j} u_j \\ \vdots \\ \sum_{j=1}^n \eta_{nj} u_j \end{pmatrix} \\ &= \Delta_2^{(A-C)} \begin{pmatrix} (\det \phi_1) u_1 \\ \vdots \\ (\det \phi_1) u_n \end{pmatrix} \oplus \Delta_1^{(A-C)} \begin{pmatrix} \sum_{j=1}^n \eta_{1j} u_j \\ \vdots \\ \sum_{j=1}^n \eta_{nj} u_j \end{pmatrix} \end{aligned}$$

$$= \Delta_2(A-C)(\det \phi_1)u \oplus \Delta_1(A-C) \begin{pmatrix} \sum_{j=1}^n \eta_{1j}u_j \\ \vdots \\ \sum_{j=1}^n \eta_{nj}u_j \end{pmatrix} \quad (4)$$

On the other hand,

$$\begin{aligned} & Z \left(\sum_{i=1}^n \oplus [(\det \phi_1) B_{j=1}^n \psi_{ij}u_j] \right) \\ &= (\det \phi_1) Z \left(\sum_{i=1}^n \oplus (B_{j=1}^n \psi_{ij}u_j) \right) \\ &= (\det \phi_1) (X \oplus Y), \end{aligned} \quad (5)$$

say, for some element $X \oplus Y$ in $\overline{\Delta_2 L^2(\mathbb{C}^n)} \oplus \overline{\Delta_1 L^2(\mathbb{C}^n)}$. Equating the first components in (4) and (5) we obtain

$$\Delta_2(A-C)(\det \phi_1)u = (\det \phi_1)X \quad (6)$$

Since $\det \phi = \psi^n \neq 0$ and $\det \phi = (\det \phi_2)(\det \phi_1)$, we have $\det \phi_1 \neq 0$.

Hence $\det \phi_1 \neq 0$ a.e. by the F. and M. Riesz Theorem. Thus (6) yields that

and
$$\Delta_2(A-C)u = X,$$

$$\begin{aligned} & Z \left(\sum_{i=1}^n \oplus (B_{j=1}^n \psi_{ij}u_j + Ct_i) \right) = Z \left(\sum_{i=1}^n \oplus (B_{j=1}^n \psi_{ij}u_j) \right) \\ & + Z \left(\sum_{i=1}^n \oplus Ct_i \right) = [(\Delta_2(A-C)u) \oplus Y] + C(\Delta_2 u \oplus v) \\ & = \Delta_2 Au \oplus (Y+Cv), \end{aligned}$$

Hence (3) can be written as

$$S^{(n)}(\phi_2 u \oplus t) = \{ \phi_2 Au \oplus Z^{-1} [\Delta_2 Au \oplus (Y+Cv)] \} \oplus (\phi_1 W \oplus \Delta W), \text{ where } W = (W_i) \in H^2(\mathbb{C}^n).$$

This shows that $S^{(n)}(\phi_2 u \oplus t) \in K$ as asserted and completes the proof.

COROLLARY 3. If T is the compression of the shift on the space $H^2 \ominus \psi H^2$, then T satisfies $\text{Alg } T = \{T\}'$.

If ψ is an outer function then more can be said.

THEOREM 4. Let T be a c.n.u. contraction having an outer scalar-valued characteristic function ψ . Then the following are equivalent to each other:

(1) $|\psi(e^{it})|=1$ on a set of positive Lebesgue measure;

(2) $\text{Alg } T = \{T\}'$;

(3) every invariant subspace for T is hyperinvariant, that is, invariant under every operator in $\{T\}'$.

The proof of this theorem follows immediately from the following lemma.

LEMMA 5. Let T be a c.n.u. contraction having an outer scalar-valued characteristic function ψ . Then T has the same lattice of hyperinvariant subspaces as the bilateral shift on $L^2(E)$, where $E = \{e^{it} : |\psi(e^{it})| < 1\}$.

In fact, Sickler [3] showed that in this case T has the same lattice of invariant subspaces as the bilateral shift on $L^2(E)$.

Our proof will be based on his descriptions of this lattice of T .

PROOF: The invariant subspaces of T are given by

$$K_S = \{f \oplus g \in H : -\Delta f + \psi g \in SH^2\} \quad K_F = \{f \oplus g \in H : -\Delta f + \psi g \in L^2(F)\}$$

where S runs over all unimodular functions and F over all measurable subsets of the unit circle (cf. [3], Theorem 5.1). In light of Sickler's result (Theorem 6.1 in [3]) it suffices to show that each K_F is hyperinvariant and each K_S is not. Since the operators in $\{T\}'$ are of the form

$$S = P \begin{pmatrix} A & 0 \\ B & C \end{pmatrix},$$

where $A \in H^\infty$ and $B, C \in L^\infty$ satisfy $B\psi + C\Delta = \Delta A$ a.e., we have

$$S \begin{pmatrix} f \\ g \end{pmatrix} = P \begin{pmatrix} A & 0 \\ B & C \end{pmatrix} \begin{pmatrix} f \\ g \end{pmatrix} = P \begin{pmatrix} Af \\ Bf + Cg \end{pmatrix} = \begin{pmatrix} Af - \psi w \\ Bf + Cg - \Delta w \end{pmatrix}$$

for some $w \in H^2$. Note that

$$\begin{aligned} -\Delta(Af - \psi w) + \psi(Bf + Cg - \Delta w) &= -\Delta Af + \psi Bf + \psi Cg \\ &= -(B\psi + C\Delta)f + \psi Bf + \psi Cg = C(-\Delta f + \psi g). \end{aligned}$$

Hence $f \oplus g \in K_F$ implies that $S(f \oplus g) \in K_F$, which proves our first assertion.

To prove the second assertion we may assume that $|\psi(e^{it})| < 1$ a.e., for otherwise all of the invariant subspaces are of the form K_F (cf. [3], Theorem 7.2). Assume that K_G is hyperinvariant for some unimodular function S . From the calculation above we conclude that for any operator S in $\{T\}'$ $-\Delta f + \psi g \in SH^2$ implies that $C(-\Delta f + \psi g) \in SH^2$. Since the set $\{-\Delta f + \psi g : f \oplus g \in K_G\}$ is dense in SH^2 (cf. the proof of Theorem 6.1 in [3]), the multiplication by C defines an operator on SH^2 . Thus $C \in H^\infty$. Let $\alpha_n = \{e^{it} : \frac{1}{n} \leq |\psi(e^{it})| \leq 1 - \frac{1}{n}\}$ for $n=2, 3, \dots$. Then $\bigcup_n \alpha_n =$ the unit circle. Consider $A=0$, $B_n = \chi_{\alpha_n}$ and $C_n = \frac{-\psi \chi_{\alpha_n}}{\Delta}$. It is easily seen unit $A \in H^\infty$ and $B_n, C_n \in L^\infty$ satisfy $B_n \psi + C_n \Delta = \Delta A$ a.e. Such choices of A , B_n and C_n give rise to an operator S_n in $\{T\}'$ for each n . Thus $C_n \in H^\infty$ for all n . Note that α_n^c , the complement of α_n with respect to the unit circle, has Lebesgue measure zero, for at least one $n \geq 2$, for otherwise $C_n = -\frac{\psi \chi_{\alpha_n}}{\Delta} = 0$ on α_n^c implies, by the F and M. Riesz Theorem and the fact that $\psi \not\equiv 0$, that α_n has Lebesgue measure zero for all n , which is absurd. Thus $\frac{1}{n} \leq |\psi| \leq 1 - \frac{1}{n}$ a.e. for some $n \geq 2$, and $C = \frac{-\psi}{\Delta}$ a.e. is an element of H^∞ . It follows that Δ is an element of H^∞ . Since Δ is real-valued, it can be easily shown that ψ is a constant function, the trivial case which we can exclude in the beginning. This completes the proof.

Proof of Theorem 4. The implication (3) \Rightarrow (1) is essentially given in the proof of Lemma 5.

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ABSTRACT — This paper is a review of the Gaussian laser beam propagation properties from the point of view of the Geometrical optics. It is meant to be tutorial in nature and useful in scope. And emphasis is placed on formulations, derivations and examples which lead to basic understanding and on results which bear practical significance.

I. Introduction

The field distribution of a fundamental mode gas laser beam is Gaussian. The equations to express the properties of Gaussian laser beam propagating in free space and along the axis on an optical system, in general, differ from the familiar first order equations of geometrical optics. In this paper we will discuss how the Gaussian beams of gas laser are transformed on their passage through free space (Sec. II), thin lens (Sec. III), spherical refracting lens surface (Sec. IV), spherical reflecting mirror surface (Sec. V), mode match (Sec. VI), and the transformation of the Gaussian beam in lens or lens-like systems (Sec. VII). And most of the above discussions are illustrated with numerical examples.

II. Propagation of Gaussian Laser Beam in a Homogeneous Medium [1-3]

In general the Gaussian spherical wave which propagates along the +Z direction can be written in the compressed form

$$(2-1) \quad \begin{aligned} \psi(x, y) &= \sqrt{\frac{2}{\pi}} \frac{1}{q} \exp(-jz - \frac{x^2 + y^2}{2q}) \\ &= \sqrt{\frac{2}{\pi}} \frac{1}{q} \exp(-jz - \frac{x^2 + y^2}{2q}) \end{aligned}$$

where q is implicitly defined as a complex radius of curvature, given

Since the set $\{e^{i\theta} : \theta \in \mathbb{R}\}$ is dense in S^1 (cf. the proof of Theorem 6.1 in [3]), the multiplication by C defines an operator on S^1 . Thus $C \in H^\infty$. Let $\alpha_n = \int_{-\pi}^{\pi} |e^{i\theta}|^2 |e^{i\theta}|^2 d\mu \leq 1 - \frac{1}{n}$ for $n=2,3,\dots$. Then $\alpha_n \rightarrow 1$ as $n \rightarrow \infty$. Consider $A=0$, $B = \alpha_n^{-1} x_{\alpha_n}$ and $C = \frac{-\beta x_{\alpha_n}}{\alpha}$. It is easily seen that A, B, C satisfy $A^2 + B^2 + C^2 = \lambda$ a.e. Such choices of A, B, C are possible for each α_n in $(T)'$ for each n . Thus $C \in H^\infty$ for all n . Note that $\alpha_n \rightarrow 1$ with respect to the unit circle. The Lebesgue measure zero, for at least one $s \geq 2$, for otherwise $C = \frac{-\beta x_{\alpha_n}}{\alpha} \rightarrow 0$ on S^1 implies, by the Fatou and Lebesgue Theorem and the fact that $\alpha_n \rightarrow 1$, that α_n has Lebesgue measure zero for all n , which is absurd. Thus $\frac{1}{n} < |e^{i\theta}|^2 < \frac{1}{n}$ a.e. for some $s \geq 2$, and $C = \frac{-\beta x_{\alpha_n}}{\alpha}$ a.e. is an element of H^∞ . It follows that λ is an element of H^∞ . Since λ is real-valued, it can be easily shown that λ is a constant function, the trivial case which we exclude in the beginning. This completes the proof.

Proof of Theorem 4. The implication (3) \Rightarrow (1) is essentially given in the proof of Lemma 5.

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