# Asymmetric Carrier Conduction Mechanism by Tip Electric Field in $WSiO_X$ Resistance Switching Device

Yong-En Syu, Ting-Chang Chang, Tsung-Ming Tsai, Geng-Wei Chang, Kuan-Chang Chang, Jyun-Hao Lou, Ya-Hsiang Tai, Ming-Jinn Tsai, Ying-Lang Wang, and Simon M. Sze

Abstract—Resistance random access memory (RRAM) is a great potential candidate for next-generation nonvolatile memory due to the outstanding memory characteristic. However, the resistance switching mechanism is still a riddle nowadays. In this letter, the switching mechanism is investigated by current–voltage (I-V) curve fitting in the TiN/WSiO<sub>X</sub>/Pt RRAM device. The asymmetric phenomenon of the carrier conduction behavior is explained at the high-resistance state in high electric field. The switching behavior is regarded to tip electric field by localizing the filament between the interface of top electrode and insulator.

*Index Terms*—Nonvolatile memory, resistance switching, tip electric field, tungsten silicide (WSi).

#### I. INTRODUCTION

**M**ODERN semiconductor nonvolatile memories are scaled constantly to achieve large capacity while device features are approaching the sub-100-nm regime. Nevertheless, for conventional charge-storage-based memories, such as Flash memory, the increasing demands for device densities by scaling

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Y.-E. Syu and J.-H. Lou are with the Department of Physics, National Sun Yat-Sen University, Kaohsiung 804, Taiwan (e-mail: syu.yong.en@gmail.com; bike731123@hotmail.com).

T.-C. Chang is with the Department of Physics, National Sun Yat-Sen University, Kaohsiung 804, Taiwan, and also with the Advanced Optoelectronic Technology Center, National Cheng Kung University, Tainan 701, Taiwan (e-mail: tcchang@mail.phys.nsysu.edu.tw).

T.-M. Tsai and K.-C. Chang are with the Institute of Materials Science and Engineering, National Sun Yat-Sen University, Kaohsiung 804, Taiwan (e-mail: tmtsai@faculty.nsysu.edu.tw; doubleccc@yahoo.com.tw).

G.-W. Chang and Y.-H. Tai are with the Department of Photonics and the Institute of Electro-Optical Engineering, National Chiao Tung University, Hsinchu 300, Taiwan (e-mail: b922030049@gmail.com; yhtai@mail.nctu. edu.tw).

M.-J. Tsai is with the Nanoelectronic Technology Division, Electronics and Optoelectronics Research Laboratories, ITRI, Hsinchu 31040, Taiwan (e-mail: mjtsai@itri.org.tw).

Y.-L. Wang is with the Taiwan Semiconductor Manufacturing Company, Hsinchu 300, Taiwan (e-mail: ylwang@tsmc.com).

S. M. Sze is with the Department of Physics, National Sun Yat-Sen University, Kaohsiung 804, Taiwan. He is also with the Department of Electronics Engineering, National Chiao Tung University, Hsinchu 300, Taiwan, and also with the Department of Electrical Engineering, Stanford University, Stanford, CA 94305-4085 USA (e-mail: simonsze@faculty.nctu.edu.tw).

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the dimension are expected to be a major challenge due to the technical and physical limitations. In order to overcome the issue, alternative memory technologies have been widely investigated [1]–[3]. Resistance random access memory (RRAM) is a great potential candidate for next-generation nonvolatile memory due to its superior properties such as simple structure, fast operation speed, low cost, and nondestructive readout [4]–[11]. The WSiO<sub>X</sub> is chosen for the RRAM switching layer because tungsten and silicon are compatible with the current complementary metal-oxide-semiconductor process. Heretofore, the resistance switching mechanism is unclear and difficult to confirm. In general, the resistance switching behavior is considered as a result of the formation and rupture of localized conduction filaments in oxide materials [10]. In this letter, the carrier transport behavior is analyzed by fitting the current-voltage (I-V) curves to identify the switching mechanism. The asymmetric carrier conduction behavior is observed at the high-resistance state (HRS) in high electric field and explained by the schematic energy band diagram. By the analysis results, the conduction filament theory can be verified.

### II. EXPERIMENT

The 30-nm WSiO<sub>X</sub> film was deposited on a TiN/SiO<sub>2</sub>/Si substrate as the resistance switching layer by RF magnetron sputtering using the WSi (W : Si = 1 : 1) target, in Ar/O<sub>2</sub> = 30 sccm/10 sccm mixed gas ambient with a working pressure of 6 mtorr. The Pt top electrode was deposited on an insulator WSiO<sub>X</sub> thin film by dc magnetron sputter. Finally, the photolithography and lift-off technique were employed to shape and complete the Pt/WSiO<sub>X</sub>/TiN memory cells. The electrical characteristics were measured for the memory devices with  $1-\mu m^2$  square cell area by an Agilent B1500 semiconductor parameter analyzer. During the electrical measurements, dc sweeping and pulse bias were applied to the bottom electrode (TiN) while the top electrode (Pt) was grounded.

# **III. RESULTS AND DISCUSSION**

The irreversible forming process that formed the conduction path is required to make the as-fabricated memory cells active, operating in dc voltage sweep mode with a compliance current of 100  $\mu$ A to prevent the device hard breakdown, shown in the inset of Fig. 1. The repeatable resistance switching



Fig. 1. Typical 100-cycle resistive switching I-V curves of Pt/WSiO<sub>X</sub>/TiN device. Inset shows the forming voltage of the device.



Fig. 2. Distributions of (a) the resistance states of HRS and LRS and (b) the switch voltage during 100 resistance switching cycles. The reliability of the device: (c) Resistance measurements after voltage pulse cycling and (d) retention behaviors of the HRS and LRS at 125  $^\circ$ C.

between HRS and low-resistance state (LRS) is obtained after the forming process. By sweeping the bias to negative over reset voltage ( $V_{\text{reset}}$ ), a gradual current decrease is presented because the conduction filaments are ruptured. The resistance state of the device transforms from LRS to HRS, called "reset process." Conversely, the conduction filaments are formed, and the cell turns back to LRS while applying a larger positive bias than set voltage ( $V_{\text{set}}$ ), i.e., "set process." A compliance current of 10 mA is assigned to prevent permanent breakdown. Fig. 1 shows the typical 100-cycle I-V characteristic of the Pt/WSiO<sub>X</sub>/TiN device by dc sweep voltage measurements, and the device exhibits the bipolar behavior by applying the bias on the TiN electrode.

The distributions of the resistance states and switch voltage are counted with continuous I-V sweep measurement of 100 cycles, as shown in Fig. 2(a) and (b). Resistance state distribution shows the stable resistance switching behaviors



Fig. 3. Curve fitting of carrier transport analysis in different bias regions and resistance states of  $WSiO_X$  switching layer.

and displays the large enough  $R_{\rm HRS}/R_{\rm LRS}$  ratio over two orders apart for the periphery circuits to probe the different resistance states. In addition, the switch voltage distribution shows the over 70% set voltage concentration from 1 to 1.2 V and 79% reset voltage concentration from -0.7 to -0.8 V. Furthermore, the endurance and retention of the  $Pt/WSiO_X/TiN$ device are necessary to investigate for nonvolatile memory application. Fig. 2(c) shows the resistance of the HRS and LRS versus the number of pulses. The HRS was set to the LRS by applying a setting voltage of 2 V for 550 ns, while the LRS was reset to HRS by applying -2 V for 2  $\mu$ s. Both HRS and LRS resistance values were extracted at 0.1 V. The RRAM device can be repeatedly operated over 10<sup>5</sup> cycles, and the  $R_{\rm HRS}$  and  $R_{\rm LRS}$  remain two orders apart. Fig. 2(d) shows the variation of the resistance with time at both the HRS and LRS for the Pt/WSiO<sub>X</sub>/TiN device at 125 °C. After 10000 s, the variation of HRS and LRS is very slight. It indicates that the Pt/WSiO<sub>X</sub>/TiN device has both excellent reliability and durability.

In order to ascertain the switching mechanism, the I-Vcurves are fitted to analyze carrier transport of the switching layer, as shown in Fig. 3. Moreover, we explained the mechanism by energy band diagram, as shown in Fig. 4. The carrier transport mechanism in LRS is considered as ohmic conduction because the carriers are transported through the continuous filament. The carrier transport mechanism in HRS with lower electric field is dominated by ohmic conduction because the current originates from the internal carrier in the  $WSiO_X$ , as shown in Fig. 4. The bias rises continuously, and the relation between  $\ln(I/T^2)$  and the square root of the applied voltage  $(V^{1/2})$  is linear. Schottky emission is considered as the main transport mechanism [12]. The major leakage current results from the electron leap across the potential energy barrier by the thermionic effect [as shown in Fig. 4(b) and (e)], while the energy band bending is insignificant. Therefore, the carrier transport mechanism at this bias region complied with the Schottky emission mechanism.

The asymmetric phenomenon of the carrier conduction behavior is found at HRS in high electric field. At the positive-bias



Fig. 4. Corresponding schematic energy band diagram in different bias regions and resistance states.

region, the current transportation mechanism is space-chargelimited current (SCLC). The current versus voltage (I-V) was plotted in measure range in inset (c) of Fig. 3 with a double logarithmic scale, which shows a linear dependence on voltage with a slope of 1.81 [12]. At the negative-bias region, the current transportation mechanism is Fowler-Nordheim (F-N) tunneling. The  $\ln(1/V^2)$  versus  $V^{1/2}$  was plotted in measure range in inset (f) of Fig. 3, which shows a linear dependence on voltage [12]. From the results of curve fitting analysis, we suppose that the asymmetric carrier conduction behavior is caused by tip electric field effect. The carrier conduction behavior in the HRS is explained by the schematic energy band diagram, as shown in Fig. 4(c) and (f). As shown in Fig. 4(c), at the positive-bias region, the tip filament would cause a strong electrical effect. A significant energy band bending at the W filament tip causes considerable quantity of electrons tunneling from the Pt/W top electrode through the narrow energy barrier area into  $WSiO_X$ . However, the mitigating electric field near the flat TiN electrode would cause the electrons to accumulate in  $WSiO_X$ . The conduction of tunneling electrons would be limited so as to form the carrier transport mechanism of SCLC. The expectable F–N tunneling is observed in the high-electricfield region [as shown in Fig. 4(f)]. Because of the enough large electric field, electrons can be injected from the TiN bottom electrode through the triangle potential barrier of  $WSiO_X$  into the Pt/W top electrode. Moreover, the tunneling electrons will not accumulate because the electrons can be transported easily by the strong electric field near the W filament tip. Therefore, the carrier conduction behavior in the HRS of positive bias is different from that in the HRS of negative bias.

#### **IV. CONCLUSION**

In summary, this letter has investigated the bipolar resistance switching characteristics of the tungsten-based RRAM with the Pt/WSiO<sub>X</sub>/TiN structure. The switching properties and reliability results indicate that the device possesses excellent memory characteristics. Additionally, after I-V curves were fitted to analyze carrier transport of the switching layer, the asymmetric phenomenon of the carrier conduction behavior is found at the HRS in high electric field. It is illustrated that the asymmetric carrier conduction behavior is generated due to the tip electric field effect expressed by the schematic energy band diagram.

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