



## Numerical analysis of nanoaluminum combustion in steam



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### ABSTRACT

The comprehensive analysis of chain mechanism development in the Al–H<sub>2</sub>O system is performed on the base of novel reaction mechanism taking into account quantum chemistry studies of potential energy surfaces of the elementary reactions with Al-containing species and estimations of rate constants of corresponding reaction channels. As well the physical properties of Al-containing species involved in the reaction mechanism and needed for the calculation of their transport coefficients are reported. The developed reaction mechanism makes it possible to describe with reasonable accuracy the experimental data on ignition temperature in Al–O<sub>2</sub>–Ar and Al–H<sub>2</sub>O systems and obtain the qualitative agreement with measured value of laminar flame speed. The two-stage regime of ignition in the Al–H<sub>2</sub>O reacting system was revealed both when the aluminum is in the liquid phase and when it comes into steam environment in the gas phase. It was shown that decreasing the ignition temperature one can increase the hydrogen yield in the combustion exhaust.

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### 1. Introduction

Water vapor is a strong oxidizer for several metals. So, the reaction of aluminum with steam yields very high energy density. The other important issue is that the combustion of Al + H<sub>2</sub>O mixture produces a significant amount of molecular hydrogen as a combustion product, and reaction of Al with H<sub>2</sub>O is considered as a very promising method of hydrogen production [1–3]. Though aluminum is widely utilized as a fuel ingredient in solid propellants and explosives to increase the heat release and, as a result, the specific impulse of rocket engines as well as to reduce the combustion instability [4–6], and a number of researches was addressed the study of kinetics of aluminum combustion (see, for example, [7–11]), the majority of previous works were focused on micron-sized particle burning. The main disadvantage of such particle combustion is that, in this case, large particle can burn only due to diffusion of oxidizer through the metallic oxide shell that forms on the surface of particle, and rate of combustion is limited by the rate of this process [12].

The novel very attractive topic in the combustion of energetic materials is the use of nano-sized Al particles. Such particles exhibit the distinguishing properties compared with micron-sized ones. They feature lower ignition temperature and faster burning rate [8,13–15]. It is believed that such particles undergo fast gasification, and, in this case, atomic aluminum reacts with oxidizer in

a gas phase [11]. This supposition has been recently proved by the detection of atomic aluminum during ignition of Al nanoparticles with  $d = 80$  nm behind the reflected shock wave [16]. The particles with larger size can be in a liquid phase, and, frequently (see, for example, [17]), the formation of aluminum in gas phase is modeled by simple one step reaction  $\text{Al}(l) \leftrightarrow \text{Al}(g)$ . In line with the melt-dispersion mechanism [18], at high rate of particle heating, the melting of Al inside the particle is accompanied by volume expansion which creates large pressure in the molten Al. This leads to alumina shell dynamic spallation. As a result, Al liquid core is dispersed into small bare clusters that fly with high velocity. Oxidation of these clusters is not limited by diffusion through the initial alumina shell and occurs in a kinetic mode.

Therefore, it is of great importance to build an adequate gas phase reaction mechanism describing properly the ignition and combustion in the Al–H<sub>2</sub>O mixture. It should be emphasized that, in past decade, only a few studies [11,17] were addressed modeling the processes in the Al–H<sub>2</sub>O system. However, for the past years, the novel reaction paths in Al–H<sub>2</sub>O system were revealed on the base of theoretical calculations [19–21]. This makes it possible to build a more reliable kinetic model for the ignition and combustion of aluminum in steam environment and reconsider the features of chain mechanism development in such a system. The present paper is focused on solving precisely these problems.

### 2. Kinetic model

The simple estimations show that, for nonoxidized particles with the radius  $r < 25$  nm, the time of heating of particle core

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caused by energy release in the course of aluminum oxidation and formation of aluminum oxide ( $\text{Al}_2\text{O}_3$ ) shell is much higher than the characteristic time of phase transformation. This leads to the rapid destruction of surface layer consisting of  $\text{Al}_2\text{O}_3$ . This conclusion was proved by molecular dynamic calculations [22]. Therefore, the model assumes that nano-sized particles boil or gasify very rapidly and aluminum comes into the environment in the gas or liquid phase. Note that, in accordance with the results of work [18], small liquid Al clusters can form during oxidation even larger sized particles. Certainly, such a model can be applied only to nonoxidized particles.

Nowadays, all developed kinetic mechanisms treated the reaction of atomic aluminum with  $\text{H}_2\text{O}$  molecule as a principal chain initiation reaction. In line with the data reported in [23], both the reaction mechanism of Huang et al. [17] and the mechanism of Washburn et al. [11] considered this reaction path as a complex processes consisting of two channels



and



and the rate constant of R1 channel at high temperature ( $T \geq 1500$  K) is more than 15 times smaller than that of R2 channel.

However, recent theoretical studies [19,20] revealed that the main reaction path of aluminum with  $\text{H}_2\text{O}$  molecule is the R1 channel, whereas the R2 reaction path is not an elementary one, and rate constant of R1 path is much higher than that of R2 one. Moreover, the reaction mechanisms of Washburn [11] and Huang et al. [17] do not include the important reaction paths with  $\text{Al}(\text{OH})_2$ ,  $\text{Al}(\text{OH})_3$ ,  $\text{AlO}_2\text{H}$  and other species that were investigated theoretically in recent studies [21,24]. Note that these species were also included in the consideration upon thermodynamic analysis of the composition of combustion exhaust in the premixed  $\text{Al-H}_2\text{-O}_2\text{-N}_2$  flame [25]. In our reaction mechanism we included all these reactions.

In addition, we involved the novel reaction  $\text{AlO} + \text{AlH} = \text{AlOH} + \text{Al}$ . The rate constant for this reaction was calculated on the base of *ab initio* study of potential energy surface with the use of program package Firefly V. 7.1.G [26]. In accordance with thermochemistry, the activation energy of the reaction  $\text{AlOH} + \text{Al} = \text{AlO} + \text{AlH}$  must not be lower than 21,200 K. The preliminary quantum-chemical calculations with the use of the unrestricted variant of the second-order Møller–Plesset perturbation theory (UMP2) [26] for solving the Schrödinger equation in relatively large 6-311+G(d,p) basis set revealed that the backward process  $\text{AlO} + \text{AlH} \rightarrow \text{AlOH} + \text{Al}$  is nearly barrierless ( $E_a \sim 670$  K). The application of more sophisticated quantum-chemical methods that treat electronic correlation more accurately proved this supposition. So, when applying the fourth-order Møller–Plesset perturbation theory in the MP4(SDQ) modification [26] with the same basis set, a slightly lower value of activation energy ( $E_a \sim 440$  K) for  $\text{AlO} + \text{AlH} = \text{AlOH} + \text{Al}$  was obtained. However, when utilizing the coupled-cluster method for solving the Schrödinger equation that proved to be one of the most accurate predictive tools in quantum chemistry, even zero activation barrier value was revealed. So, the calculations using the unrestricted coupled-cluster method including single, double and triple excitations (UCCSD(T)) [26] with 6-311+G(d,p) basis set, suggest that the transition state of this reaction lies more than 4000 K lower than  $\text{AlO} + \text{AlH}$ . Thus, it is worthwhile to include the reaction  $\text{AlO} + \text{AlH} = \text{AlOH} + \text{Al}$  in the present kinetic model with the rate constant  $k(T) = 2.54 \times 10^{13} T^{0.17} \text{ cm}^3 \text{ mol}^{-1} \text{ s}^{-1}$  determined in accordance with methodology used for barrierless reaction [21].

In order to describe the formation of liquid  $\text{Al}_2\text{O}_3$ , in according with [8,17] we treated the reaction of the transformation of

gaseous  $\text{Al}_2\text{O}_3$  to the liquid phase  $\text{Al}_2\text{O}_3(\text{g}) \leftrightarrow \text{Al}_2\text{O}_3(\text{l})$  with the rate constant of  $10^{14} \text{ s}^{-1}$ . Besides these reaction paths, the developed kinetic model must comprise the submechanism that describes the reactions in the  $\text{H}_2\text{-O}_2$  system. Earlier, the authors built the kinetic mechanism, which reproduces the large set of experimental data in  $\text{H}_2\text{-O}_2\text{-Ar}$  and  $\text{H}_2\text{-air}$  mixtures with rather high accuracy [27,28]. Precisely this reaction mechanism was utilized as a submechanism in the kinetic model for the ignition and combustion in the  $\text{Al-H}_2\text{O}$  system. Thus, the developed reaction mechanism comprises 59 reversible reactions with following species: O,  $\text{O}_2$ , H,  $\text{H}_2$ , OH,  $\text{H}_2\text{O}$ ,  $\text{HO}_2$ ,  $\text{H}_2\text{O}_2$ , Al, AlO,  $\text{AlO}_2$ ,  $\text{Al}_2\text{O}$ ,  $\text{Al}_2\text{O}_2$ ,  $\text{Al}_2\text{O}_3$ ,  $\text{Al}_2\text{O}_3(\text{l})$ , AlH,  $\text{AlH}_2$ ,  $\text{AlH}_3$ , AlOH,  $\text{Al}(\text{OH})_2$ ,  $\text{Al}(\text{OH})_3$ ,  $\text{AlO}_2\text{H}$ . Table 1 lists the elementary reactions involved in the model and coefficients needed for the calculations of forward reaction rate constants  $k_q^+$  in line with Arrhenius formula  $k_q = A_q T^{n_q} \exp(-E_q/T)$ , where  $A_q$  is the Arrhenius coefficient,  $E_q$  is the activation energy of  $q$ th reaction and  $n_q$  is the power coefficient. The coefficients for the reactions with Al-containing species presented in Table 1 were chosen on the base of recommendations [8,11,20,21,29]. It should be emphasized that we reconsidered the rate constants for the reactions



reported previously by Swihart et al. [29]. The rate constants were estimated on the base of the simplified model of triple collisions for the reaction (R3) and in accordance with methodology used for barrierless reaction for the reaction (R4) [21]. The rate constants of backward reactions  $k_q^-$  were calculated with the use of detailed balancing principle. Thermodynamic properties of Al-containing species were taken from the databases reported in [21,30].

It should be emphasized that for some species there are no data on their transport properties. These data are strongly needed to model the deflagrative mode of combustion and the speed of laminar flame propagation. In this work, the coefficients of molecular diffusion, thermal conductivity and viscosity were estimated with the usage of known formulas [31]. The ionization energy  $E_i$ , dipole moment  $\mu_i$ , polarizability  $\alpha_i$  and Van-der-Waals collision diameters  $\sigma_i$ , needed for these estimations, were taken from [21,32] or obtained on the base of density functional theory (DFT) calculations at the B3LYP/6-311+G(d,p) level of theory [26]. These parameters as well as the estimated well depth of Lennard-Jones potential for Al-containing species are presented in Table 2.

As in [8,17], we supposed that transport properties of condensed phase of Al and  $\text{Al}_2\text{O}_3$  are identical to their gas phase counterparts. Though, it should be emphasized that this supposition can be invalid, and the transport properties of rather large clusters  $(\text{Al})_n$  and  $(\text{Al}_2\text{O}_3)_n$  with  $n > 4$ , that form the liquid phase, can differ substantially from those of Al and  $\text{Al}_2\text{O}_3$  monomers. This issue requires special consideration on the basis of *ab initio* calculations of structure and physical properties of such clusters.

### 3. Results and discussions

Despite the long history of the studies of aluminum particle burning, the researches of the features of reaction in aluminum–air and aluminum–steam mixtures with the use of detailed gas phase reaction mechanism are fairly limited. These researches allow one to gain an insight both in nano- and in micro-sized aluminum particle combustion. Because the behavior of reaction in Al–air and Al– $\text{H}_2\text{O}$  mixtures was examined earlier by Huang et al. [8,17], it is reasonable to consider the features of chain mechanism development in such systems, investigated on the base of updated kinetic model, in comparison with that reported in

**Table 1**

List of reactions and coefficients for calculation of rate constants (cm, mole, s, K).

No.	Reaction	$k_q^+$			$k_q^-$			Ref.
		$A_q$	$n_q$	$E_q$	$A_q$	$n_q$	$E_q$	
1.	$H_2O + H = OH + H_2$	8.4E+13	0	10,116	2E+13	0	2600	[27]
2.	$O_2 + H = OH + O$	2.2E+14	0	8455	1.3E+13	0	350	[27]
3.	$H_2 + O = OH + H$	1.8E+10	1	4480	8.3E+9	1	3500	[27]
4.	$O_2 + M = O + O + M$	5.4E+18	-1	59,400	6.0E+13	0	-900	[27]
5.	$H_2 + M = H + H + M$	2.2E+14	0	48,300	9E+17	-1	0	[27]
6.	$H_2O + M = OH + H + M$	1E+24	-2.2	59,000	2.2E+22	-2	0	[27]
7.	$OH + M = O + H + M$	8.5E+18	-1	50,830	7.1E+18	-1	0	[27]
8.	$H_2O + O = OH + OH$	5.8E+13	0	9059	5.3E+12	0	503	[27]
9.	$H + O_2(+M) = HO_2(+M) k_0$	3.5E+16	-0.41	-565				[27]
	$F_c = 0.5$	1.48E+12	0.6	0				
	$k_\infty$	7.39E+5	2.43	26,926				[27]
10.	$H_2 + O_2 = H + HO_2$	4.76E+11	0.372	28,743	1E+13	0	540	[27]
11.	$H_2O + O = H + HO_2$	1.5E+15	0.5	36,600	3E+14	0	0	[27]
12.	$H_2O + O_2 = OH + HO_2$	1.2E+13	0	20,200	2.5E+14	0	950	[27]
13.	$OH + OH = H + HO_2$	1.3E+13	0	28,200	5E+13	0	500	[27]
14.	$H + H_2O_2 = H_2 + HO_2$	1.7E+12	0	1900	6E+11	0	9300	[27]
15.	$H + H_2O_2 = H_2O + OH$	5E+14	0	5000	2.4E+14	0	40,500	[27]
16.	$HO_2 + HO_2 = H_2O_2 + O_2$	1.8E+13	0	500	3E+13	0	21,600	[27]
17.	$HO_2 + H_2O = H_2O_2 + OH$	1.8E+13	0	15,100	1E+13	0	910	[27]
18.	$OH + HO_2 = H_2O_2 + O$	5.2E+10	0.5	10,600	2E+13	0	2950	[27]
19.	$H_2O_2 + M = OH + OH + M$	1.2E+17	0	22,900	9.1E+14	0	-2650	[27]
20.	$Al + O + M = AlO + M$	3E+17	-1	0				[11]
21.	$H_2O/2.8/ O_2/1.1/ H_2/1.1/$	2.31E+13	0.17	0				[21]
22.	$Al + O_2 = AlO + O$	7.12E+12	0.5	13,150				[21]
23.	$AlO + O_2 = AlO_2 + O$	1E+15	0	44564.6				[11]
24.	$AlO_2 = AlO + Al$	1E+15	0	67035.7				[11]
25.	$Al_2O = AlO + AlO$	1E+15	0	59335.7				[11]
26.	$Al_2O_2 = Al + AlO_2$	1E+15	0	74937.1				[11]
27.	$Al_2O_2 = Al_2O + O$	1E+15	0	52,466				[11]
28.	$Al_2O_3 = Al_2O_2 + O$	3E+15	0	49144.4				[11]
29.	$Al_2O_3 = AlO_2 + AlO$	3E+15	0	63915.4				[11]
30.	$Al + H_2O = AlOH + H^a$	1.96E+14	-0.09	3744				[20]
		2.78E+06	2.06	438				
31.	$Al + 2H_2O = Al(OH)_2 + H_2^b$	1.16E+15	0.5	0				[24], estimated
32.	$Al + 2H_2O = AlOH + H + H_2O^b$	1.16E+15	0.5	1260				[24], estimated
33.	$Al + HO_2 = AlO + OH$	1.33E+13	0.17	0				[21]
34.	$Al + HO_2 = AlH + O_2$	1.33E+13	0.17	0				[21]
35.	$AlO + H + M = AlOH + M$	1.99E+15	0.5	0				[21]
36.	$Al + OH + M = AlOH + M$	2.16E+15	0.5	0				[21]
37.	$AlOH + H = AlO + H_2$	2.66E+08	0.82	7844				[20]
38.	$Al + H + M = AlH + M$	9.43E+14	0.5	0				estimated
39.	$AlH_2(+M) = AlH + H(+M) k_0$	9.68E+14	0	19,962				[29]
	$k_\infty$	1.46E+15	0	23,376				
	$F_c = -4.1 * \exp(-T/21.6) + 5.1 * \exp(-T/493) + \exp(-942/T)$							
40.	$AlH_3(+M) = AlH + H_2(+M) k_0$	1.01E+15	0	27,089				[29]
	$k_\infty$	1.48E+13	0	30,756				
	$F_c = 0.94 * \exp(-T/885) + 0.06 * \exp(-T/552) + \exp(-3807/T)$							
41.	$AlH + H = Al + H_2$	7.15E+13	0.17	0				estimated
42.	$AlH_2 + H = AlH + H_2$	2E+13	0	0				[29]
43.	$AlH_3 + H = AlH_2 + H_2$	4.75E+09	1.5	0				[29]
44.	$AlOH + O = AlO + OH$	7.53E+12	0.5	4450				[21]
45.	$AlO + OH + M = AlO_2H + M$	2.62E+15	0.5	0				[21]
46.	$AlO + AlH = AlOH + Al$	2.54E+13	0.17	0				estimated
47.	$AlO_2 + H + M = AlO_2H + M$	2.19E+15	0.5	0				[21]
48.	$AlOH + O + M = AlO_2H + M$	2.94E+15	0.5	0				[21]
49.	$AlO + HO_2 = AlOH + O_2$	2.19E+14	-0.08	-35				[21]
50.	$AlOH + HO_2 = AlO_2H + OH$	3.76E+13	0.14	0				[21]
51.	$AlO_2 + H_2O = AlO_2H + OH$	2.63E+02	3.26	3430				[21]
52.	$AlO_2 + H_2 = AlO_2H + H$	4.62E+10	1.39	2940				[21]
53.	$AlO_2 + OH = AlO_2H + O$	2.57E+13	0.17	0				[21]
54.	$AlO_2H + O = AlOH + O_2$	2.14E+13	0.17	0				[21]
55.	$AlO_2H + H = Al(OH)_2$	7.39E+13	0.17	0				[21]
56.	$AlOH + OH = Al(OH)_2$	2.48E+13	0.16	-23				[21]
57.	$Al(OH)_2 + OH = Al(OH)_3$	2.54E+13	0.15	-48				[21]
58.	$Al_2O_3(g) = Al_2O_3(l)$	1E+14	0	0				[8]

<sup>a</sup> The rate constant is the sum of two Arrhenius dependencies.<sup>b</sup> Reaction products were analyzed in [24], the rate constant was estimated on the base of the simplified model of triple collisions [21].

**Table 2**  
Parameters taken for the calculation of species transport properties used in CHEMKIN Program, as well as ionization energy needed for the calculation of Lennard-Jones potential well depth.

Species	Molecular geometry <sup>a</sup>	Lennard-Jones potential well depth <sup>b</sup> , $\epsilon/k_B$ (K)	Lennard-Jones collision diameter <sup>c</sup> , $\sigma$ (Å)	Dipole moment, $\mu$ (Debye)	Polarizability, $\alpha$ (Å <sup>3</sup> )	Rotational relaxation collision number $Z_{rot}$ at 298 K	Ionization energy, $E_i$ (eV)
Al	0	2836	2.655	0 <sup>d</sup>	8.34 <sup>d</sup>	1	5.986 <sup>d</sup>
AlO	1	1381.2	3.203	4.32 <sup>d</sup>	6.7 <sup>d</sup>	1	9.46 <sup>d</sup>
AlO <sub>2</sub>	1	984	3.564	0 <sup>c</sup>	9.33 <sup>c</sup>	1	9.71 <sup>c</sup>
Al <sub>2</sub> O	1	2231	3.344	0	12.6	1	8.3
Al <sub>2</sub> O <sub>2</sub>	2	915.2	4.089	5.9	11.1	1	9.5
Al <sub>2</sub> O <sub>3</sub>	1	352	4.264	0	9.5	1	10
AlH	1	1101.8	2.975	0.177 <sup>d</sup>	6.264 <sup>d</sup>	1	8.233 <sup>d</sup>
AlH <sub>2</sub>	2	526.8	3.184	0.481 <sup>d</sup>	5.645 <sup>d</sup>	1	7.222 <sup>d</sup>
AlH <sub>3</sub>	2	119.8	4.001	0 <sup>d</sup>	4.39 <sup>d</sup>	1	10.793 <sup>d</sup>
AlOH	2	668.7	3.374	0.81 <sup>c</sup>	6.71 <sup>c</sup>	1	9.04 <sup>c</sup>
Al(OH) <sub>2</sub>	2	133.2	4.224	1.32 <sup>c</sup>	6.43 <sup>c</sup>	1	7.22 <sup>c</sup>
Al(OH) <sub>3</sub>	2	48.5	4.526	0	4.27	1	9.68
AlO <sub>2</sub> H	2	216	4.217	5.34 <sup>c</sup>	5.19 <sup>c</sup>	1	9.8 <sup>c</sup>

<sup>a</sup> 0 – atom, 1 – linear molecule, 2 – nonlinear molecule.

<sup>b</sup> Lennard-Jones potential well depth calculated in line with CHEMKIN format.

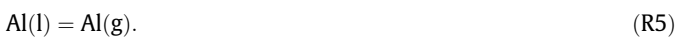
<sup>c</sup> Data were taken from [21].

<sup>d</sup> Data were taken from NIST Computational Chemistry Comparison and Benchmark Database [32].

[8,17]. All computations were carried out using the CHEMKIN Program package and the PREMIX subroutine [33].

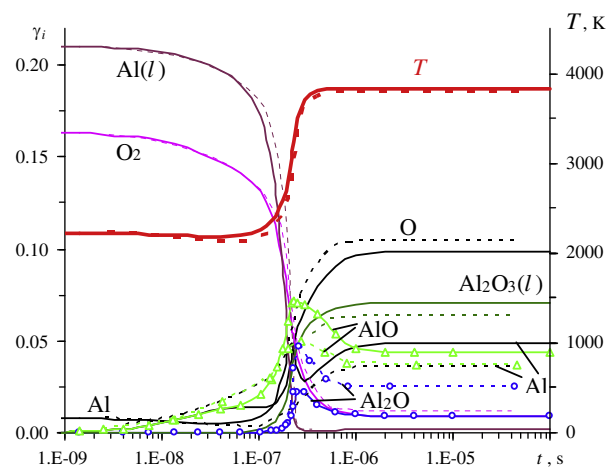
### 3.1. Ignition in the Al–air mixture

The main supposition of this work model as well as the model of Huang et al. [8] is that the liquid aluminum comes through the spalled alumina shell into environment. Certainly this supposition is needed for some proofs. Following [8] we assumed that liquid aluminum passes into gas phase in the course of reaction



Huang et al. introduced the rate constant for backward reaction  $k_{-R5}(T) = 10^{14} \text{ s}^{-1}$ . The rate constant of forward reaction (R5) was estimated in line with thermochemistry. Following the reference [8] we modeled the ignition process in adiabatic constant-pressure environment. Note that preliminary theoretical calculations exhibited that the rate constant of the reaction  $\text{Al} + \text{N}_2 = \text{AlN} + \text{N}$  (its energy barrier is approximately equal to 80,000 K or even higher) is much smaller than that for the reaction  $\text{Al} + \text{O}_2 = \text{AlO} + \text{O}$  and, therefore, it can be assumed that reactions with  $\text{N}_2$  molecules do not influence the chain mechanism of aluminum oxidation and nitrogen can be considered as a diluent. Though, it is worth noting that the reactions involving N-containing species ( $\text{N}_2$ , N, NO,  $\text{NO}_2$ ,  $\text{N}_2\text{O}$ , etc.) should be included in the reaction mechanism in order to describe the formation of aluminum nitrides and  $\text{NO}_x$  in the combustion exhaust. However, this is out the scopes of the present paper. Figure 1 shows the temporal profiles of species mole fractions and temperature calculated in [8] and computed with the usage of the model of the present work during ignition of a stoichiometric (fuel/air equivalence ratio  $\phi = 1$ ) Al(l)–air mixture at initial temperature  $T_0 = 2300 \text{ K}$  and pressure  $P_0 = 1 \text{ atm}$ . As in [8], we supposed that aluminum is in a liquid phase. As is seen, the curves predicted by both kinetic models are very close to each other. Some differences are observed for the concentrations of Al, AlO,  $\text{Al}_2\text{O}$ . This appears to be caused by the difference in thermodynamic properties of  $\text{Al}_2\text{O}$  molecule utilized in kinetic model [8] and the model of the present work.

Recently in [16] the experimental data on the ignition of oxidized 80 nm Al particles in  $\text{Al}/\text{O}_2/\text{Ar}$  mixture obtained with the shock tube technique were reported. The authors of this work detected the ignition of such particles at  $T \geq 1200 \text{ K}$  and pressure  $P_0 = 7 \text{ atm}$  in the shock tube with working time of 1–2 ms. Shown



**Fig. 1.** Evolution of species mole fractions and temperature during ignition in the stoichiometric Al(l)–air mixture calculated in [8] (dotted curves) and computed with the usage of the present work model (solid curves).  $T_0 = 2300 \text{ K}$ ,  $P_0 = 1 \text{ atm}$ .

in Fig. 2 is the variation of temperature vs. time in stoichiometric Al(l)/ $\text{O}_2$ /Ar mixture at different values of initial temperature computed with the use of the model of this work assuming the aluminum in a liquid phase. One can see that the model predicts the ignition at  $t \sim 1 \text{ ms}$  at the temperature  $T = 1100 \text{ K}$  that is close to  $T_{ign} = 1200 \text{ K}$  reported in [16].

### 3.2. Flame propagation in the Al–air mixture

Shown in Fig. 3 are the spatial profiles of species mole fractions and temperature in the laminar Al(l)–air stoichiometric ( $\phi = 1$ ) atmospheric pressure flame predicted by both kinetic models considered. And again there exists some discrepancy between the Al, AlO and  $\text{Al}_2\text{O}$  mole fractions calculated with the usage of kinetic model [8] and the model of the present work. However, the values of flame speed  $U_n$  predicted by these models are similar. The calculation of Huang et al. [8] gives  $U_n = 5.8 \text{ m/s}$  and the model of this work predicts  $U_n = 5.6 \text{ m/s}$ . This is not surprised because the Al–O submechanisms involved in the model of the present work and in the model of Huang et al. [8] are very close to each other and provide the identical profiles of atomic oxygen concentration in the flame front. So as the laminar flame speed is determined mainly by the rate of diffusion of the lightest mixture component,

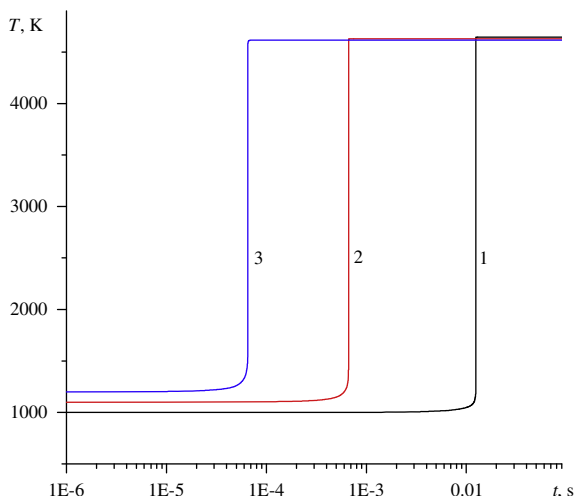


Fig. 2. Temporal profiles of temperature in stoichiometric Al(l)/O<sub>2</sub>/Ar mixture (40% O<sub>2</sub>/60% Ar) at P<sub>0</sub> = 7 atm and different initial temperatures T<sub>0</sub> = 1000, 1100 and 1200 K (curves 1–3).

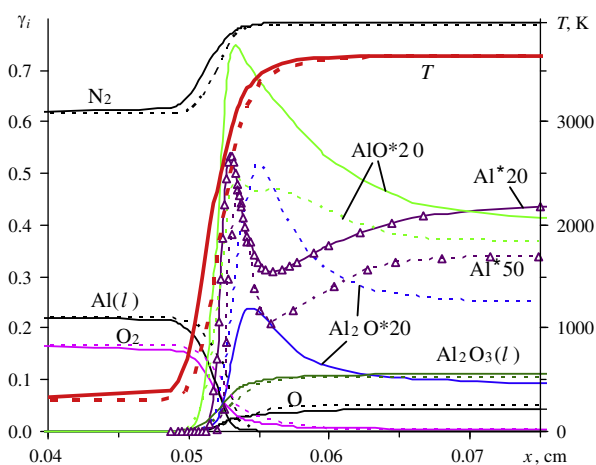


Fig. 3. Variation of species mole fractions and temperature in the stoichiometric Al(l)-air laminar flame calculated in [8] (dotted curves) and computed with the usage of the present work model (solid curves). T<sub>0</sub> = 300 K, P<sub>0</sub> = 1 atm.

atomic oxygen, from the hot flame region to the cold region, one can expect that the U<sub>n</sub> values predicted by these two models should be very close.

The sensitivity analysis demonstrated (see Fig. 4) that, both in the model [8] and in the model of the present work, the U<sub>n</sub> value is the most sensitive to the rate constant of the reaction Al + O<sub>2</sub> = AlO + O. The second important reaction in model [8] is Al<sub>2</sub>O = AlO + Al, while in the model of the present work, the other two reactions Al<sub>2</sub>O<sub>3</sub> = Al<sub>2</sub>O<sub>2</sub> + O and Al<sub>2</sub>O<sub>2</sub> = 2AlO contribute significantly to the sensitivity of U<sub>n</sub> value.

### 3.3. Ignition in the Al–H<sub>2</sub>O system

As was examined above for the aluminum–air mixture, further we will consider the ignition of liquid aluminum in steam environment. Shown in Fig. 5 is the variation of temperature vs. time for the atmospheric stoichiometric (φ = 1) Al(l)–H<sub>2</sub>O system at different T<sub>0</sub> values calculated with the use of the model of the present work and model [17]. One can see that the predictions of the kinetic models under study differ significantly. So, at high initial temperatures (T<sub>0</sub> ≥ 2000 K), the model [17] predicts shorter ignition delay than that of the model of the present work, and, at smaller temperature, the situation is inverted. Moreover, the

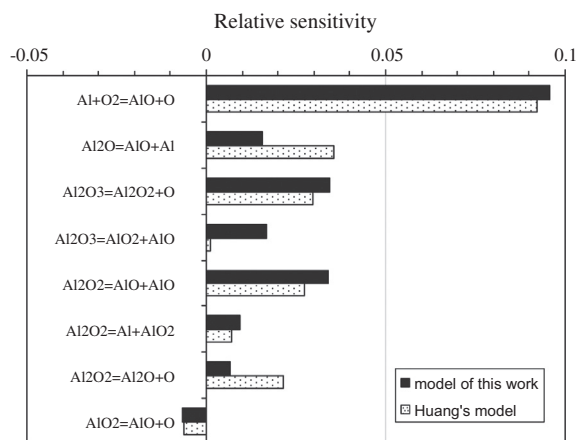


Fig. 4. Sensitivity of laminar flame speed to the reaction rates for the stoichiometric Al(l)-air mixture at T<sub>0</sub> = 300 K, P<sub>0</sub> = 1 atm calculated in [8] and computed with the usage of the model of the present work.

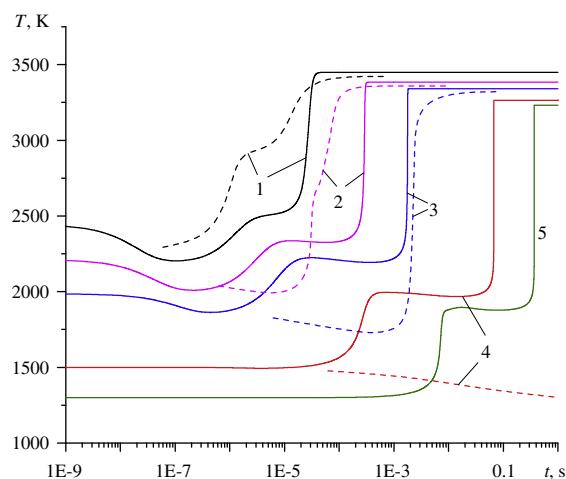


Fig. 5. Evolution of temperature during ignition in the stoichiometric Al(l)-H<sub>2</sub>O mixture calculated in [17] (dotted curves) and computed with the usage of the present work model (solid curves) at T<sub>0</sub> = 2800, 2300, 2000, 1500 and 1300 K (curves 1–5), P<sub>0</sub> = 1 atm.

model [17] does not provide the ignition at all at T<sub>0</sub> = 1500 K, while the model of the present work predicts the ignition with τ<sub>in</sub> ~ 0.1 s at T<sub>0</sub> = 1500 K and with τ<sub>in</sub> ~ 0.4 s at T<sub>0</sub> = 1300 K. In addition, at T<sub>0</sub> ~ 2000 K, the behaviors of τ<sub>in</sub>(T<sub>0</sub>) dependences predicted by two considered models differ principally. The kinetic model [17], after the initial decrease of gas temperature, predicts the continuous growth of T value, while the model of this work gives that after somewhat increase of the temperature, which follows the stage of temperature decrease, the T value falls down slightly again (or does not change at all), and the decrease in the temperature continues up to the ignition event.

When studying numerically the combustion of micrometer-sized aluminum particles with steam, Washburn et al. [11] included detailed chemical reaction mechanism in their model. This mechanism is practically identical to that developed by Huang et al. [17]. The difference between reaction mechanisms [11] and [17] concerns the account for the sensitivity coefficients for some recombination reactions and the reactions of alumina liquid phase formation. Our calculations have shown that these discrepancies do not influence the chain mechanism development and ignition of Al–H<sub>2</sub>O mixture.

It should be emphasized that predictions of the model of the present work are in reasonable agreement with the experimental



data on the registration of selfignition temperature for nano-sized Al particles in the Al-water vapor system. So Parr et al. [34] detected the ignition temperature in the range of 1325–1360 K for 24–192 nm Al particles. Low temperature ignition of micron-sized particles in a water vapor medium was also observed by Schoenitz et al. [35]. Our model provides the ignition of aluminum in steam at  $T = 1300$  K for atmospheric pressure that is close to the ignition temperature detected in the experiments.

In order to explain a distinction in the predictions of the models under study, let's consider the temporal profiles of species mole fractions during ignition calculated with the use of both considered kinetic mechanisms. Figures 6 and 7 depict such profiles for the atmospheric stoichiometric Al(l)–H<sub>2</sub>O mixture at  $T_0 = 2300$  and 1500 K, respectively. It follows from the plots shown in Fig. 6 that, at the initial stage of ignition process, the evolution of reagents predicted by two considered models differs considerably. In order to clarify this fact, consider the chain mechanism development providing by both models. The sequence of reaction steps in Al(l) + H<sub>2</sub>O reacting system provided by these models is depicted in Fig. 8.

At high initial temperature  $T_0 = 2300$  K (Fig. 6), the model [17] predicts identical rates of Al(l) and H<sub>2</sub>O concentrations decrease and the rate of H<sub>2</sub> concentration increase at the initial stage. It means that the process  $\text{Al} + \text{H}_2\text{O} = \text{AlO} + \text{H}_2$  prevails. This reaction is slightly exothermic and results in the temperature increase. As there is no oxidizer for molecular hydrogen in the mixture, its concentration monotonically increases, while the concentration of aluminum oxide AlO diminishes due to recombination processes  $\text{AlO} + \text{Al} = \text{Al}_2\text{O}$  and  $2\text{AlO} = \text{Al}_2\text{O}_2$ . Atom O, which is needed for Al<sub>2</sub>O<sub>3</sub> formation in the reaction  $\text{Al}_2\text{O}_2 + \text{O} = \text{Al}_2\text{O}_3$ , mainly forms due to the dissociation process  $\text{Al}_2\text{O}_2 = \text{Al}_2\text{O} + \text{O}$ .

In the model of the present work, the occurrence of another chain-initiation reaction  $\text{Al} + \text{H}_2\text{O} = \text{AlOH} + \text{H}$  (it is also exothermic) immediately results in the formation of active H atoms, which in the course of two successive processes  $\text{H} + \text{H}_2\text{O} = \text{H}_2 + \text{OH}$  and  $\text{H} + \text{OH} = \text{H}_2 + \text{O}$  give OH radical and O atom. On the one hand, the recombination of AlOH with OH leads to the formation of Al(OH)<sub>2</sub> and Al(OH)<sub>3</sub> as well as AlO<sub>2</sub>H in the reaction  $\text{Al(OH)}_2 = \text{AlO}_2\text{H} + \text{H}$ . These species are long-lived and delay the ignition. On the other hand, O atoms react with AlOH and H<sub>2</sub> generating AlO (and, as consequence, Al<sub>2</sub>O<sub>2</sub> and Al<sub>2</sub>O<sub>3</sub>) and H<sub>2</sub>O. As well, AlO forms in the course of reaction  $\text{AlOH} + \text{Al} = \text{AlO} + \text{AlH}$ . This reaction decreases the rate of AlOH, Al(OH)<sub>2</sub>, Al(OH)<sub>3</sub> and AlO<sub>2</sub>H production. From the plots shown in Figs. 6 and 7 one can see that, after the rapid decrease, the concentration of water vapor begins to increase up to the ignition time instant, and this promotes the ignition. These two contradicted factors result in longer delay time at high

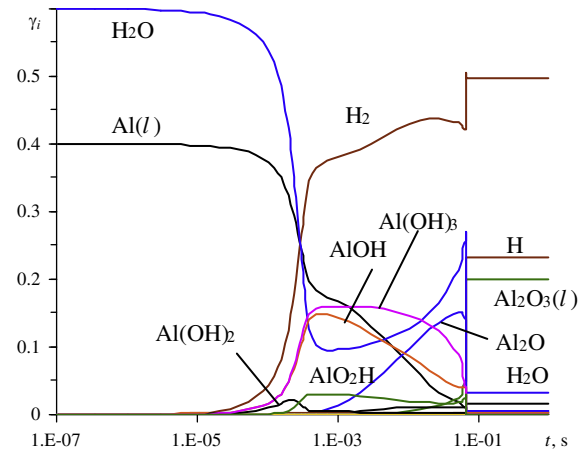


Fig. 7. Variation of species mole fractions during the ignition of stoichiometric Al(l)–H<sub>2</sub>O mixture computed with the usage of the present work model at  $T_0 = 1500$  K,  $P_0 = 1$  atm.

temperatures and, on the contrary, in shorter one at low temperatures in the predictions of our model compared to those of the model [17]. It should be noted that at  $T_0 \leq 1500$  K (see Fig. 5) the model [17] does not predict the ignition at all and, hence, cannot describe the experimental data on the ignition temperature [34].

Computations showed that the smaller the initial temperature is, the higher is the yield of molecular hydrogen in the combustion products. Thereby, in order to produce more hydrogen during combustion of nanoaluminum in water vapor, it is necessary to use small nanoparticles that allow one to decrease the ignition temperature. The difference in the H<sub>2</sub> yield in the cases with  $T_0 = 1300$  and 2300 K achieves 25%.

Let us consider now how ignition temperature and induction time change when aluminum in the gas phase reacts with steam. Remind that such regime of combustion can be realized for small nonoxidized nano-sized particles [22]. Figure 9 shows the temporal profiles of gas temperature for the stoichiometric Al(g)–H<sub>2</sub>O mixture at different initial temperatures, predicted by the model [17] and the model of the present work. One can see that, in this case, the ignition temperature and induction time are smaller than those for the liquid aluminum combustion, and both models under study provides this tendency. These results are qualitatively consistent with the experimental data (see [8]) that demonstrates the decrease in ignition temperature with decreasing particle diameter. As it follows from the plots shown in Fig. 9, both models predict that the stoichiometric Al(g)–H<sub>2</sub>O mixture ignites at

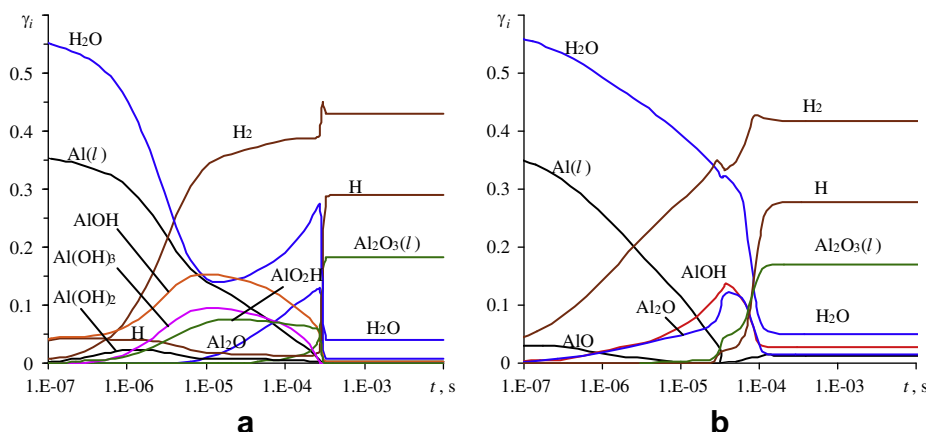
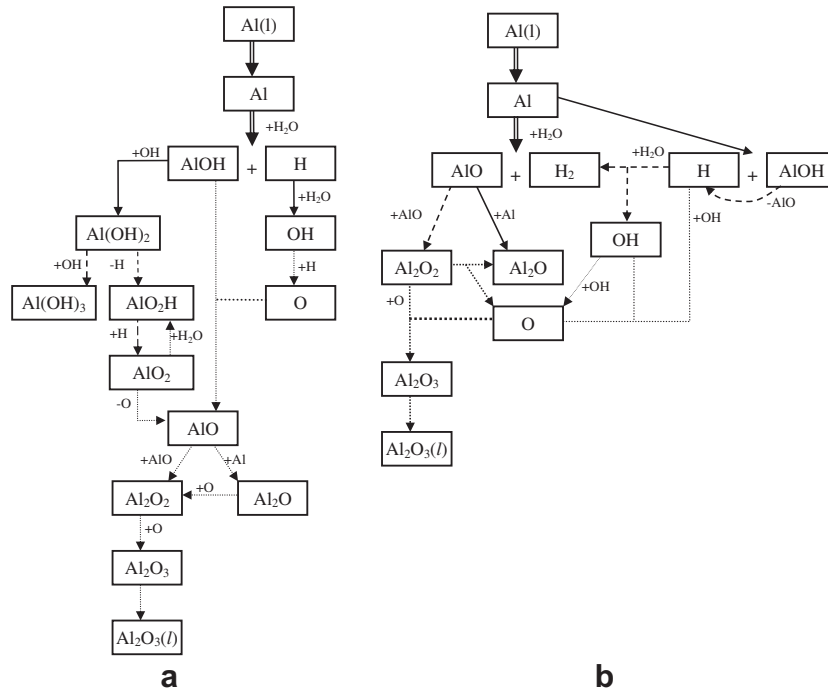
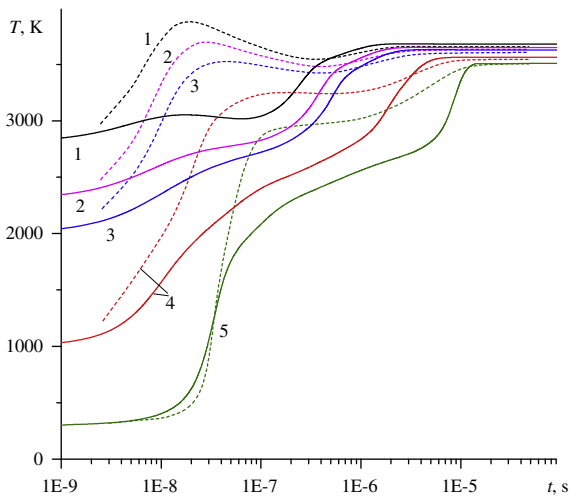


Fig. 6. Variation of species mole fractions during the ignition of the stoichiometric Al(l)–H<sub>2</sub>O mixture computed with the usage of the present work model (a) and calculated in [17] (b) at  $T_0 = 2300$  K,  $P_0 = 1$  atm.



**Fig. 8.** The schemes of chain-mechanism development in the Al(l)-H<sub>2</sub>O reacting system in accordance with our model (a) and the model [17] (b). Dotted, dashed and solid arrows denote the reaction channels with different rates (in ascending order).



**Fig. 9.** Evolution of temperature during ignition in the stoichiometric Al(g)-H<sub>2</sub>O mixture calculated in [17] (dotted curves) and predicted by the model of the present work (solid curves) at  $T_0 = 2800, 2300, 2000, 1000$  and  $300$  K (curves 1–5),  $P_0 = 1$  atm.

$T_0 = 300$  K. However the temperature profiles calculated with the use of model [17] and the model of the present work differ notably. The model [17] predicts the existence of maximum in the temperature behavior at  $T_0 \geq 2000$  K, while the model of the present work does not give such a type of temperature profile.

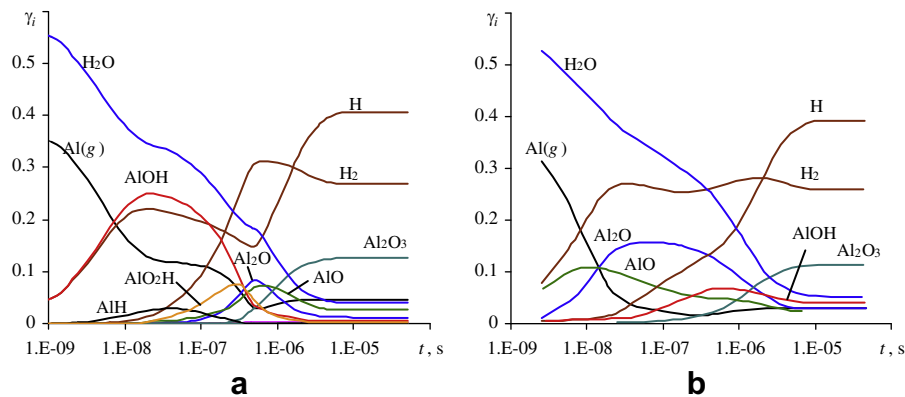
Figure 10 depicts the time history of species mole fractions calculated with the use of both models under study for the stoichiometric Al(g)-H<sub>2</sub>O mixture at  $T_0 = 2300$  K and  $P_0 = 1$  atm. One can see that the model [17] gives the appearance of AlO and H<sub>2</sub> at the first stage due to the reaction of Al with H<sub>2</sub>O. Because this reaction is slightly exothermic, the temperature, at the initial stage, increases. Further, in line with model [17], AlO reacts with aluminum atom  $\text{AlO} + \text{Al} = \text{Al}_2\text{O}$ . As a result, the concentration of Al<sub>2</sub>O rises drastically and becomes greater than that of AlO. This reaction occurs

with high energy release ( $\Delta H \sim 560$  kJ/mol). That is why the model [17] provides a strong increase in the temperature up to the time instant  $t \sim 2 \cdot 10^{-8}$  s. At this time instant the concentration of H<sub>2</sub> achieves its local maximum and then varies only slightly due to competition of the process of H<sub>2</sub> formation in the course of reaction  $\text{Al} + \text{H}_2\text{O} = \text{AlO} + \text{H}_2$  and the process of H<sub>2</sub> elimination owing to its dissociation  $\text{H}_2 + \text{M} = \text{H} + \text{H} + \text{M}$ . Because the dissociation process is endothermic, the temperature, at first stage, decreases and, then, starts to increase due to formation of condensed phase of Al<sub>2</sub>O<sub>3</sub>.

The model of the present work predicts the fast arise of AlOH and atomic hydrogen at initial stage. The absence of AlO does not lead to the fast formation of Al<sub>2</sub>O and, consequently, to the great energy release. The temperature rises monotonically mostly due to heat release in the course of reaction  $\text{AlO}_2 + \text{H} + \text{M} = \text{AlO}_2\text{H} + \text{M}$ , and then, at the final stage, owing to formation of Al<sub>2</sub>O<sub>3</sub> in the liquid phase. Thus, one can conclude that when aluminum both in the liquid and in the gas phases reacts with the water vapor, there exists two-stage regime of ignition. The final temperature in the combustion exhaust is higher when we deal with the combustion of small sized particles which can gasify very rapidly and, as a result, gaseous aluminum reacts with steam.

### 3.4. Flame propagation in the Al-H<sub>2</sub>O mixture

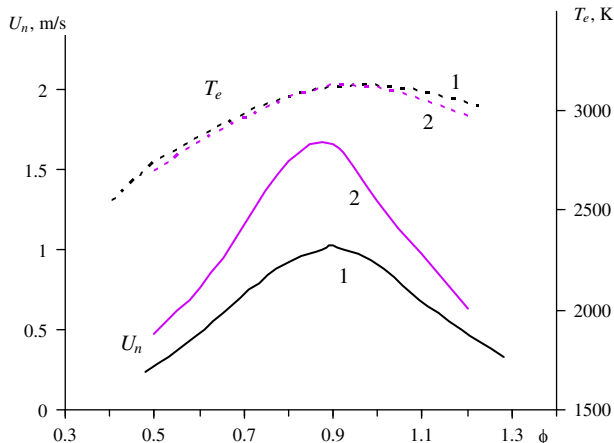
Because the equilibrium temperature achieved during combustion of Al-air system is notably higher than that for the Al-H<sub>2</sub>O mixture, the flame speed in the former case must be greater than that in the latter one. This conclusion was proved by calculations of Huang et al. [17]. Our computations, based on the reaction mechanism of the present work, exhibit the similar results. However, the reaction mechanism of the present work provides higher values of the speed of laminar flame propagation  $U_n$  than the Huang et al. mechanism [17], while the values of final flame temperature are identical to those predicted by Huang et al. [17]. Depicted in Fig. 11 is a comparison of  $U_n$  and  $T_b$  predicted by the mechanisms of [17] and the present work for Al(l)-H<sub>2</sub>O system at various  $\phi$  values. As is seen, the difference in the flame speed calculated with the



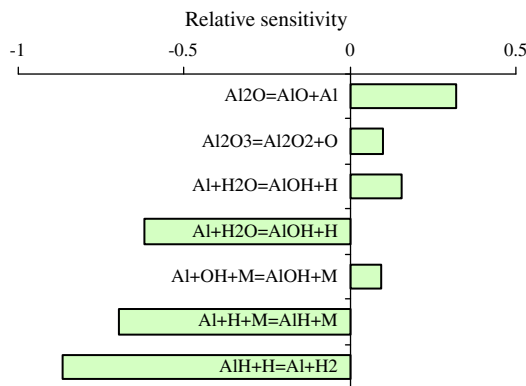
**Fig. 10.** Variation of species mole fractions during the ignition of the stoichiometric Al(g)-H<sub>2</sub>O mixture calculated by using the model of the present work (a) and calculated in [17] (b) at T<sub>0</sub> = 2300 K, P<sub>0</sub> = 1 atm.

use of reaction mechanisms under consideration can achieve a factor of 1.6. The maximal value of U<sub>n</sub> predicted by the model of this work is achieved at φ = 0.9 (T<sub>0</sub> = 800 K, P<sub>0</sub> = 1 atm) and equals to 1.66 m/s. The sensitivity analysis showed that U<sub>n</sub> value is the most sensitive to the rate constants of the reactions AlH + H = Al + H, Al + H + M = AlH + M and Al + H<sub>2</sub>O = AlOH + H (see Fig. 12). It should be emphasized that precisely rate coefficients of these reactions were reconsidered compared to the model [17]. Moreover, the Al + H<sub>2</sub>O reaction channel, included in the Huang et al. reaction mechanism [17], differs principally from that involved in the model of the present work. This leads to substantial distinction in the variation of species concentrations of mixture components and, especially, the concentration of H atoms in the flame front predicted by the model of Huang et al. [17] and the model of the present work. Because the laminar flame speed for the Al-H<sub>2</sub>O reactive system is determined mainly by the rate of the diffusion of H atoms, the U<sub>n</sub> values predicted by these two models must be different. At the same time, the value of final temperature in the combustion products is determined by the thermodynamic properties of individual species taken in the kinetic model. These properties do not differ significantly for both models under study. That is why the T<sub>e</sub> values predicted by the models of [17] and the present work for Al(l)-H<sub>2</sub>O system are practically identical despite the substantial difference in the reaction mechanisms.

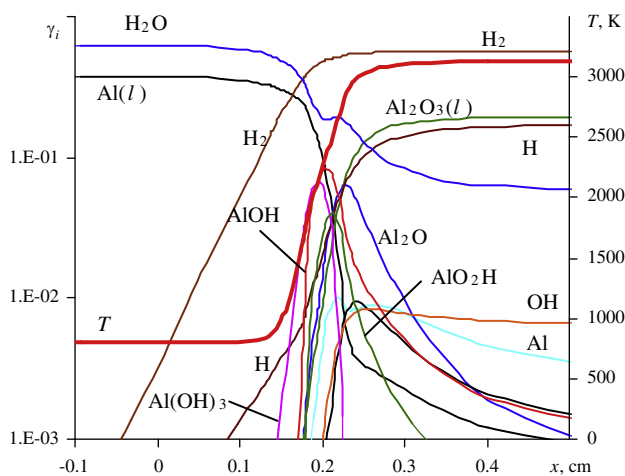
The variation of species mole fractions and temperature in the flame front region for the Al(l)-H<sub>2</sub>O system with φ = 0.9 predicted



**Fig. 11.** Flame speed U<sub>n</sub> and final flame temperature T<sub>e</sub> values (solid and dotted lines) vs. fuel/oxidizer equivalence ratio φ of the Al(l)-H<sub>2</sub>O mixture calculated in [17] and predicted by the model of the present work (curves 1 and 2, respectively) at T<sub>0</sub> = 800 K and P<sub>0</sub> = 1 atm.



**Fig. 12.** Sensitivity of laminar flame speed to the reaction rates for the stoichiometric Al(l)-H<sub>2</sub>O mixture at T<sub>0</sub> = 800 K, P<sub>0</sub> = 1 atm calculated with the usage of the model of the present work.



**Fig. 13.** Variation of species mole fractions and temperature in the Al(l)-H<sub>2</sub>O laminar flame predicted by the model of the present work at T<sub>0</sub> = 800 K, P<sub>0</sub> = 1 atm, φ = 0.9.

by the model of the present work is depicted in Fig. 13. One can see that the main combustion products are H<sub>2</sub> and liquid Al<sub>2</sub>O<sub>3</sub>, and the final temperature does not exceed 3250 K. The mole fraction of molecular hydrogen γ<sub>H<sub>2</sub></sub> in the combustion products, at such conditions, is as large as 0.55. The content of Al<sub>2</sub>O<sub>3</sub> in the liquid phase is much smaller (γ<sub>Al<sub>2</sub>O<sub>3</sub></sub> = 0.2). The third component with



notable mole fraction is atomic hydrogen ( $\gamma_{\text{H}} = 0.15$ ). The concentrations of other mixture components are much smaller.

Unfortunately, until now there have been no experimental data on the flame speed in the Al–H<sub>2</sub>O system with dust particles of different sizes. Only the  $U_n$  value for the system of aluminum particles being in the liquid water were measured [13]. Nevertheless, it is interesting to compare the measurements [13] with the predictions of the model of this work. Certainly, the comparison can be considered as qualitative only. In the experiments [13], the oxidized Al nanoparticles with diameter of 38 nm were used, and the concentration of Al<sub>2</sub>O<sub>3</sub> in the whole particle volume was equal to 54.3 wt%. The system was primarily ignited and the change in the position of flame front was detected.

We assumed that burning of Al nanoparticles occurs in steam environment after vaporization of liquid water and took into consideration that Al and Al<sub>2</sub>O<sub>3</sub> in a particle can be both in solid and in liquid phases. Therefore, we considered the combustion of the mixture with following composition Al(s,l)/Al<sub>2</sub>O<sub>3</sub>(s,l)/H<sub>2</sub>O(g) = 1/0.2228/1.5. In addition, we supposed that particles burn in water vapor, if they are in a liquid phase. So far as, in line with the theory reported in [36], the melting temperature for Al particles with  $d = 38$  nm is approximately equal to 840 K [8]. The thermodynamic properties for combined solid-liquid phase Al(s,l) and Al<sub>2</sub>O<sub>3</sub>(s,l) were taken from Janaf tables [37]. In this case, the reactions Al(l) = Al(g) and Al<sub>2</sub>O<sub>3</sub>(l) = Al<sub>2</sub>O<sub>3</sub>(g) were replaced as Al(s,l) = Al(g) and Al<sub>2</sub>O<sub>3</sub>(s,l) = Al<sub>2</sub>O<sub>3</sub>(g). In experiments [13] at  $P_0 = 100$  atm the flame speed was detected in the range of 5.5–8.5 m/s. Our model predicts  $U_n = 8.9$  m/s, that is in reasonable coincidence with experimental data, though, in this case, we can say only about qualitative agreement, because the experimental conditions cannot be modeled exactly.

#### 4. Conclusions

The extended reaction mechanism for the Al–H<sub>2</sub>O system combustion treated the recent *ab initio* and theoretical studies of the elementary reactions in such system was built. In contrast, previous kinetic models of Huang et al. [17] and Washburn et al. [11], this mechanism suggests the novel channel of chain initiation Al + H<sub>2</sub>O = AlOH + H and involves the reactions with additional species Al(OH)<sub>2</sub>, Al(OH)<sub>3</sub> and AlO<sub>2</sub>H. The developed reaction mechanism allows one to describe with reasonable accuracy the experimental data on ignition temperature obtained for combustion of nanoaluminum in oxygen and in water vapor.

The computations showed that in the Al–H<sub>2</sub>O mixture when aluminum is both in the liquid and in the gas phase, there exists two-stage regime of ignition and this fact is determined by the features of chain mechanism development. If the aluminum comes into the steam environment in a liquid phase, the ignition occurs at higher temperature ( $T_{\text{ign}} = 1300$  K for stoichiometric Al–H<sub>2</sub>O mixture) than that when Al reacts with water vapor in the gas phase. The decrease in the initial temperature leads to the increase in the yield of molecular hydrogen in combustion exhaust. In the laminar atmospheric flame in Al(l)–H<sub>2</sub>O mixture close to stoichiometric one, the mole fraction of H<sub>2</sub> can be as large as 0.55, while the mole fractions of the other major components, liquid Al<sub>2</sub>O<sub>3</sub> and atomic hydrogen, are equal to 0.2 and 0.15 respectively. Because the aluminum in the gas phase reacts with steam much faster even at low temperatures, it is reasonable in the process of hydrogen production to use small nano-sized non-oxidized Al particles that, due to their rapid gasification, can provide the burning of aluminum in the gas phase.

#### Acknowledgments

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#### References

- [1] E. Shafirovich, V. Diakov, A. Varma, *Combust. Flame* 144 (1–2) (2006) 415–418.
- [2] M.S. Vlaskin, E.I. Shkolnikov, A.V. Lisicyn, A.V. Bersh, A.Z. Zhuk, *Int. J. Hydrogen Energy* 35 (5) (2010) 1888–1894.
- [3] F. Franzoni, M. Milani, L. Montorsi, V. Golovitchev, *Int. J. Hydrogen Energy* 35 (4) (2010) 1548–1559.
- [4] V. Yang, T.B. Brill, W.Z. Ren (Eds.), *Solid propellant chemistry, combustion, and motor interior ballistics*, in: AIAA Progress in Aeronautics and Astronautics, 2000, p. 663.
- [5] A. Dokhan, E.W. Price, J.M. Seitzman, R.K. Sigman, *Proc. Combust. Inst.* 29 (2) (2002) 2939–2945.
- [6] A. Ingenito, C. Bruno, *J. Propul. Power* 20 (6) (2004) 1056–1064.
- [7] J. Servaites, H. Krier, J.C. Melcher, R.L. Burton, *Combust. Flame* 125 (1–2) (2001) 1040–1054.
- [8] Y. Huang, G.A. Risha, V. Yang, R.A. Yetter, *Combust. Flame* 156 (1) (2009) 5–13.
- [9] M.T. Swihart, L. Catoire, *J. Phys. Chem. A* 105 (1) (2001) 264–273.
- [10] M.W. Beckstead, Y. Liang, K.V. Pudduppakkam, *Combust. Explos. Shock Waves* 41 (6) (2005) 622–638.
- [11] E.B. Washburn, J.N. Trivedi, L. Catoire, M.W. Beckstead, *Combust. Sci. Technol.* 180 (8) (2008) 1502–1517.
- [12] R.A. Yetter, F.L. Dryer, *Metal Particle Combustion and Classification, Micro-Gravity Combustion: Fire in Free Fall*, Academic Press, 2001, p. 419.
- [13] G.A. Risha, J.L. Sabourin, V. Yang, R.A. Yetter, S.F. Son, B.C. Tappan, *Combust. Sci. Technol.* 180 (12) (2008) 2127–2142.
- [14] P. Escot Bocanegra, V. Sarou-Kanian, D. Davidenko, C. Chauveau, I. Gökalp, *Prog. Propul. Phys.* 1 (2009) 47–62.
- [15] N. Moallemi, M. Bidbadi, M. Jadidi, M.E. Hosseini, *Aust. J. Basic Appl. Sci.* 5 (12) (2011) 1991–1998.
- [16] P. Lynch, G. Fiore, H. Krier, N. Glumac, *Combust. Sci. Technol.* 182 (7) (2010) 842–857.
- [17] Y. Huang, G.A. Risha, V. Yang, R.A. Yetter, *Analysis of nano-aluminum particle dust cloud combustion in different oxidizer environments*, in: 43rd AIAA Aerospace Sciences Meeting and Exhibit 10–13 January 2005, Reno, Nevada, AIAA 2005-738.
- [18] V.I. Levitas, *Combust. Flame* 156 (2) (2009) 543–546.
- [19] S. Álvarez-Barcia, J.R. Flores, *Chem. Phys.* 382 (1–3) (2011) 92–97.
- [20] A. Sharipov, N. Titova, A. Starik, *J. Phys. Chem. A* 115 (17) (2011) 4476–4481.
- [21] A.S. Sharipov, N.S. Titova, A.M. Starik, *Combust. Theory Model.* 16 (5) (2012) 842–868.
- [22] T.J. Campbell, G. Aral, S. Ogata, R.K. Kalia, A. Nakano, P. Vashishta, *Phys. Rev. B* 71 (20) (2005) 205413.
- [23] R.E. McClean, H.H. Nelson, M.L. Campbell, *J. Phys. Chem.* 97 (38) (1993) 9673–9676.
- [24] S. Álvarez-Barcia, J.R. Flores, *Chem. Phys.* 374 (1–3) (2010) 131–137.
- [25] J. Guo, J.M. Goodings, A.N. Hayhurst, *Combust. Flame* 150 (1–2) (2007) 127–136.
- [26] A.A. Granovsky, *Firefly*. <<http://classic.chem.msu.su/gran/firefly/index.html>>; version 7.1.G.
- [27] A.M. Starik, N.S. Titova, A.S. Sharipov, V.E. Kozlov, *Combust. Explos. Shock Waves* 46 (5) (2010) 491–506.
- [28] A.M. Starik, N.S. Titova, A.S. Sharipov, *Kinetic mechanism of H<sub>2</sub>–O<sub>2</sub> ignition promoted by singlet oxygen O<sub>2</sub>(<sup>a</sup>Δ<sub>g</sub>)*, in: G.D. Roy, S.M. Frolov (Eds.), *Deflagrative and Detonative Combustion*, Torus Press, Moscow, 2010, pp. 19–42.
- [29] M.T. Swihart, L. Catoire, B. Legrand, I. Gökalp, C. Paillard, *Combust. Flame* 132 (1–2) (2003) 91–101.
- [30] E. Goos, A. Burcat, B. Ruscic, *Ideal gas thermochemical database with updates from active thermochemical tables*. <<ftp://ftp.technion.ac.il/pub/supported/aetdd/thermodynamics>>.
- [31] J.O. Hirschfelder, C.F. Curtiss, R.B. Bird, *Molecular Theory of Gases and Liquids*, Wiley, N.Y., 1954.
- [32] R.D. Johnson III, NIST Computational Chemistry Comparison and Benchmark Database, NIST Standard Reference Database, Number 101, Release 15b, August 2011. <<http://cccbdb.nist.gov>>.
- [33] R.J. Kee, F.M. Rupley, J.A. Miller, M.E. Coltrin, J.F. Grcar, E. Meeks, et al. *CHEMKIN Release 4.0*, Reaction Design Inc., San Diego, CA, 2004.
- [34] T.P. Parr, C. Johnson, D. Hanson-Parr, K. Higa, K. Wilson, *Evaluation of advanced fuels for underwater propulsion*, in: 39th JANNAF Combustion Subcommittee Meeting, December 2003.
- [35] M. Schoenitz, C.-M. Chen, E.L. Dreizin, *J. Phys. Chem. B* 113 (15) (2009) 5136–5140.
- [36] C.R.M. Wronski, *Br. J. Appl. Phys.* 18 (12) (1967) 1731–1737.
- [37] M.W. Chase Jr., C.A. Davies, J.R. Downey, D.J. Frurip Jr., R.A. McDonald, A.N. Syverud, *NIST JANAF THERMOCHEMICAL TABLES*, 1985, Version 1.0. <<http://kinetics.nist.gov/janaf/>>