# Schottky-Barrier S/D MOSFETs With High-K Gate Dielectrics and Metal-Gate Electrode

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Abstract—This letter presents a low-temperature process to fabricate Schottky-barrier silicide source/drain transistors (SSDTs) with high- $\kappa$  gate dielectric and metal gate. For p-channel SSDTs (P-SSDT) using PtSi source/drain (S/D), excellent electrical performance of  $I_{\rm on}/I_{\rm off} \sim 10^7 - 10^8$  and subthreshold slope of 66 mV/dec have been achieved. For n-channel SSDTs (N-SSDTs) using DySi<sub>2-x</sub> S/D,  $I_{\rm on}/I_{\rm off}$  can reach  $\sim 10^5$  at  $V_{\rm ds}$  of 0.2 V with two subthreshold slopes of 80 and 340 mV/dec. The low-temperature process relaxes the thermal budget of high- $\kappa$  dielectric and metal-gate materials to be used in the future generation CMOS technology.

*Index Terms*—High- $\kappa$ , metal gate, MOSFET, Schottky.

### I. INTRODUCTION

T HE SERIES resistance of ultrashallow source/drain (S/D) junctions is a serious issue for future CMOS transistor scaling. Schottky-barrier silicide S/D (SSDT) structure has been suggested as a potential solution, due to the sharp silicide/silicon interface and low sheet resistance of silicide [1]–[5]. SSDT is particularly attractive when a metal-gate/high- $\kappa$  gate stack is employed, as it avoids the use of a high-temperature annealing process required for implanted S/D junctions and poly gate, hence, eliminating the thermal stability issues associated with high- $\kappa$  gate stack [6]. In this letter, we successfully demonstrate bulk SSDTs with HfO<sub>2</sub> high- $\kappa$  dielectric, HfN metal gate, and PtSi (for pMOS) and DySi<sub>2-x</sub> (for nMOS) silicide S/D using a simplified low temperature process. The process can

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be easily extended to ultrathin body (UTB) silicon-on-insulator (SOI) structure to further improve the SSDT performance.

# II. MOS DEVICE FABRICATION

Si (100) wafers of both n- and p-type with resistivity of 4–8  $\Omega$  · cm were used as the starting substrate. First, ~6-nm HfO<sub>2</sub> film was deposited on Si at 400 °C using Hf[OC(CH<sub>3</sub>)<sub>3</sub>]<sub>4</sub> and O<sub>2</sub> in a metal-organic chemical vapor deposition system, followed by an in situ annealing in N2 ambient at 700 °C. Then HfN( $\sim$ 50 nm) and TaN( $\sim$ 100 nm) were deposited sequentially in a sputtering system with base pressure of  $1.5 \times 10^{-7}$  torr. TaN is used as a capping layer to reduce the gate sheet resistance  $(\sim 10 \Omega \text{ sq})$  [7]. The wafers were patterned using photolithography and reactive ion etching procedures. Immediately after dipping in a diluted hydrogen fluoride solution (DHF), the wafers were loaded into the sputtering system again for platinum (~100 nm) (for P-SSDT) or dysprosium (~100 nm)/HfN  $(\sim 70 \text{ nm})$  stack (for N-SSDT) deposition. Since Dy is easily oxidized during ex situ anneal, a capping layer of thermally stable HfN [7] is used to prevent Dy oxidization. Silicidation of Pt or Dy was performed by forming gas anneal at 420 °C for 1 h. Then unreacted Pt was removed by hot diluted Aqua Regia etch. The HfN capping layer and unreacted Dy were etched by DHF and a diluted HNO<sub>3</sub> solution sequentially.

Although PtSi has flat surface, the surface of  $DySi_{2-x}$  film is quite rough. This is because  $DySi_{2-x}$ , as well as other rare earth (RE) silicides, is not formed through layer-by-laye, but islanded-preferred during the solid-state reaction with substrate Si [8]. No improvement to the  $DySi_{2-x}$  film morphology was found by using other capping layers such as Pt, Ti, Ta, Ru, and Al, etc. or by *in situ* vacuum anneal without a capping layer. HfN is found to be the most suitable capping layer in our study, due to its easy removal and no contamination in  $DySi_{2-x}$ .

## III. DEVICE CHARACTERIZATION AND DISCUSSION

For the PtSi/n-Si Schottky contact, the electron barrier height  $(\Phi_{\rm bN})$  was measured to be 0.85 and 0.86 eV from current– voltage (*I–V*) and capacitance–voltage (*C–V*), respectively. The corresponding hole barrier height ( $\Phi_{\rm bP} \cong E_g - -\Phi_{\rm bN}$ ) is about 0.26 eV, close to the reported value [9]. For the DySi<sub>2-x</sub>/p–Si diode, Fig. 1 shows that  $\Phi_{\rm bP}$  are 0.66 eV (*I–V*) and 0.88 eV (*C–V*), compared to the reported value of ~0.74 eV for the DySi<sub>2-x</sub> film formed by ultrahigh vacuum evaporation [9]. The large difference in  $\Phi_{\rm bP}$  from *C–V* and *I–V*, and the larger than



Fig. 1. I-V curves of three different  $\text{DySi}_{2-x}/\text{p-Si}$  diodes. (1) Normal  $\text{DySi}_{2-x}$  formed by Dy/Si silicidation at 420 °C for 1 h. (2) Adding an ultrathin intermediate Ge layer (~1 nm). (3) A very thin  $\text{DySi}_{2-x}$  with an Al capping layer. The barrier heights  $\Phi_{\text{bP}}$  and ideality factors deduced from the thermal emission model are (0.66 eV, 1.10), (0.72 eV, 1.05), and (0.82 eV, 1.46), respectively. The inset is the C-V curves of diode 1. The deduced  $\Phi_{\text{bP}}$  and N<sub>A</sub> are 0.88 eV and 2.1 × 10<sup>15</sup> cm<sup>-3</sup>, respectively. Other two diodes have the similar C-V results within the experimental error.



Fig. 2.  $I_d$ - $V_{ds}$  and  $I_d$ - $V_g$  curves of P-SSDT with PtSi S/D. The channel width and length are 400 and 4  $\mu$ m. Equivalent oxide thickness (EOT) of HfO<sub>2</sub> gate dielectric is 2.0 nm,  $V_t = -0.56$  V. The inset is the XTEM image of the device structure. A "hole" between S/D and channel acts as a sidewall spacer.

unity ideality factor (*n*) imply significant barrier height inhomogeneity of the  $\text{DySi}_{2-x}/\text{Si}$  interface [10], associated with the rough surface and interface of  $\text{DySi}_{2-x}/\text{Si}$  as observed by atomic force microscopes and cross-sectional transmission electron microscope (XTEM) measurements. Fig. 2 and the inset show the *I*-*V* characteristics of P-SSDT and the device structure. The drain current at small  $V_g(I_{\text{off}})$  is between 1–10 pA



Fig. 3.  $I_d-V_g$  and  $I_d-V_g$  curves of N-SSDT with  $\text{DySi}_{2-x}$  S/D. The channel width and length are 400 and 4  $\mu$  m. EOT is 1.5 nm, S = 80 mV/dec at  $V_g < 0.23$  V, and 340 mV/dec at  $V_g > 0.23$  V,  $V_t = 0.56$  V. The inset is the schematic band diagram of Schottky barrier at source/channel. If the device is already turned on before  $V_g$  reaches  $V_{\rm sbfb}$ , it shows one subthreshold slope as conventional device. Otherwise, it shows two subthreshold slopes because the slope reduces dramatically when  $V_g > V_{\rm sbfb}$ .

for devices with drain area of  $1 \times 10^{-4}$  cm<sup>2</sup>, close to the theoretical value of the reverse current of the PtSi/n-Si diode with barrier height of 0.86 eV ( $\sim$ 3.8 pA).  $I_{on}$  depends on process parameters more sensitively. In our simplified one-mask process, a "hole" between the S/D and gate is formed during DHF dipping because HfN can be etched by DHF while TaN cannot. The "hole" acts as a spacer to separate the S/D and the gate. The "hole" size can be controlled by DHF dipping time. For the  $W/L = 200/2.5 - \mu m$  devices, the  $I_{\rm on}$  at  $V_g = V_{\rm ds} = -2.5$  V increases from 3.5 to 6.2 mA by reducing the DHF dipping time from 105 to 60 s. Our P-SSDTs have a  $I_{\rm on}/I_{\rm off}$  ratio  $\sim 10^8$  at  $V_q = V_{ds} = -1.0$  V, with the subthreshold slope of 66 mV/dec. Their performance is similar to or even slightly better than the best reported data for SiO<sub>2</sub>/poly-Si gate P-SSDT [11]. N-SSDT has  $\sim$ three orders of magnitude larger  $I_{\rm off}$  than P-SSDT. Fig. 3 shows that the  $I_{\rm on}/I_{\rm off}$  ratio at a small  $V_{\rm ds}$  (0.2 V) is about 10<sup>5</sup>, close to the recently reported value of N-SSDT on SOI [12].  $I_{\rm off}$  increases significantly at large  $V_{\rm ds}$  due to the quick increase of the  $DySi_{2-x}/p$ -Si diode leakage current with the increase of the reverse bias (see Fig. 1). The transfer  $I_d$ - $V_q$  curves of N-SSDT show two different subthreshold slopes  $\sim$ 80 mV/dec at  $V_q < 0.23$  V ( $V_{\rm sbfb}$ : source-body flatband voltage) and  $\sim 340$ mV/dec at  $V_q > 0.23$  V. This behavior, as well as the large  $I_{\rm off}$  observed in N-SSDT, can be explained by the relatively low hole barrier height  $\Phi_{\rm bP}$  of the DySi<sub>2-x</sub>/p–Si contact as compared to the electron barrier height  $\Phi_{bN}$  of the PtSi/n-Si contact in P-SSDT. As shown in the inset of Fig. 3, the tunneling



Fig. 4. Transfer curves of N-SSDTs with  $\text{DySi}_{2-x}$  S/D formed at different conditions, (also see Fig. 1). The device with very thin  $\text{DySi}_{2-x}$  has lower  $I_{\text{off}}$  and  $I_{\text{on}}$ . The  $I_{\text{on}}/I_{\text{off}}$  ratio at  $V_{\text{ds}} = 50$  mV is improved about 4 to 5 times by adding a very thin (~1 nm) Ge intermediate layer.

current  $(J_{\rm TN})$  in N-SSDT through the electron Schottky barrier  $\Phi_{\rm bN} \cong E_g - \Phi_{\rm bP} = 0.38$  eV between source and channel dominates in the  $V_g > V_{\rm sbfb}$  region, resulting in two slopes below threshold [5], [13]. It exhibits only one subthreshold slope if  $V_{\rm sbfb}$  is larger than the threshold voltage  $(V_{\rm th})$ , which requires  $\Phi_{\rm bN} < E_g/2 - (kT/q) \ln(N_A/n_i)$  (~0.26 eV in our case) and this condition is not satisfied for DySi<sub>2-x</sub>/Si Schottky junction. On the other hand, PtSi–Si contact meets the requirement of  $\Phi_{\rm bP} < E_g/2 - (kT/q) \ln(N_D/n_i)$  (~0.27 eV) for P-SSDT. Therefore, only one subthreshold slope of 66 mV/dec is observed in P-SSDT.

Low barrier height is a key issue for SSDT. RE silicides have the lowest electron barrier height among the known silicides [9]. However, their barrier height is not low enough to meet the N-SSDT requirement. Even worse, the large barrier height inhomogeneity of the RE silicide/Si contact causes  $I_{\rm off}$  (~ 25  $\mu$ A/cm<sup>2</sup>) larger than the theoretical value of a uniform barrier height (~  $1.1 \,\mu$ A/cm<sup>2</sup> for DySi<sub>2-x</sub>/p–Si diode with  $\Phi_{bP} = 0.74 \text{ eV}$ ). Improving the RE silicide quality and reducing its barrier height are big challenges for N-SSDT. We found that adding a very thin ( $\sim 1$  nm) deposited amorphous Ge intermediate layer can improve the  $DySi_{2-x}$  film morphology significantly, because the amorphous Ge may suppress the crystallization of  $DySi_{2-x}$  during its growth. The better diode performance is shown in Fig. 1 with  $\Phi_{\rm bP}$  of 0.72 eV, *n* of 1.05, and about half order of magnitude lower reverse current than the normal ones. In Fig. 4, the transistor performance is also improved. The  $I_{\rm on}/I_{\rm off}$  ratio at  $V_{\rm ds} = 0.05$  V increases about 4 to 5 times. Very thin  $\text{DySi}_{2-x}$  film has quite smooth surface, which can be formed by low temperature (300 °C, 1 h) or short time (400 °C, 1 min) anneal. In this case, an additional metal (such as Al) film on the DySi<sub>2-x</sub> is necessary to reduce the series resistance. Fig. 1 shows that  $\Phi_{\rm bP}$  and n are 0.82 eVand 1.46, respectively. Fig. 4 shows that  $I_{\rm off}$  at small  $V_g$  and  $V_{\rm ds}$  is reduced by 8 to 9 times. The value of ~ 2.0  $\mu$ A/cm<sup>2</sup> is close to the theoretical prediction. However,  $I_{\rm on}$  is also reduced due to the large S/D series resistance. Further improvement of the electrical performance of N-SSDT needs  $\Phi_{\rm bN}$  to be reduced to less than 0.2 eV with quite smooth interface. The use of UTB-SOI also helps due to the reduction of the Schottky contact area.

In conclusion, N- and P-SSDTs with high- $\kappa$  dielectric and metal gate were demonstrated for the first time using a low temperature process. The transistors have comparable electrical performance to the best data of SSDT reported so far.

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