

# Positive Oxide Charge-Enhanced Read Disturb in a Localized Trapping Storage Flash Memory Cell

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**Abstract**—Read disturb-induced erase-state threshold voltage instability in a localized trapping storage Flash memory cell with a poly-silicon–oxide–nitride–oxide–silicon (SONOS) structure is investigated and reported. Our results show that positive trapped charge in bottom oxide generated during program/erase (P/E) cycles play a major role. Both gate voltage and drain voltage will accelerate the threshold voltage ( $V_t$ ) drift. Hot-carrier caused disturb effect is more severe in a shorter gate length device at low temperature. A model of positive charge-assisted electron tunneling into a trapping nitride is proposed. Influence of channel doping on the  $V_t$  drift is studied. As the cell is in an “unbiased” storage mode, tunnel detrapping of positive oxide charges is responsible for the threshold voltage shift, which is insensitive to temperature.

**Index Terms**—Cycling-induced oxide charges, Flash EEPROM, hot-carrier effect, MXVAND, NROM, PHINES, positive charge-assisted electron tunneling, read disturb, poly-silicon–oxide–nitride–oxide–silicon (SONOS), tunnel detrapping, threshold voltage ( $V_t$ ) instability.

## I. INTRODUCTION

RECENTLY, localized charge-trapping storage, two-bits-per-cell Flash memory technology with a poly-silicon–oxide–nitride–oxide–silicon (SONOS) structure has attracted much attention [1]–[6]. In addition to its simpler fabrication process, the smaller cell size and higher packing density, the absence of drain turn-on and overerase are considered to be the advantages over current floating-gate Flash memory technologies [7]. Fig. 1 shows a typical SONOS memory cell, which is an nMOSFET with an oxide–nitride–oxide (ONO) stack as the gate dielectric. Based on this device structure, various methods of operation have been proposed. For example, NROM [2], multiplex virtual-grained AND (MXVAND) [3] (hot-electron programming and hot-hole erasing) and programming by hot hole injection nitride electron storage (PHINES) [5] [hot-hole programming and Fowler–Nordheim (FN) tunnel erasing] have been demonstrated to be promising candidates for Flash EEPROM technology. Charge loss characteristics of these cells have been discussed, and corresponding models have been proposed [3], [8]–[11]. Various ONO processing technologies [12] and alternative trapping layer materials [13] have been investigated to further improve their charge retention.

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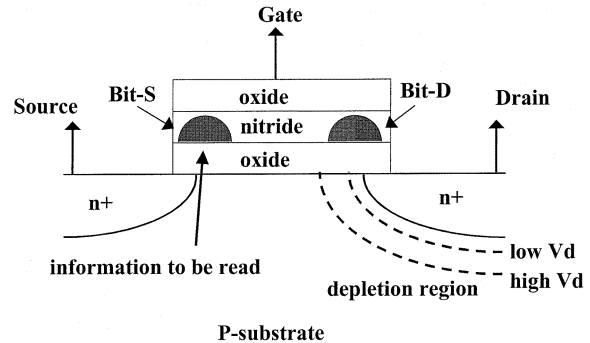


Fig. 1. Schematic representation of the cell structure and localized charge storage. Concept of reverse read is depicted by the depletion region in which the effect of injected charges is screened out.

The unique feature of two-bits-per-cell storage is due to the localized charge injection, nonconducting property of charge storage material, and a reverse read scheme [2]. As depicted in Fig. 1, a drain bias ( $V_d$ ) is necessary to read out the information stored in bit-S. The applied  $V_d$  must be large enough to “screen” out the injected charges at the drain bit (Bit-D). The read current (or threshold voltage) is then controlled by the charge state near the source (Bit-S), almost regardless of the charges at Bit-D. Typically, the applied bit-line voltage at read is around 1.6 V, which is higher than the read voltage in a floating-gate Flash memory cell (typically around 1 V). Fig. 2 shows the threshold voltage ( $V_t$ ) shift of a 10-K program/erase (P/E) cycled cell during read operation, whereas the  $V_t$  drift is not observed in a noncycled cell. In the figure, the applied gate bias ( $V_g$ ) is 2.5 V, and the drain bias is 1.6 V. This  $V_t$  shift, which will degrade the memory window and result in failure of cell operation, however, is rarely explored [14]. The purpose of this work is to find the dominant mechanisms that cause such erase-state  $V_t$  instability.

## II. EXPERIMENTAL

The samples used in this paper are made of an nMOSFET-like device, with an ONO gate dielectric structure, in a virtual ground array. The source and drain are formed by buried-diffusion bit-lines [2]. The thickness of each ONO layer is 9 (top oxide), 6, and 7 nm. Devices with gate length from 0.5–0.3  $\mu\text{m}$  and initial threshold voltage ( $V_{ti}$ ) from 1.2–3 V are characterized. The gate width is around 0.35  $\mu\text{m}$ . Temperature effect from  $-20^\circ\text{C}$  to  $85^\circ\text{C}$  is studied. Channel hot-electron injection and band-to-band tunneling-induced hot-hole injection are utilized for programming and erasing, respectively. Read operation is

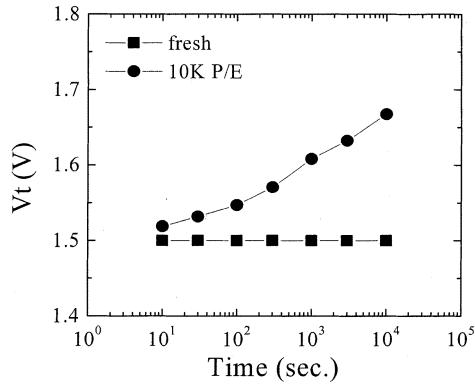


Fig. 2. Read-disturb characteristics of a fresh cell and a 10-K P/E cycled cell.  $L_g = 0.5 \mu\text{m}$  and  $V_g/V_d/V_s = 2.5\text{V}, /1.6\text{V}, /0\text{V}$ .

TABLE I  
OPERATION PRINCIPLES AND BIAS CONDITIONS UTILIZED  
DURING CELL OPERATIONS

	Bit-D			Bit-S		
	Vg	Vd	Vs	Vg	Vd	Vs
Program (CHE)	11V	5V	0V	11V	0V	5V
Erase (BTBT HH)	-3V	8V	0V	-3V	0V	8V
Read (Reverse)	2.5V	0V	1.5V	2.5V	1.5V	0V

achieved with a reverse read scheme. The bias conditions are summarized in Table I. Program time is around  $5 \mu\text{s}$ , and erase time is about 5 ms. With such conditions, 2-V difference between program-state and erase-state  $V_t$  can be obtained. Three disturb modes are investigated for P/E cycled cells in erase state. In the first one, gate and drain biases are applied and the source is grounded. It is the conventional read disturb (RD) mode. The second one has an applied gate bias, and the source and drain are grounded. It is referred to as gate disturb (GD) mode, which is to emulate the gate bias effect on cells sharing the same word-line with the cell to be read. Finally, the room temperature drift (RT) mode, in which the  $V_t$  drift is insensitive to temperature [10], is the special case in which the cell is in “unbiased” storage (without any applied bias).

### III. MODELING OF POSITIVE OXIDE CHARGE EFFECT ON READ DISTURB

Fig. 3 shows the  $V_t$  shift of a 10-K P/E cycled cell in a two-phase measurement, e.g., RT and RD are measured in sequence or vice versa. It is found that after RT (or RD), the following RD (or RT) is reduced. This gives us a hint that RT and RD may have the same cause, since they influence each other. It also means that  $V_t$  shift due to RD is dependent on the sampled time. Since positive oxide charge detrapping (with the coexistence of residual negative charges in nitride) has been clarified to be the major cause of RT [3], [10], these positive oxide charges created during P/E cycling are considered to affect RD also. It also explains the cycling effect on the observed erased-state  $V_t$  instability [3]. Table II shows the  $V_t$  drift models at various bias conditions and the corresponding time evolutions. In each case, the cycling-induced positive oxide charge plays a major role.

As the cell is stored without applied and  $V_d$  (RT mode), tunnel detrapping of the cycling-generated positive oxide charges ( $J_h$

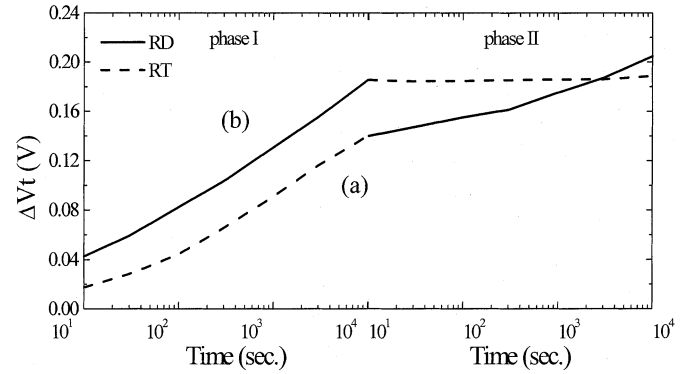


Fig. 3. Two-phase measurement of the temporal evolution  $V_t$  shift of a 10-K P/E cycled cell. In curve (a), the cell is at RT in the first  $10^4$  s and is at RD in the following  $10^4$  s. The sequence is reversed in curve (b). Bias condition of RT and RD are  $V_g/V_d/V_s = 0\text{V}, 0\text{V}, 0\text{V}$  and  $V_g/V_d/V_s = 3\text{V}, 1.6\text{V}, 0\text{V}$ , respectively.  $L_g = 0.5 \mu\text{m}$ .

in Table II) is the dominant mechanism. According to the tunnel front model [15], this detrapping current follows a  $t^{-1}$  time dependence. Its  $V_t$  evolution would have a logarithmic time dependence ( $\Delta V_t \propto \log(t)$ ) accordingly.

If a sufficiently high  $V_g$  is applied (e.g., for cells belonging to the same word-line of a selected-reading cell, GD mode), an inversion layer is induced beneath the gate. The positive oxide charges will cause the inversion electrons to tunnel into the trapping nitride at such high vertical field. According to the positive charge-assisted tunneling (PCAT) model [16], this injection current ( $J_{\text{cat}}$  in Table II) will follow the  $t^{-p}$  law where  $t$  is the elapsed time, and  $p$  is around 0.7. The  $V_t$  shift, thus, has a power-law time dependence  $t^n$  with  $n = 1 - p \approx 0.3$ .

As for the read cell, in which the read  $V_g$  and  $V_d$  are applied (RD mode), channel electrons will be accelerated by the lateral field near the drain region. Both the high drain field and the positive oxide charges will enhance electron injection into nitride. This positive oxide charge-enhanced hot-electron injection will cause a positive  $V_t$  drift with the same time dependence as GD, i.e.,  $\Delta V_t \propto t^n$ .

## IV. RESULTS AND DISCUSSIONS

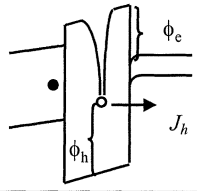
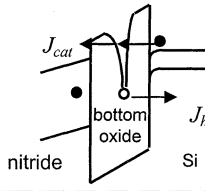
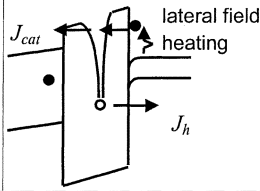
### A. Bias Dependence of Read Disturb

Fig. 4 shows the GD characteristics of a 10-K P/E cycled cell. The applied gate bias is from  $V_g = 0\text{V}$  (RT mode) to  $V_g = 4\text{V}$ . It is found that  $V_t$  shift shows linear dependence on  $\log(t)$  at  $V_g = 0\text{V}$ . This has been well modeled by hole tunnel detrapping [15]. However, accelerated  $V_t$  drift, which exhibits a power-law time dependence  $t^n$ , is observed at  $V_g = 4\text{V}$ . It would be clear if we show the  $\Delta V_t$  versus disturb time on a log-log plot, as indicated by the open square and the right-axis in Fig. 4. Since the time evolution of  $V_t$  drift would show a power-law time dependence  $t^n$  if PCAT dominates, the accelerated  $V_t$  drift with its unique time dependence confirms that positive charge-assisted electron tunneling into the trapping nitride takes place at high vertical field, in addition to hole tunnel detrapping. The  $V_g$  dependence of  $\Delta V_t$  is shown in Fig. 5.

Fig. 6 shows the RD characteristics of a 10-K P/E cycled cell with  $L_g = 0.5 \mu\text{m}$ . The drain bias is from 1.7–2.3 V, and the gate bias is 3 V.  $V_t$  shift versus drain bias is shown in Fig. 7. Since the

TABLE II

SCHEMATICS OF BAND DIAGRAMS AND CARRIER TRANSPORT MECHANISMS AT VARIOUS BIAS CONDITIONS. SOLID CIRCLES REPRESENT THE RESIDUAL ELECTRONS IN TRAPPING NITRIDE AND CHANNEL ELECTRONS IN SILICON, AND OPEN CIRCLES REPRESENT THE POSITIVE TRAPPED CHARGES IN BOTTOM OXIDE.  $\phi_h$  AND  $\phi_e$  REPRESENT THE HOLE AND ELECTRON TUNNELING BARRIER HEIGHTS, RESPECTIVELY.  $m_h$  AND  $m_e$  REPRESENT THE HOLE AND ELECTRON TUNNELING EFFECTIVE MASSES, RESPECTIVELY.  $N_h$  IS THE EFFECTIVE POSITIVE OXIDE CHARGE VOLUME DENSITY. THREE CARRIER TRANSPORT PATHS ARE ILLUSTRATED. (a) POSITIVE OXIDE CHARGE-TUNNEL DETRAPPING AT LOW OXIDE FIELD ( $J_h$ ). (b) POSITIVE OXIDE CHARGE-ASSISTED ELECTRON INJECTION INTO NITRIDE ( $J_{cat}$ ), IN ADDITION TO  $J_h$ , AT HIGH OXIDE FIELD. (c) POSITIVE OXIDE CHARGE-ASSISTED ELECTRON INJECTION INTO NITRIDE ( $J_{cat}$ ), IN ADDITION TO  $J_h$ , AT HIGH LATERAL FIELD. HOT CARRIERS ARE GENERATED VIA LATERAL FIELD HEATING. TIME EVOLUTIONS OF THEIR INJECTION CURRENTS AND  $V_t$  SHIFTS ARE ALSO FORMULATED

Hole tunnel detrapping	Positive charge-assisted electron tunneling (PCAT)	
		
Low vertical field Low lateral field	High vertical field Low lateral field	High vertical field High lateral field
Room temp. drift $V_g/V_d=0V/0V$	Gate disturb $V_g/V_d=3V/0V$	Read disturb $V_g/V_d=3V/2V$
$J_h \propto \frac{N_h}{\sqrt{m_h \phi_h}} t^{-1}$ [15]	$J_{cat} \propto N_h t^{-p}$	$p = \sqrt{\frac{m_e \phi_e}{m_h \phi_h}} \sim 0.7$ [16]
$\Delta V_t^h(t) \propto \frac{N_h}{\sqrt{m_h \phi_h}} \log(t)$	$\Delta V_t^{cat}(t) \propto N_h t^n$	$n = 1-p \sim 0.3$

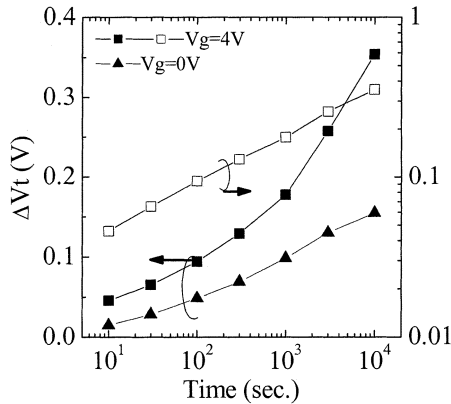


Fig. 4.  $V_t$  shift versus disturb time of a 10-K P/E cycled cell at various gate biases. The drain and source are grounded and  $L_g = 0.5 \mu\text{m}$ .

nitride conduction band edge is 2.1 eV above the silicon conduction band edge, sufficient lateral-field heating is necessary to make the channel electrons inject into the trapping nitride, if the vertical field is low. The drastically enhanced RD found at high drain bias ( $V_d = 2.3 \text{ V}$ ) can be explained by the channel hot-electron injection into the trapping nitride. In addition, the  $t^n$  time dependence of  $V_t$  drift also indicates that PCAT process is involved. As  $V_d$  decreases, the hot-carrier effect is too weak to make the electrons inject into the trapping nitride. The  $V_t$  shift is then dominated by hole tunnel detrapping, which is the same as RT (Fig. 4,  $V_g \approx 0 \text{ V}$ ). It is shown that for  $V_d < 2 \text{ V}$ , the  $V_t$  shift due to read disturb is almost the same (Fig. 7), and  $\Delta V_t$  is proportional to  $\log(t)$  (Fig. 6). Furthermore, a smaller  $V_t$  shift at  $V_g/V_d = 3 \text{ V}/1.7 \text{ V}$  than at  $3 \text{ V}/0 \text{ V}$  can be explained by reduced vertical field in the former case [17].

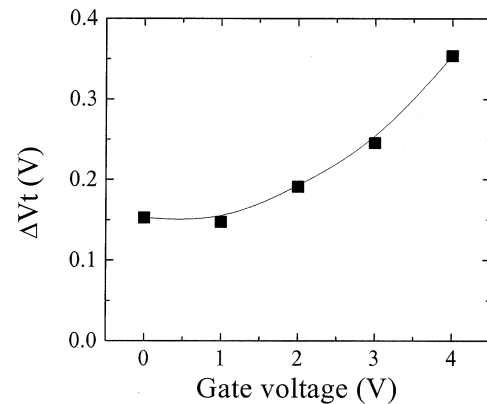


Fig. 5.  $V_t$  shift versus applied gate bias of a 10-K P/E cycled cell. The drain and source are grounded, and the disturb time is  $10^4 \text{ s}$ ,  $L_g = 0.5 \mu\text{m}$ .

$V_d$  acceleration is widely used for hot-carrier lifetime projection. As shown in Fig. 8,  $V_d$  acceleration effect is observed at high  $V_d$ . However, gate-enhanced read disturb, instead, dominates at low drain bias where the hot-carrier effect is insignificant. The  $V_t$  shift is almost independent of  $V_d$  in this regime (e.g.,  $V_d < 2 \text{ V}$  with  $L_g = 0.5 \mu\text{m}$ ). The extrapolated lifetime based on the results at high  $V_d$  regime will be overestimated if we use a low drain voltage during read operation, where the hot-carrier effect is very weak.

### B. Channel Length Scaling Effect

It is known that the hot-carrier effect will become more serious in a short-channel device. Fig. 9 shows the RD results of three cells with gate length ( $L_g$ ) of 0.5, 0.4, and 0.3  $\mu\text{m}$ , respectively. Increased RD is observed in a shorter gate length cell. In addition, the temporal evolution of the  $V_t$  shift follows a  $\log(t)$  time

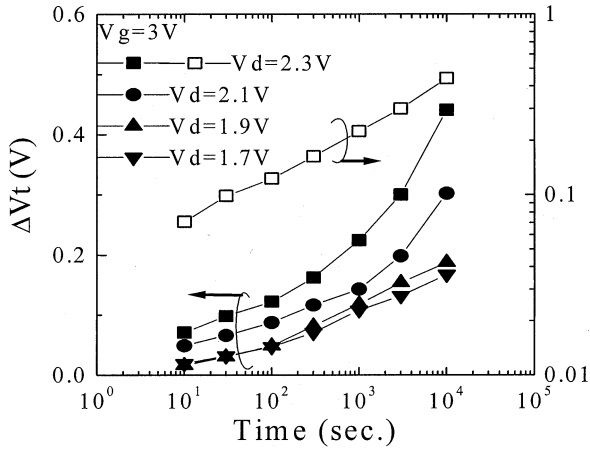


Fig. 6.  $V_t$  shift versus disturb time of a 10-K P/E cycled cell at various drain biases.  $V_g/V_s = 3\text{ V}/0\text{ V}$  and  $L_g = 0.5\ \mu\text{m}$ .

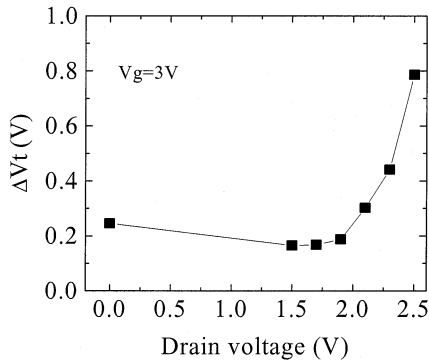


Fig. 7. Read disturb-induced  $V_t$  shift versus applied drain bias of a 10-K P/E cycled cell. The disturb time is  $10^4\text{ s}$ ,  $V_g/V_s = 3\text{ V}/0\text{ V}$  and  $L_g = 0.5\ \mu\text{m}$ .

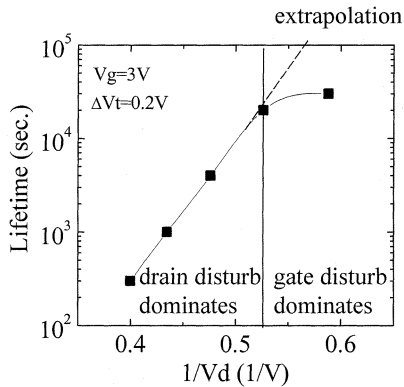


Fig. 8. Lifetime projection by  $V_d$ -acceleration approach. Lifetime is defined as the disturb time resulting in a  $V_t$  shift of 0.2 V.  $V_g/V_s = 3\text{ V}/0\text{ V}$  and  $L_g = 0.5\ \mu\text{m}$ . As shown in the figure, lifetime is overestimated in the low- $V_d$  regime where gate-enhanced disturb dominates.

dependence for  $L_g = 0.5\ \mu\text{m}$ , and exhibits a  $t^n$  time dependence for  $L_g = 0.3\ \mu\text{m}$ . It implies that strong hot-carrier enhanced RD dominates the  $V_t$  shift as the channel length is scaled down.

### C. Temperature Effect

We also investigate the temperature effect on RD. It is reported that hot-carrier-induced gate current injection increases with decreasing temperature, even with a low drain bias (e.g.,  $V_d < 2\text{ V}$ ) [18]. Fig. 10 shows that RD is almost independent of temperature in a long-channel device ( $L_g = 0.5\ \mu\text{m}$ ), in which

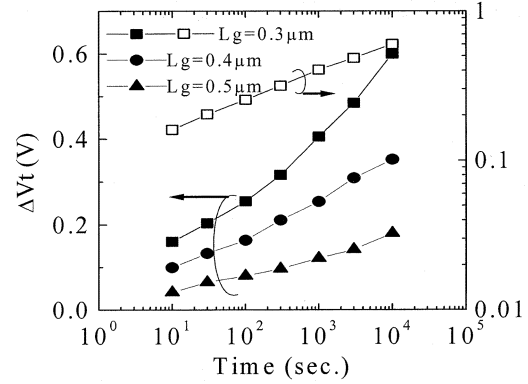


Fig. 9.  $V_t$  shift versus disturb time of 1-K P/E cycled cells with different channel lengths.  $V_g/V_d/V_s = 3\text{ V}/1.6\text{ V}/0\text{ V}$ .

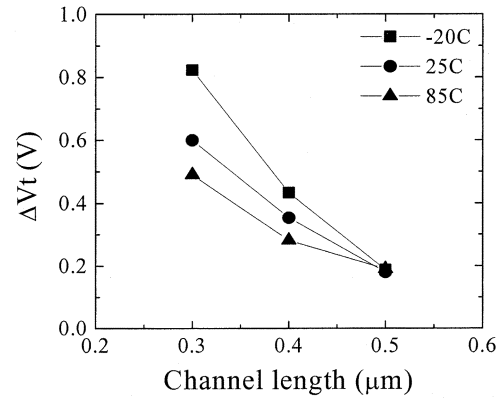


Fig. 10. Temperature effect on read disturb of 1-K P/E cycled cells. Devices with gate lengths of  $0.5\ \mu\text{m}$ ,  $0.4\ \mu\text{m}$ , and  $0.3\ \mu\text{m}$  are characterized. The applied gate and drain biases are 3 and 1.6 V, respectively, and the disturb time is  $10^4\text{ s}$ .

hole tunnel detrapping dominates [so it is termed as room-temperature drift (RT)]. On the other hand, enhanced RD is observed at low temperature in a shorter device ( $L_g = 0.3\ \mu\text{m}$ ), where hot-carrier-enhanced degradation dominates.

### D. Channel Doping Effect

Channel doping is usually used to control short-channel effect and to improve hot-electron injection efficiency. It is reported that GD is worse in a high- $V_{ti}$  cell [19]. RT, GD, and RD are characterized for cells with initial threshold voltages of 1.2, 2, and 3 V, respectively, in our study. During GD and RD measurements, we use  $(V_g - V_{ti}) = 1.5\text{ V}$ , which is the same gate overdrive that conducts the same read current for each cell. The read drain bias is 1.6 V. The results are plotted in Fig. 11.

As shown in Fig. 11, two major features are observed. First, RT is almost the same for these three cells, while GD and RD increase with increased  $V_{ti}$  (channel doping). Second,  $V_t$  shift is strongly enhanced by  $V_g V_d$  in high- $V_{ti}$  cells; however, it is less affected by  $V_g V_d$  in a low- $V_{ti}$  cell. More severe RD in a high- $V_{ti}$  cell is due to a stronger hot-carrier effect, since the channel doping concentration is higher. As tunneling detrapping of positive oxide charges is responsible for the  $V_t$  drift in low vertical/lateral field, RT is expected to be almost the same in these three cells (solid squares in Fig. 11). Since the same gate overdrive implies the same inversion charge density in these cells, we need to find out the cause that results in the enhanced GD in a high- $V_{ti}$  cell.

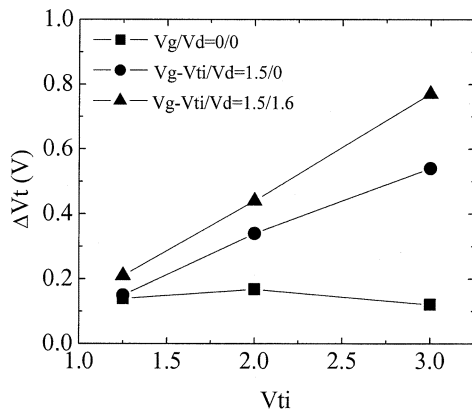


Fig. 11. Initial threshold voltage versus room-temperature drift, gate disturb, and read disturb. The cells undergo 10-K P/E cycles. The bias conditions for GD and RD are  $V_g - V_{ti}/V_d/V_s = 1.5 \text{ V}/0 \text{ V}/0 \text{ V}$  and  $1.5 \text{ V}/1.6 \text{ V}/0 \text{ V}$ , respectively, and  $V_g = V_d = V_s = 0 \text{ V}$  for RT. The disturb time is  $10^4 \text{ s}$  and  $L_g = 0.5 \mu\text{m}$ .

The oxide (vertical) field can be formulated as (without trapped charges)

$$E_{\text{ox}} = \frac{V_g - V_{ti}}{T_{\text{ox}}} + \frac{\sqrt{2\varepsilon_{\text{si}}(2\phi_f)qN_A}}{\varepsilon_{\text{ox}}}$$

in strong inversion, where  $E_{\text{ox}}$  is the oxide field,  $T_{\text{ox}}$  is the equivalent oxide thickness of the ONO stack,  $N_A$  is the effective channel doping concentration, and  $\phi_f$  is the potential difference between the Fermi level (with acceptor concentration of  $N_A$ ) and the intrinsic Fermi level.  $\varepsilon_{\text{ox}}$  and  $\varepsilon_{\text{si}}$  are the dielectric constants of oxide and silicon, respectively. Obviously, even with the same gate overdrive ( $V_g - V_{ti}$ ), the vertical field is stronger in a cell with higher channel doping  $N_A$ . The higher vertical field will enhance the charge tunneling rate [17]. This explains the larger  $V_t$  shift in a high- $V_{ti}$  cell.

## V. CONCLUSION

RD-induced  $V_t$  instability is investigated in this work. It is found that at low vertical/lateral field, tunneling detrapping of cycling-induced positive oxide charges dominates the  $V_t$  shift, which is insensitive to temperature. A hot-carrier effect (drain field heating) dominates the  $V_t$  shift in a short-channel device, especially at low temperature. Gate voltage-enhanced  $V_t$  drift occurs even in a long-channel device, and it is much worse in a high- $V_{ti}$  cell. All of the results are attributed to positive oxide charges, which are created during P/E stress, that enhance electron injection into the trapping nitride. Process improvements and operational optimizations have been shown effective in suppressing these disturbs [20], [21].

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