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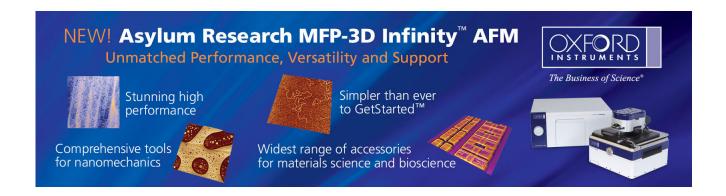
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## Shrinkage- and refractive-index shift-corrected volume holograms for optical interconnects

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The Bragg mismatching condition for volume holograms occurs because of the changes in the thickness and the refractive index of holographic recording materials during the recording and reconstruction procedures. We propose an improved compensation method to physically correct these effects in the fabrications of volume holograms for optical interconnects. In order to show the validity of this method, Slavich photographic plate VRP-M is used to fabricate optical interconnects. The correction of the Bragg diffraction angle shift of about 2.10°, which is induced by 6.14% film shrinkage and 0.06 refractive index shift, is successfully demonstrated with the surface-normal configuration. A shrinkage- and refractive-index shift-corrected volume hologram with 23% diffraction efficiency is experimentally confirmed. The methodology proposed is applicable to other phase media when the associated film shrinkage and refractive-index shift data are experimentally determined. © 2002 American Institute of Physics. [DOI: 10.1063/1.1502022]

Thick recording materials such as silver halide emulsion, dichromated gelatin, and photopolymer, can be used to record volume holograms for optical interconnects. The thickness and the refractive index of these recording materials can change significantly after optical exposure and postprocessing. According to the coupled-wave theory, the diffraction properties of a volume hologram are strongly dependent on the wavelength of the reconstructed light, the thickness, and the refractive index of the recording material. Thus, the operation characteristics of thick volume holograms are strongly influenced by the shifts in the thickness and the refractive index of the photographic film, especially when the operating wavelength is different from that of the recording beam such as for optical interconnect applications. There are several articles proposed to reduce the influences of the thickness shift.<sup>2-6</sup> Some are related to the procedures of the postprocessing.<sup>2-5</sup> Recently, Zhao et al.<sup>6</sup> proposed a corrected condition in the geometry of optical setup for recording volume holograms, and he demonstrated its feasibility with Dupont photopolymer HRF-600X001. There is no reference related to how to simultaneously correct the influences of the thickness shrinkage as well as the refractiveindex shift. In this letter, an improved compensation method to physically correct the influences due to both the thickness shift and the refractive-index shift is proposed, based on Zhao's technique. In addition, volume holograms for optical interconnects are fabricated with silver halide emulsion to demonstrate its validity.

The geometry of optical configuration for recording a volume hologram is shown in Fig. 1. Two light beams  $R_1$  and  $S_1$  with wavelength  $\lambda_1$  are incident on a recording ma-

terial at  $\theta_{r1}$  and  $\theta_{s1}$ . The recording material consists of a substrate with refractive index  $n_s$  and a photographic film with refractive index  $n_{f1}$  (at  $\lambda_1$ ) and thickness  $d_1$ . Their refractive angles in the photographic film are  $\theta'_{r1}$  and  $\theta'_{s1}$ , and the spacing of the interference fringes is  $\Lambda_1$ . After the exposure, the recording material is postprocessed for developing. Then the thickness of the photographic film and its refraction index may shift to different values. Assume the thickness, the spacing of the interference fringes, and the slant angle are changed from  $(d_1, \Lambda_1, \psi_1)$  to  $(d_2, \Lambda_2, \psi_2)$ , as shown in Fig. 2. In order to satisfy the Bragg matching condition for volume holograms, the recording and reconstruction conditions must be carefully designed. For optical interconnect applications, this problem becomes more serious because volume hologram are always fabricated by the technique of shorter-wavelength construction for longerwavelength reconstruction.7 Thus, the reconstructed light with wavelength  $\lambda_2$  ( $>\lambda_1$ ) must be incident at a specially designed angle  $\theta_{r2}$ , as shown in Fig. 3. If the outgoing signal

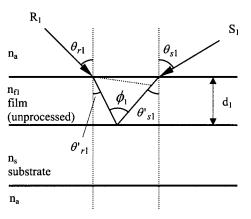


FIG. 1. Geometry for recording a volume hologram.

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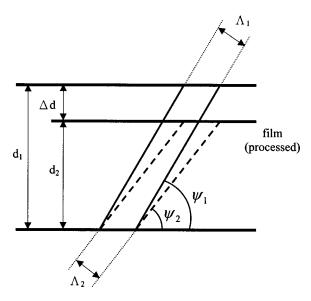


FIG. 2. Schematic representation of thickness shift of photographic film before and after processing.

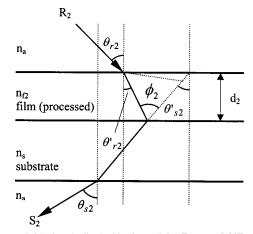
wave  $S_2$  with diffracted angle  $\theta_{s2}$  is desired, then the angles  $\theta'_{r2}$  and  $\theta'_{s2}$  can be calculated with Snell's law, and the angles  $\theta'_{r1}$  and  $\theta'_{s1}$  can be derived with the K-vector diagram for Bragg matching condition, as shown in Fig. 4. Here,  $K_1$  is the grating vector of the volume hologram before processing, its initial and end points are on the circular arc with radius  $2\pi/\Lambda_1$ . And  $K_2$  is the grating vector of the volume hologram after processing. Since a light beam with wavelength  $\lambda_2$  is used to reconstruct this volume hologram, its original and end points are located on another circular arc with radius  $2\pi/\Lambda_2$ . Finally, the values of  $\theta_{r1}$  and  $\theta_{s1}$  can be calculated by using Snell's law again. Hence, we have

$$\theta_{r1} = \sin^{-1} \left\{ \frac{-n_{f1}}{n_a} \cos \left[ \psi_1 + \sin^{-1} \left( \frac{\lambda_1}{2n_{f1}\Lambda_1} \right) \right] \right\}, \tag{1}$$

and

$$\theta_{s1} = \sin^{-1} \left\{ \frac{n_{f1}}{n_a} \cos \left[ \psi_1 - \sin^{-1} \left( \frac{\lambda_1}{2n_{f1}\Lambda_1} \right) \right] \right\}, \tag{2}$$

$$\psi_1 = \tan^{-1} \left\{ \frac{(-1)d_1}{d_2} \cot \left[ \frac{1}{2} (\theta'_{r2} - \theta'_{s2}) \right] \right\},\tag{3}$$



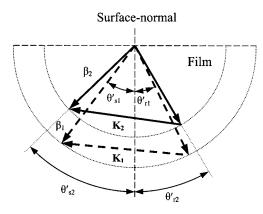


FIG. 4. K-vector diagram for the compensation study.

$$\Lambda_1 = \frac{\lambda_2 \sin \psi_1}{n_{f2} (\sin \theta'_{r2} + \sin \theta'_{s2})},\tag{4}$$

$$\theta_{r2}' = \sin^{-1}\left(\frac{n_a}{n_{f2}}\sin\theta_{r2}\right),\tag{5}$$

$$\theta_{s2}' = \sin^{-1}\left(\frac{n_a}{n_{f2}}\sin\theta_{s2}\right),\tag{6}$$

 $n_{f1}$  and  $n_{f2}$  are the refractive indices at  $\lambda_1$  and  $\lambda_2$ , respectively. It is obvious from Eqs. (1) and (2) that  $\theta_{r1}$  and  $\theta_{s1}$  can be calculated if the experimental conditions of  $n_a$ ,  $n_{f1}$ ,  $d_1$ ,  $n_{f2}$ ,  $d_2$ ,  $\theta_{r2}$ , and  $\theta_{s2}$  are specified. These equations serve as design formulas for the precompensated volume holograms.

Optical interconnects with the specifications of  $\theta_{r2}$ =0°,  $\theta_{s2}$ =45°, and  $\lambda_2$ =632.8 nm are fabricated. A solidstate laser with 532 nm wavelength is used to record the interference fringes on the Slavich photographic plate (VRP-M) whose optimal exposure is 30  $\mu$ J/cm<sup>2</sup>. After the optical exposure, it is processed with the technique presented by Sazonov and Kumonko<sup>8</sup> (with SM-6 developer and PBU-Amidol bleach). An ellipsometer (Model: nkd-6000<sup>TM</sup>, Aquila Instruments Ltd.) and a surface profiler (model: Dektak 3030, Veeco Instruments Inc., Sloan) are used to measure the refractive indices and the thickness of this film before and after processing, respectively. The measured results are shown in Fig. 5. From this figure, we obtain  $n_{f1} = 1.6046$  (at 532 nm),  $n_{f2}$  = 1.6647 (at 632.8 nm),  $d_1$  = 5.70  $\mu$ m, and  $d_2$ 

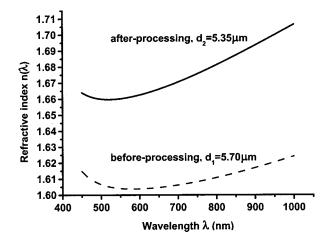


FIG. 5. The measured relation curves between the refractive index vs the wavelength of Salvich VRP-M film before and after processing.

TABLE I. Calculated values for four possible uncompensated conditions.

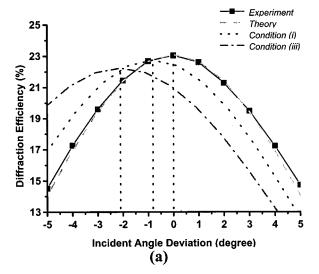
Conditions	$\theta_{r1}$	$\theta_{s1}$	$\theta_{s2}$
(i)	-2.20°	39.26°	46.15°
(ii)	$-2.18^{\circ}$	39.23°	46.11°
(iii)	-3.39°	40.81°	48.04°
(iv)	-3.38°	40.80°	48.03°

=5.35  $\mu$ m, respectively. Substituting these data and  $n_a$ =1 into Eqs (1)–(6), the conditions  $\theta_{r1}$ = -1.46° and  $\theta_{s1}$ = 38.31° are obtained.

Let us consider the conventional cases. First, assume we consider only the thickness shrinkage of the photographic film, then we have two conditions (i)  $d_1 = 5.70 \,\mu\text{m}$ ,  $d_2$ = 5.35  $\mu$ m,  $n_{f1} = n_{f2} = 1.6046$ , and (ii)  $d_1 = 5.70 \mu$ m, and  $d_2 = 5.35 \ \mu \text{m}, \ n_{f1} = n_{f2} = 1.6647$ . In addition, if we even do not consider both the thickness shrinkage and the refractive index shift of the photographic film, then we may have another two conditions (iii)  $d_1 = d_2 = 5.70 \,\mu\text{m}$ ,  $n_{f1} = n_{f2}$ = 1.6046, and (iv)  $d_1 = d_2 = 5.70 \mu \text{m}$ ,  $n_{f1} = n_{f2} = 1.6647$ . We use these four possible uncompensated conditions to calculate the values of  $\theta_{r1}$ ,  $\theta_{s1}$ , and  $\theta_{s2}$ . And they are calculated and listed in Table I. In these four cases,  $\theta_{s2}$  does not equal to 45° as originally designed. We have plotted the relation curves of the diffraction efficiency versus the deviation of the incident angle from the surface-normal direction, as shown in Fig. 6(a) for conditions (i) and (iii) and Fig. 6(b) for (ii) and (iv), respectively. The theoretical and experimental curves have been added in both Figs. 6(a) and 6(b) for comparisons. It is seen that the experimental curve shows good correspondence with the theoretical curve. It is noted in Fig. 6 that the incident angle for matching the Bragg condition (maximum diffraction efficiency) are deviated from the surface-normal direction by  $-0.81^{\circ}$ ,  $-0.78^{\circ}$ ,  $-2.10^{\circ}$ , and  $-2.09^{\circ}$  for conditions (i), (ii), (iii), and (iv), respectively. If both the thickness shrinkage and the refractive index shift have been taken into account, then there should be no angle deviation.

According to the coupled-wave theory, it can be calculated that if the product of the index modulation strength n' and the thickness  $d_2$  (in micron unit) of the photographic film is larger than 0.24, then the diffraction efficiency will be larger than 90%. In our experiments, the thickness of the photographic film is so thin that the maximum diffraction efficiency is only about 23%. If the thickness  $d_2$  could be increased to 18.77  $\mu$ m and other data would remain unchanged, the diffraction efficiency would be enhanced to 90%.

The influences owning to shifts in the thickness and the refractive index of holographic recording materials have been discussed. An improved compensation method to physically correct these effects in the fabrication of volume holograms for optical interconnects has been proposed. Moreover, volume holograms for optical interconnects are fabricated with silver halide emulsion photographic material



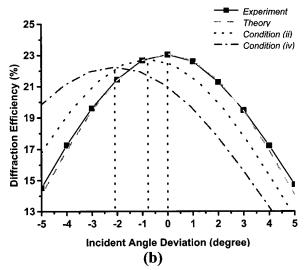


FIG. 6. Variations of diffraction efficiency as function of the angular deviation from the surface-normal Bragg condition for (a) conditions (i) and (iii); and (b) conditions (ii) and (iv).

to show the validity of this method. The methodology proposed is applicable to other phase media when the associated film shrinkage and refractive-index shift data are experimentally determined.

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