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# The variation of the transition temperature of high $T_c$ superconductors by electric field effect

Tzong-Jer Yang<sup>a</sup>\* and Wen-Der Lee<sup>b</sup>

<sup>a</sup>Dept. of Electrophysics, National Chiao-Tung University, Hsinchu, Taiwan 300, Rep. of China

<sup>b</sup>Dept. of Electrical Engineering, Lee-Ming Institute of Technology, Tai-Shan 24314, Taiwan, ROC

High  $T_c$  superconductor is considered as an alternately stacked metal and insulating layers. The electric field effect in this layer superconductor is investigated by the Thomas-Fermi(TF) approximation for the screening of an applied electric field and in the frame of Ginzburg-Landau(GL) theory with the coupling between metallic layers for the transition temperature. The change of the transition temperature related to the coupling constant and applied voltage is derived for two and three  $CuO_2$  layers. The transition temperature for more than three layers is also briefly discussed.

## 1. INTRODUCTION

Electric field induced charge effects to superconductors were studied back to 1960. The changes in conventional superconducting transition temperature  $T_c$  were extremely small due to the electric field penetration length of less than 1Å. The discovered high  $T_c$  oxide superconductors lightened a new wave to get large electric field effects due to the low carrier density and thus larger penetration length of around ten of Å and small coherence length. Therefore, the shift of the critical temperature of the high-temperature superconductors(HTSC) is easily observed.

Up to the present time, the mechanism causing the large electric field effects is most accepted to be the charge modulation model.[1] The usual approach is to use BCS weak coupling model, TF model and Poisson equation for a weaker anisotropic HTSC. Sakai[2]used TF model and regarded HTSC as consisting of alternately stacked N superconducting and N insulating layers(thickness d). The present short report uses TF model,  $g_3$  interlayer coupling model[3] and plasmon BCS-like mechanism to simulate the field effect of the change of  $T_c$  in HTSC. Such an approach is an attempt to relate the electric field effect to the superconducting mechanism. Our approach is briefly described in Sec-II. The results and discussion is presented in Sec-III. Finally, a brief conclusion is made.

## 2. THEORY

Consider a system of a MIS like structure[2], there is a gate insulator with thickness L between the gate metal and HTSC. Assume, there is also an insulating layer  $I_0$  with thickness s between the gate insulator and the topmost  $CuO_2$  layers. The  $CuO_2$  superconducting layer is considered as a two dimensional.

By application of a gate voltage V, there is the charge density Q induced at the gate electrode is  $Q = (v - \phi_1) \frac{\varepsilon_i \varepsilon_b}{\varepsilon_b L + \varepsilon_i s}$ , where  $\phi_1$  is the potential of the topmost  $CuO_2$  layer of HTSC,  $\varepsilon_i$  and  $\varepsilon_b$  are the dielectric constant of the gate insulator and the insulator layer  $I_o$ , respectively.

In the TF model, the induced charge density on the pth  $CuO_2$  layer is  $\delta n_p = e\phi_F \phi(r_{||}, z_p)$ , where  $\phi_F$  is density of states at the Fermi energy and  $z_p = L + S + (p-1)d$ . The origin of z axis is set at the gate metal and positive z axis is downward perpendicular to the surface of HTSC. The electrostatic potential  $\phi(r_{||}, z)$  at the position z may be determined by the Poisson equation[4]

$$\frac{d^2\phi}{dz^2} = -\frac{Q}{\varepsilon_i}\delta(z) + \frac{\varepsilon_s Q - \varepsilon_b (Q - e^2\phi_F\phi(z))}{\varepsilon_b\varepsilon_s\delta(z - z_1)}$$

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$$+\frac{e^2\phi_F}{\varepsilon_s}\sum_{p=2}^N\phi(z)\delta(z-z_p)+\frac{\phi_N(z_N)}{d}\delta(z-z_{N+1})$$

where  $\varepsilon_s$  is the dielectric constant of the insulator inside HTSC.  $\phi_N$  is the potential between Nth  $CuO_2$  layer and (N + 1) th  $CuO_2$  layer. While  $z \neq 0$  and  $z \neq z_p$ , the Poisson equation becomes Laplace equation. The solution of this equation is  $\phi_p(z) = A_p(z-z'_p) + B_p$ , where  $z'_p$  is a reference coordinate and may be taken to be  $z'_p = z_p + \frac{d}{2}$ , By using Gauss's law, and the continuity of the potential

$$\phi_p(z_{p+1}) = \phi_{p+1}(z_{p+1}) = \phi(z_{p+1})$$

then we have the recurrence relation among  $A_p, B_p, A_1$  and  $B_1$ . To understand the change of  $T_c$  of HTSC with the applied electric field, several situations will be discussed separately as following.

(i) only one  $CuO_2$  layer

The Fermi energy will be modified in the TF approximation as  $E_F = E_F^0 + e\phi_1$ , where  $E_F^o$  is the Fermi energy of HTSC without applied voltage.  $T_c$  of HTSC can be gotten by plasmon mechanism.  $T_c = 1.134E_F \exp(-1/\lambda)$ , where  $\lambda$  is a coupling constatn depending on the Fermi energy. The essential change of  $T_c$  with applied electric field shows a linear behavior.

(ii) only two  $CuO_2$  superconducting layers

In this situation, we use th GL theory approach to get  $T_c$  of HTSC. The free energy density f of HTSC according to the  $g_3$  model[3] is

$$f = \alpha \sum_{i=1}^{2} (T - T_{ci}) |\psi_i(z_i)|^2 - g_3[\psi_1^*(z_1)\psi_2(z_2) + c.c.]$$

where  $T_{ci} = 1.134 E_{Fi} \exp(-1/\lambda_i)$ ,

 $E_{Fi} = E_F^o + e\phi_i$ ,  $i = 1, 2, ; \psi_1(z_1)$  and  $\psi_2(z_2)$  are order parameter on 1st and 2nd  $CuO_2$  layer respectively.  $g_3$  is interlayer coupling constant and assumed to be negligible for field effect. The  $T_c$ can be obtained by minimizing free energy with respect to  $\psi_1^*(z_1)$  and  $\psi_2^*(z_2)$ . Then  $T_c$  is

$$T_{c} = \frac{1}{2} [(T_{c1} + T_{c2}) + \sqrt{(T_{c1} - T_{c2})^{2} + 4(g_{3}/\alpha)^{2}}]$$

(iii) More than three  $CuO_2$  layers

According to the result of Sakai[2], the electric field effect is limited within three superconducting  $CuO_2$  layers. The free energy F of more than three  $CuO_2$  layers may be easily expressed. Then minimizing F with respect to  $\psi_p^*$ , we may get the linear system equations and finally  $T_c$  can be gotten.

#### 3. RESULTS and DISCUSSION

For one  $CuO_2$  layer,  $T_c$  increases linearly as electric field E increases in plasmon mechanism. For two  $CuO_2$  layers, the interlayer coupling strength  $g_3$  affects  $T_c$ , but the sign of  $g_3$  does not change  $T_c$ . The change of  $T_c$  deviates slightly from that in one  $CuO_2$  layer. For more than three  $CuO_2$  layers, the  $T_c$  value may be calculated by Ginzbung-Landau theory and taken into account the effect of Debye screening length. The expression of  $T_c$  is not shown here. But the qualitative behavior of the change of  $T_c$  due to the applied field is essentially limited to two or three  $CuO_2$ layers. Because the Debye screening length is almost  $5\mathring{A}$  in most HTSC. Our present approach may be extended to the other mechanisms.

### 4. CONCLUSION

 $T_c$  of HTSC in an applied electric field is derived on the basis of plasmon mechanism in the framework of TF model and  $g_3$  interlayer coupling model. The qualitative behavior of the change of  $T_c$  is shown linearly with the electric field E and independent of the sign of  $g_3$  for one or two  $CuO_2$  layers. The electric field effect on HTSC is essentially limited within the Debye screening length. Our approach may be easily extended to other mechanism.

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