

# On-Wafer Electrooptic Probing Using Rotational Deformation Modulation

W. H. Chen, W. K. Kuo, S. L. Huang, *Member, IEEE*, and Y. T. Huang

**Abstract**—On-wafer electrooptic probing of two-dimensional electric-field vector (E-vector) is demonstrated by using one laser beam and one electrooptic probe. This technique utilizes both compressed-stretched deformation and rotational deformation on the index ellipsoid of the electrooptic crystal. Both experiment and simulation were performed to map the E-vector on a circuit board, and the measurement error is within 2.2%.

**Index Terms**—Nondestructive testing, polarimetry.

## I. INTRODUCTION

ELECTROOPTIC (EO) probing techniques have been advancing very quickly in recent years [1]–[5] due to their superior features in characterization of semiconductor devices and circuits. These techniques allow the measurement of the electric field at any point on the device under test (DUT) [6]. Both analog and digital IC measurements are suitable for EO probing techniques [7], [8]. Some of the EO probing techniques are able to measure not only the electric-field amplitude, but also direction of the electric field [1], [2], [5]. Valuable information such as chamfered bending transmission lines and patch antennas of radio frequency devices can be released [5]. The conventional systems often require two beams or two different EO crystals to differentiate the directions of the electric field under test because only one type of EO modulation effect is utilized in the measurement. The measurements are, thus, inaccurate and complicated because the path length and EO interaction strength of the two probing beams are different [1], [5]. Now, we demonstrate the EO probing technique with one beam and one EO crystal to extract two-dimensional (2-D) electric-field vector (E-vector) using an additional modulation effect, i.e., rotational deformation modulation (RDM). In contrast to the conventional compressed/stretched deformation on the index ellipsoid of the EO crystal, RDM utilizes the rotation of principal axes. With the inclusion of this modulation technique, 2-D E-vector measurement is realized. The commercial software Ansoft Maxwell 3-D Field Simulator is employed to verify our measurements. Good agreement is obtained between experiment and simulation results. This E-vector measurement technique is compact, accurate, and low cost.

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## II. SYSTEM DESCRIPTION

### A. Measurement Principle

1) *Rotational Deformation Modulation*: Depending upon the space group of EO crystals and the applied E-vector, the index ellipsoid of EO crystal experiences compressed/stretched deformation or rotational deformation. The compressed/stretched deformation has been applied in amplitude modulation for a long time [9], so we do not describe it in detail here. Not all EO crystals experience rotational deformation when the electric field is applied. Therefore, the EO crystal for RDM needs to be properly selected and biased.

Consider a beam incident on an electric biased EO crystal along its  $z$  axis. The cross-sectional ellipse of the index ellipsoid equation when cutting perpendicular to the propagation vector can be expressed as

$$\left(\frac{1}{n_x^2} + \alpha\right)x^2 + \left(\frac{1}{n_y^2} + \beta\right)y^2 + 2\gamma xy = 1 \quad (1)$$

where

$$\alpha = r_1 \cdot f(E_x, E_y, E_z)$$

$$\beta = r_2 \cdot f(E_x, E_y, E_z)$$

$$\gamma = r_3 \cdot f(E_x, E_y, E_z)$$

where  $n_x, n_y$  are the intrinsic indexes, and  $r_1, r_2$  and  $r_3$  are Pockels coefficients.  $f(E_x, E_y, E_z)$  is a function of  $E_x, E_y$  and  $E_z$ . The rotation angle  $\theta$ , of the principal axes can be obtained by analyzing (1). If the conditions of  $n_x = n_y$  and  $\alpha = \beta = 0$  both hold, the rotation angle is  $45^\circ$ . In general, the applied electric field is much smaller than the half-wave electric field required for typical EO crystals. Therefore, when  $n_x$  is not equal to  $n_y$ ,  $\theta$  can be estimated as

$$\theta \approx \frac{\gamma}{\frac{1}{n_x^2} - \frac{1}{n_y^2} + \alpha - \beta}. \quad (2)$$

When  $\alpha, \beta$  are both zero or  $\alpha - \beta$  is much smaller than  $(1/n_x^2) - (1/n_y^2)$ ,  $\theta$  is proportional to the electric field amplitude. An RDM system can, therefore, be constructed with this property as shown in Fig. 1. A polarizer and a Wollaston prism are used in this system. According to the analysis by using Jones matrix, the intensity of output beam,  $I_{\text{RDM}}$ , detected by the differential photodetector is

$$I_{\text{RDM}} = -\sin^2(\delta)\sin(4\theta) \quad (3)$$

where  $\delta$  is the phase retardation induced by the static birefringence of the EO crystal. For typical values of on-wafer electric

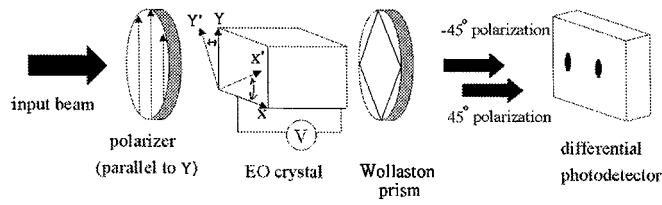


Fig. 1. The rotational deformation modulation system. As the electric field is applied on the EO crystal, the  $X$  and  $Y$  axes of index ellipsoid rotate to  $X'$  and  $Y'$ , respectively.

field,  $4\theta$  is much smaller than 1. Therefore,  $\sin(4\theta)$  can be approximated as  $4\theta$  and  $I_{\text{RDM}}$  becomes proportional to  $\theta$ . This system can, therefore, modulate the input beam linearly. The quarter-wave plate usually used for compressed/stretched deformation modulation (CSDM) system is no longer needed [9], because RDM can work without a retardation bias.

Based on the above discussion, a cubic crystal is not suitable for RDM because its  $n_x = n_y$  and  $\alpha = \beta = 0$ , while EO crystals of  $\bar{4}2m$  and  $3m$  space groups are good for RDM. LiTaO<sub>3</sub> is selected as the EO crystal in our experimental system not only because it is suitable to RDM, but it also has a good compressed/stretched deformation effect. Moreover, its rotational deformation and compressed/stretched deformation are induced by the electric field of different directions.

2) *The Principle of 2-D E-Vector Measurement:* The EO tensor of LiTaO<sub>3</sub> can be expressed as

$$[r_{ij}] = \begin{bmatrix} 0 & -r_{22} & r_{13} \\ 0 & r_{22} & r_{13} \\ 0 & 0 & r_{33} \\ 0 & r_{51} & 0 \\ r_{51} & 0 & 0 \\ -r_{22} & 0 & 0 \end{bmatrix}. \quad (4)$$

At the wavelength of 633 nm,  $n_o$  and  $n_e$  of LiTaO<sub>3</sub> are 2.176 and 2.180, respectively, and the EO coefficients,  $r_{13}$ ,  $r_{33}$ ,  $r_{22}$ , and  $r_{51}$  are  $8.4 \times 10^{-12}$ ,  $30.5 \times 10^{-12}$ ,  $-0.2 \times 10^{-12}$ , and  $22 \times 10^{-12}$  m/V, respectively.

By analyzing the index-ellipsoid deformation of LiTaO<sub>3</sub> under electric-field bias in our propagation path, as shown in Fig. 2(a), it can be proved that  $E_y$  has no RDM effect. Moreover, the  $E_y$  effect on CSDM in our experiment is negligible according to a previous study [1]. The rotation angle for the principal axes of the index ellipsoid  $\theta(E_x)$  is proportional to  $E_x$ , and the phase retardation  $\delta(E_z)$  caused by the EO effect of LiTaO<sub>3</sub> is proportional to  $E_z$ . RDM and CSDM are caused by  $E_x$  and  $E_z$ , respectively, in our probing system, as shown in Fig. 2(b). According to the calculation of Jones matrix, the intensity of detected light,  $I$  in our experimental setup is

$$I = A \cos(2\phi) \cdot \theta(E_x) + B \sin(2\phi) \cdot \delta(E_z) \quad (5)$$

where  $A$  and  $B$  are constants, and  $\phi$  is the polarization angle of the incident beam with respect to the  $z$  axis as shown in Fig. 2(c). The modulation type corresponds to the polarization of incident beam. When  $\phi$  is  $0^\circ$ , i.e., the polarization is along the  $z$  axis, (5) becomes

$$I_{0^\circ} = A \cdot \theta(E_x) = A_1 \cdot E_x \quad (6)$$

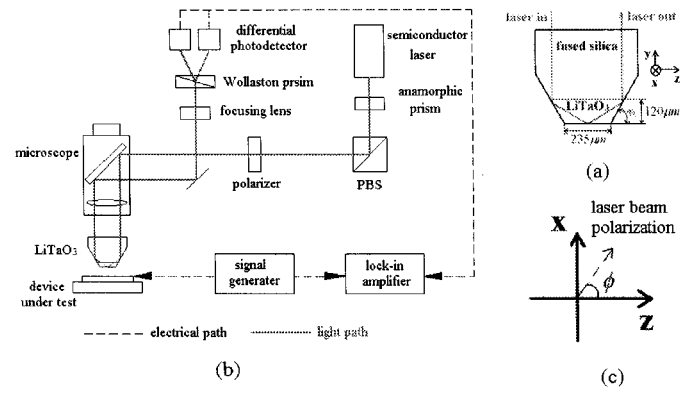


Fig. 2. (a) The propagation path of the incident beam in the LiTaO<sub>3</sub>. (b) The on-wafer EO probing system using RDM and CSDM. (c) The polarization angle of incident beam.  $z$  is the optic axis of LiTaO<sub>3</sub>.

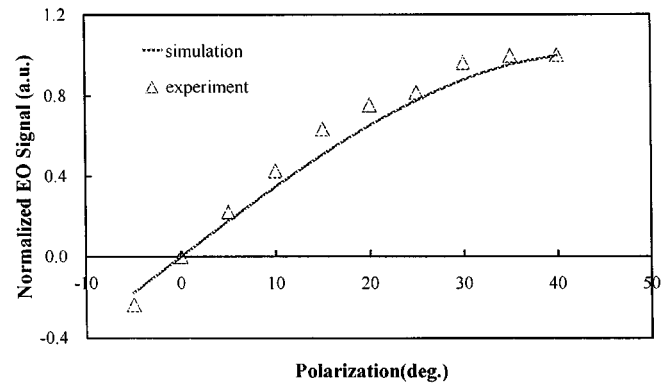


Fig. 3. The EO signal response corresponds to different polarizations.

where  $A_1$  is a constant; the intensity of detected light  $I_{0^\circ}$  is proportional to  $E_x$ , and only RDM exists. When  $\phi$  is  $45^\circ$ , the intensity of detected light becomes

$$I_{45^\circ} = B \cdot \delta(E_z) = B_1 \cdot E_z \quad (7)$$

where  $B_1$  is a constant;  $I_{45^\circ}$  is proportional to  $E_z$ , and only CSDM exists.

The  $A_1$  to  $B_1$  ratio can be determined with high accuracy with the same beam. Therefore, the E-vectors can be obtained by utilizing two different polarizations.

### B. Experimental Setup

The experimental system is shown in Fig. 2(b). A single-longitudinal-mode laser diode is used as the light source (wavelength = 635 nm), which is followed by a 3-to-1 anamorphic prism to complete its divergent angle in vertical and horizontal directions. LiTaO<sub>3</sub> is employed as the EO crystal. The beam path in the LiTaO<sub>3</sub> crystal is shown in Fig. 2(a). The beam propagates in the  $y$ - $z$  plane. A polarizer is placed in front of the EO crystal to control the polarization of the incident beam. Low-noise photodetection is achieved by using a differential photodetector and a lock-in amplifier. A 1-kHz square-wave signal is fed to the device as the test signal.

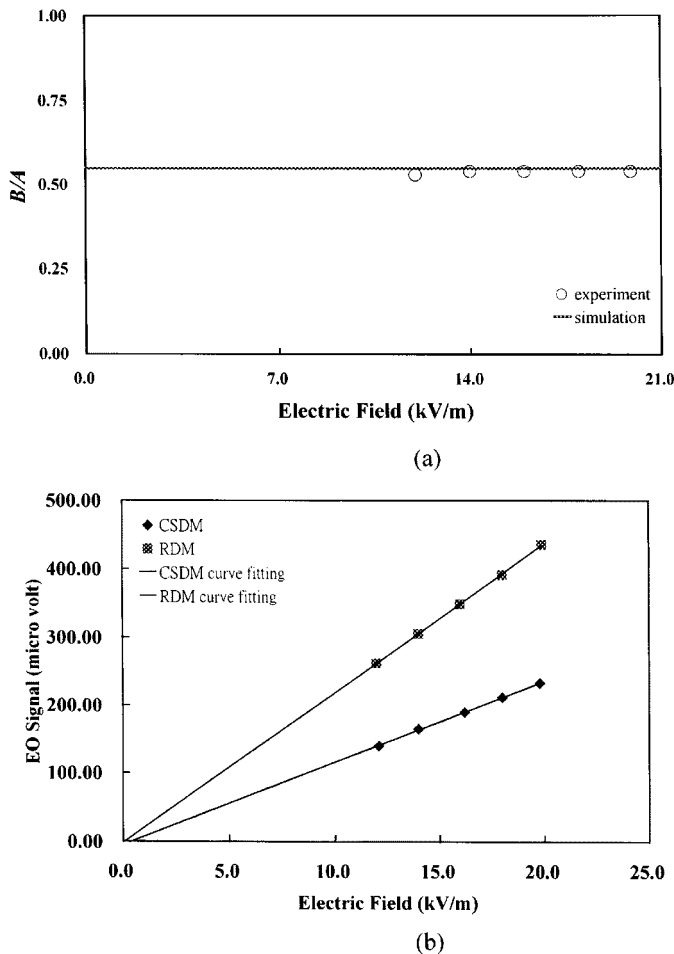


Fig. 4. (a) The simulation and experimental results of  $B$  to  $A$  ratios. (b) Using RDM and CSDM to measure the  $E_x$  and  $E_z$ .

### III. EXPERIMENTAL RESULTS

Several experiments were performed to demonstrate the one-beam EO probing technique for mapping a 2-D E-vector. First, we applied  $E_z$  onto the DUT to measure the EO signal response for different polarizations, as shown in Fig. 3. When the beam polarization was  $0^\circ$ , no EO signal was detected because of no relative phase retardation. This agrees with (5) well. Secondly, constants  $A$  and  $B$  of (5) were determined by varying the applied electric-field strength. Fig. 4(a) shows that the  $B$  to  $A$  ratio agrees well with our simulation results. The average error of this experimental result is 2.2% and the standard deviation is 0.8%. All of these results prove that the RDM technique is feasible. Fig. 4(b) shows the measurement results of RDM and CSDM, respectively. The RDM and CSDM curves both have a linear relationship to the electric field, as expected. This means CSDM and RDM both can work successfully in our experimental system.

The final experiment is a 2-D E-vector measurement of the DUT shown in Fig. 5. The DUT is an unsymmetrical metal gap, which is made on a dielectric substrate. The gap width is 0.5 mm. The measured electric field along "route 1" and "route 2" on the DUT are in good agreement with the simulation by the commercial software Ansoft Maxell 3-D Field Simulator. The  $E_z$  simulation curve along "route 1" is a little incongruous near

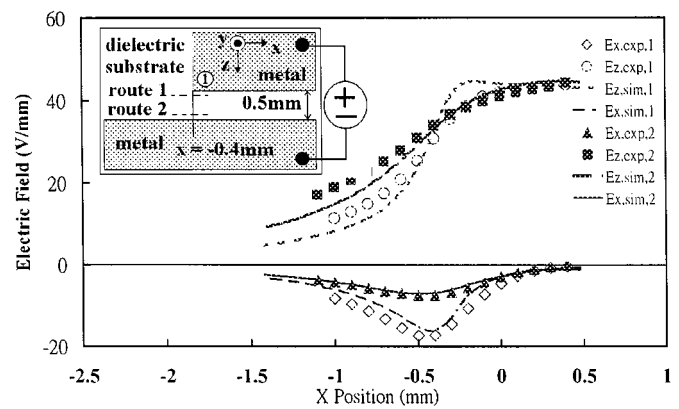


Fig. 5. The inset at the upper left corner is DUT, and "route 1" and "route 2" are the measurement paths. The marks are 2-D electric field measurement results, and the curves are simulation results. The subscripts "1" and "2" represent the data obtained from route 1 and route 2, respectively.

corner  $\circ 1$ , because the curvature of corner  $\circ 1$  was difficult to estimate accurately.

### IV. CONCLUSION

On-wafer EO probing using rotational deformation modulation has been demonstrated successfully. This new probing technique not only simplifies the conventional measurement processes and systems, it also improves measurement accuracy. Therefore, a low-cost approach for E-vector measurement can be achieved. The RDM technique can also be applied to 3-D E-vector measurement with only two beams.

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