

A New Quasi-2-D Model for Hot-Carrier Band-to-Band Tunneling Current

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Abstract—A significant mismatch occurs when we predict the gate-induced drain leakage current (GIDL) by using existing one-dimensional (1-D) models. It's found that the gate-induced drain leakage current is attributed to not only the vertical field but also the lateral field near the drain-to-gate overlap region. Therefore, a new quasi-two-dimensional (quasi-2-D) model considering both the lateral and vertical fields for predicting the gate-induced drain leakage current is proposed by using the drain-induced energy-barrier reduction in our model. The calculated results using the developed quasi-2-D model are in good agreement with measured values for a wide range of gate and drain biases. Therefore, the proposed new model can be used to simulate the hot-carrier band-to-band tunneling current for p-channel flash memory device.

Index Terms—Band-to-band tunneling, gate-induced drain leakage current (GIDL), hot carrier, two-dimensional effect.

I. INTRODUCTION

AS THE GATE-oxide thickness is scaled and the gate electrode is in OFF-state, a significant leakage current passing through the drain electrode can be observed when the drain bias is large but is not large enough for junction breakdown [1]–[5]. This leakage current is so-called “gate-induced drain leakage current (GIDL),” which is commonly attributed to band-to-band tunneling near the gate-to-drain overlap region. Meanwhile, this band-to-band tunneling induced hot-electron injection is recently proposed to be a programming method for p-channel flash memory cell [6].

Since the band-to-band tunneling current is important for flash memory applications, it is necessary that the conducting mechanism should be well-modeled and well-assessed. The model in [1] neglects the dependence of the vertical electric field on the drain doping profile and uses a fixed value of band bending. However, it is well-known that the value of band bending is strongly related to the drain doping concentration and is not a constant value when the drain and gate biases are applied. The model in [4] considers only the vertical field and neglects the lateral field near the gate-to-drain overlap region. The model in [5] only considers the built-in lateral field caused by the lateral gradient of the drain doping concentration and neglects the contribution of the external drain voltage to the lateral field. It is evident that the existing one-dimensional (1-

D) models for calculating the band-to-band tunneling current are not valid if both the gate and drain biases are changed. Therefore, a new quasi-two-dimensional (quasi-2-D) gate-induced drain leakage model is proposed in this paper to calculate the band-to-band tunneling current precisely.

II. EXPERIMENT

The experimental poly-Si gate p-channel MOSFET's were fabricated in n-well on the $\langle 100 \rangle$ p-type Si-wafer with the substrate doping concentration of about $3\text{--}5 \times 10^{15} \text{ cm}^{-3}$. The n-well was formed by the phosphorus implant with a dose of $7.5 \times 10^{13} \text{ cm}^{-2}$ at 80 KeV and the P⁺ S/D was formed by implanting BF₂ with a dose of $2.0 \times 10^{15} \text{ cm}^{-2}$ at 60 KeV. The gate-oxide thickness is 100 Å, the channel width is 25 μm, the channel lengths vary from 1 to 0.8 μm and the length of gate-drain overlap is about 0.1 μm. A 3000 Å poly-gate film with a POC1₃ diffusion is formed and the resistivity is 21.8–22.5 Ω/sq. The tested devices are fabricated using conventional P-MOSFET process without threshold-voltage adjustment.

To measure the GIDL current, different bias conditions are tested and a HP4145B is utilized to measure the drain and gate currents. The source terminal is floating. Note that the variation range of gate/drain bias is large and, therefore, the whole characteristics can be obtained. Measured results are shown in Figs. 1 and 2.

From Figs. 1 and 2, some interesting phenomena are observed. One can see that the band-to-band tunneling currents are different under the same V_{GD} . The conventional 1-D band-to-band tunneling current model [1] can be expressed as

$$I_D = A \cdot E_s \exp(-B/E_s) \quad (1)$$

where A is a preexponential constant; $B = 21.3 \text{ MV/cm}$ [1]; and E_s is expressed by

$$E_s \simeq \frac{V_{GD} - 1.2}{3T_{ox}}. \quad (2)$$

According to (1) and (2), the drain current should be the same under the same V_{GD} , regardless of different V_G or V_D . Apparently, this prediction conflicts with the characteristics shown in Figs. 1 and 2, in which the larger negative drain bias makes the larger band-to-band tunneling current under a constant V_{GD} . When the V_{GD} is fixed, the vertical field near the gate-to-drain overlap region is almost the same regardless of gate biases. However, if the gate voltage is smaller, the drain voltage must be larger under a constant V_{GD} . The larger drain voltage will render larger lateral field near the gate-to-drain overlap region and this, in turn, causes a larger band-to-band

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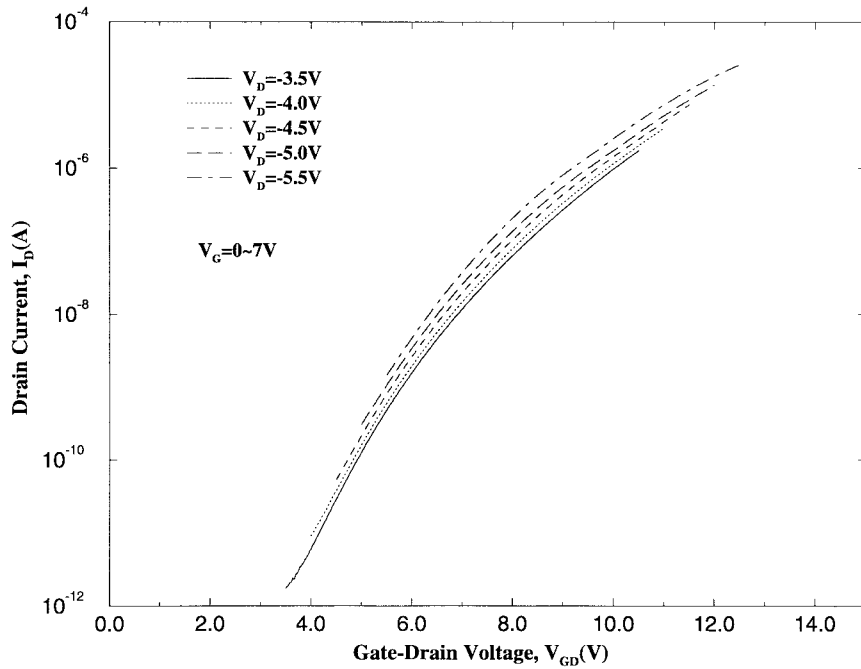


Fig. 1. Band-to-band tunneling current versus the voltage difference between the gate and drain electrode. The gate electrode is swept from 0 to 7 V and the drain electrode is biased at -3.5 , -4.01 , -4.5 , -5.0 , -5.5 V.

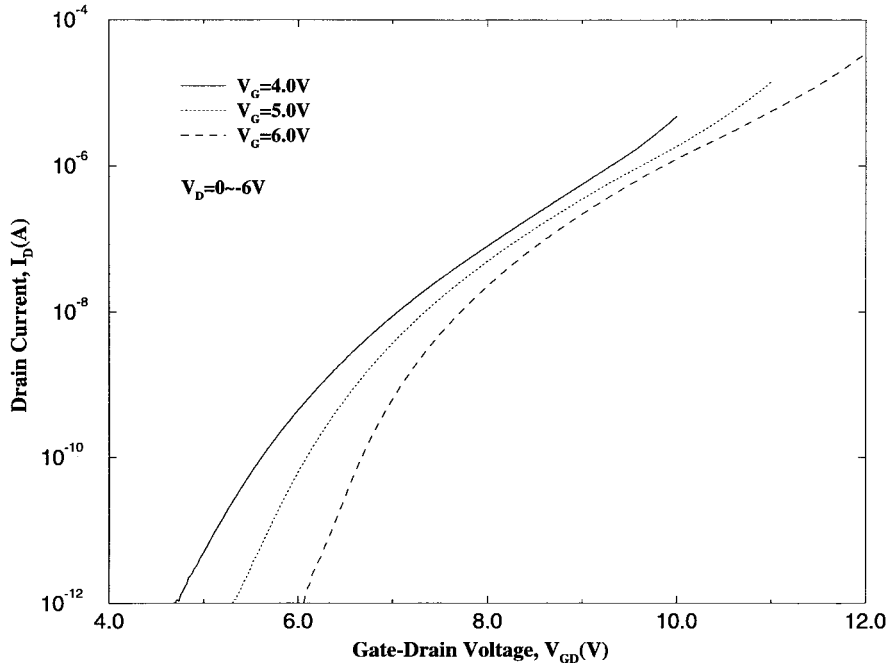


Fig. 2. Band-to-band tunneling current versus the voltage difference between the gate and drain electrode. The drain electrode is swept from 0 to -6 V and the gate electrode is biased at 4.0 , 5.0 , 6.0 V.

tunneling current. The above variations can be further verified by a two-dimensional (2-D) device simulator and the results are shown in Fig. 3. It is noteworthy that the peak vertical field is larger than the peak lateral field for these two bias conditions, and the smaller lateral field obviously influences the band-to-band tunneling unexpectedly. Therefore, if the lateral field is more influential in causing the band-to-band tunneling, the curves shown in Figs. 1 and 2 would split obviously. This fact can not be obtained from all existing

1-D models, including those in [1], [4], and [5]. Note that the polarities of the drain and gate voltages are negative and positive, respectively. For evaluating the band-to-band tunneling current correctly, the effect caused by the lateral field must be taken into consideration.

III. A NEW MODEL AND ITS SIMULATED RESULTS

There are several characteristics in our model: a) the band bending value is dependent on the drain doping concentration

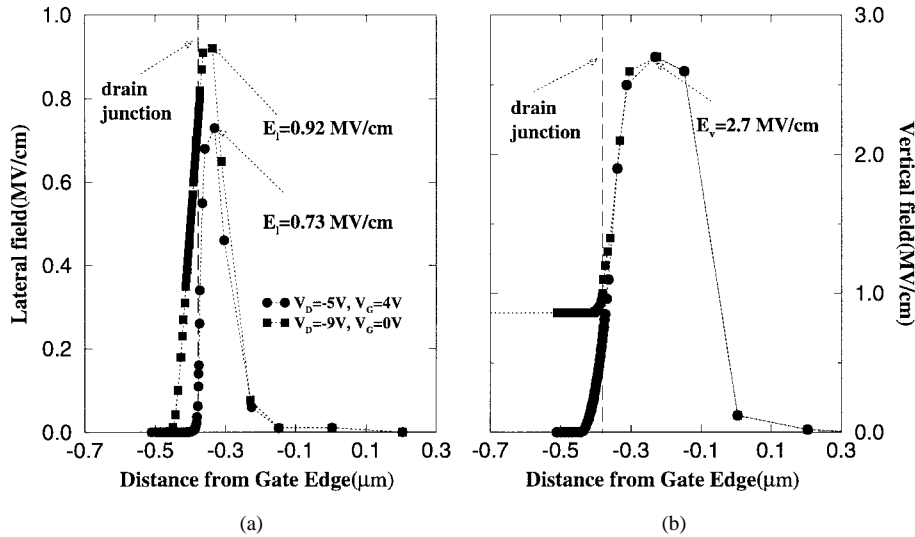


Fig. 3. Simulated (a) lateral and (b) vertical electric fields by a 2-D device simulator at two bias conditions ($V_D = -5$ V, $V_G = 4$ V, and $V_D = -9$ V, $V_G = 0$ V) for the same V_{GD} .

and the gate/drain bias instead of a fixed value, b) the band-to-band tunneling occurs at a shaded depletion region instead at a fixed position, c) the band-to-band tunneling current is a function of V_D and V_G instead of only V_{GD} as proposed in existing 1-D models, and d) the validity of our model calculation is given for a large variation of both the drain and gate biases instead of a single change in the drain bias.

A. Potential and Electric Field Distribution

In this subsection, the concept of the new accurate model is described in detail, and a device cross section is shown in Fig. 4(a). The electric field in the drain depletion region can be calculated by using the deep-depletion approximation and is formulated as follows:

$$E_{si}(Z) = \frac{qN_A}{\epsilon_{si}} Z \quad (3)$$

where

- E_{si} electric field in the depletion region;
- N_A doping concentration in the drain depletion region;
- q electronic charge;
- Z the coordinate normal to the Si-SiO₂ interface, which is equal to 0 at the edge of the drain depletion region and is positive directed from the depleted edge to the Si-SiO₂ interface, as shown in Fig. 4(c).

Note that the doping concentration is assumed to be constant in the drain depletion region for simplicity. Since the depth where the doping distribution of drain is nearly uniform is about 200–300 Å and the width of the deep depletion layer where band-to-band tunneling occurs is only about 100–200 Å. Therefore, the above assumption in our calculations is acceptable.

Moreover, from Gaussian law, the electric displacement across the Si-SiO₂ must be continuous and is described by

the following relationship:

$$\epsilon_{si} E_{si}(Z = W) = \epsilon_{ox} E_{ox} = \epsilon_{ox} \frac{(V_{GD} - V_{FB} - V_{bend})}{T_{ox}} \quad (4)$$

where E_{ox} is the electric field across SiO₂; T_{ox} is the oxide thickness; W is the depletion width, which is equal to $\sqrt{2\epsilon_{si}V_{bend}/qN_A}$; V_{FB} is the flatband voltage; V_{GD} is the potential difference between the gate and drain electrode; and V_{bend} is the band bending value.

One can see that the vertical band bending value could be larger than 1.2 eV or smaller than 1.2 eV, depending on the bias condition and the doping concentration near the gate-to-drain overlap region. According to (3) and (4), the band bending value can be calculated by the following equation:

$$\begin{aligned} V_{bend} &= V_{GD} - V_{FB} + qN_A T_{ox}^2 \epsilon_{si} / \epsilon_{ox}^2 \\ &\quad - \sqrt{(V_{GD} - V_{FB} + qN_A T_{ox}^2 \epsilon_{si} / \epsilon_{ox}^2)^2 - (V_{GD} - V_{FB})^2}. \end{aligned} \quad (5)$$

The position that occurs band-to-band tunneling is not fixed at one position but is extended by a shaded region shown in Fig. 4(c). The potential distributions of conduction and valence bands can be separately expressed by

$$\Phi_v = -\frac{V_{bend}}{W^2} Z^2 \quad (6)$$

$$\Phi_c = -\frac{V_{bend}}{W^2} Z^2 + E_g \quad (7)$$

where E_g is the energy gap of Si.

B. Band-to-Band Tunneling Current

If the drain is biased negatively, a reverse p-n junction (drain to channel) would introduce a large lateral electric

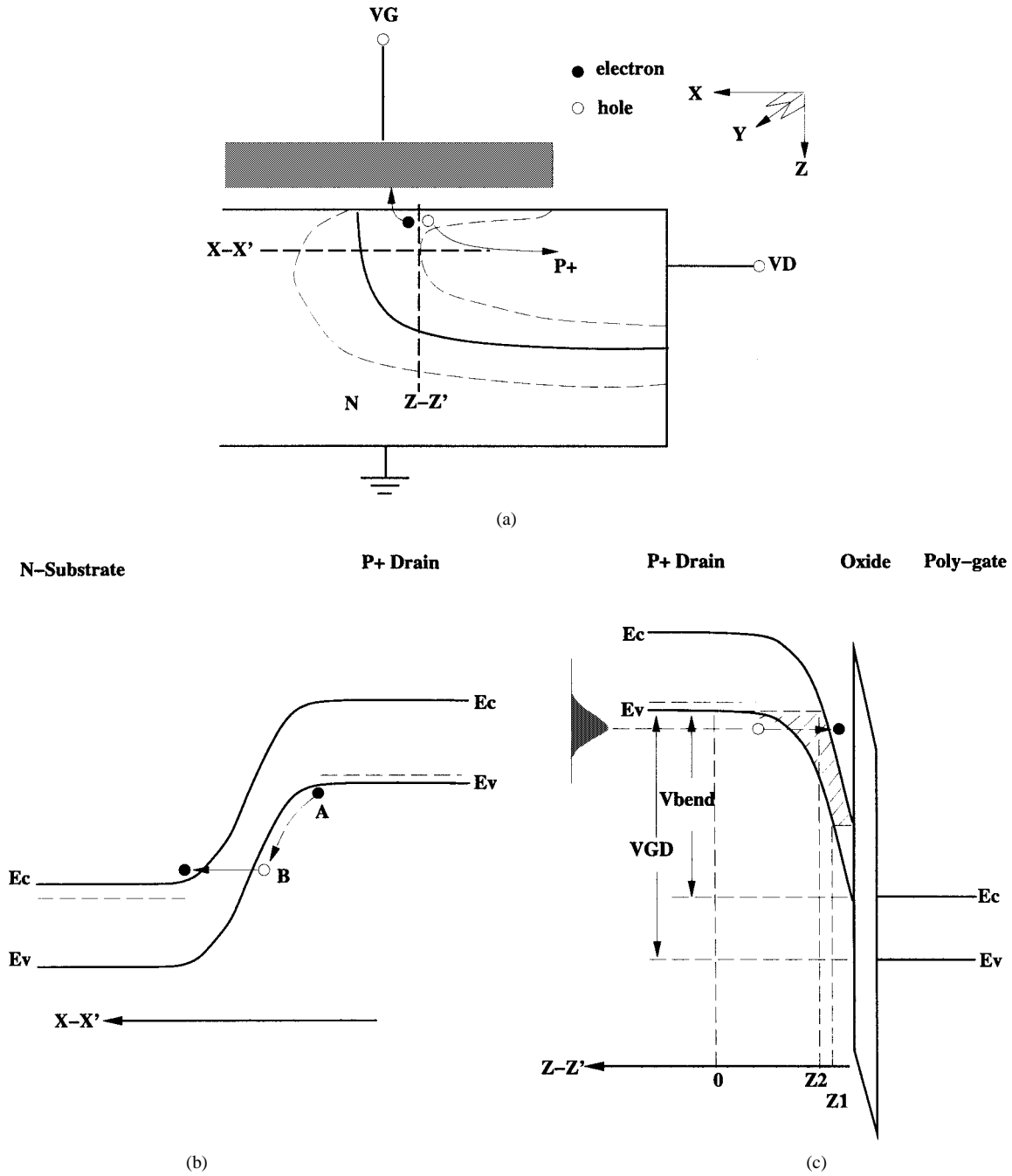


Fig. 4. (a) Cross-sectional view of a planar diode, (b) the energy band diagram in the direction of $x-x'$, and (c) the energy band diagram at the position B based on our new model.

field as shown in Fig. 4(b). This lateral electric field would render electrons in the valence band traveling horizontally, for example, from A to B as depicted in Fig. 4(b) and, in turn, these valence electrons would gain energy from the lateral electric field. If the band-to-band tunneling occurs for these hot valence electrons, the equivalent energy barrier would reduce, as shown in Fig. 4(c). Therefore, the energy barrier which is equal to energy gap for cold electrons could be reduced due to the applied drain bias. We call this phenomenon as the drain-induced energy-barrier reduction effect. Here, we introduce the drain biasing effect on the energy barrier reduction and use a fitting parameter α . The reduced energy barrier E_{gp}

for hot-electron band-to-band tunneling can be expressed as follows:

$$E_{gp} = E_g - \alpha V_D \quad (8)$$

where α is an empirical fitting parameter.

To calculate the band-to-band tunneling current, the tunneling probability as a function of electric field in Si is expressed by using the two-band tunneling theory as follows [7]:

$$P(E_{si}) = \frac{q^2 m^{1/2} E_{si}^2}{18\pi \hbar^2 E_{gp}^{1/2}} \exp\left(-\frac{\pi m^{1/2} E_{gp}^{3/2}}{2\hbar q E_{si}}\right) \quad (9)$$

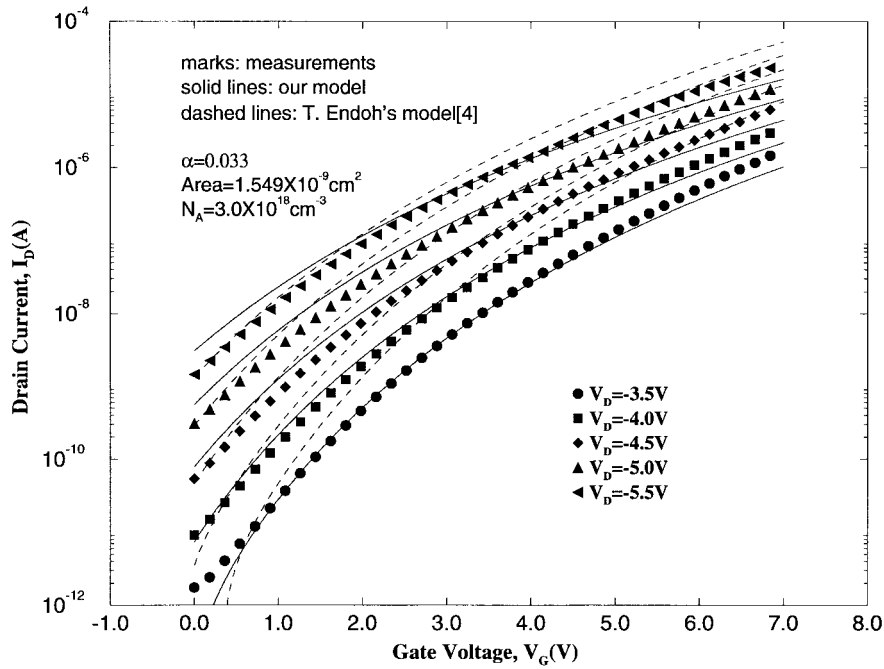


Fig. 5. Comparisons between measured results and calculated currents using our model. The dashed lines are calculated by the model in [4]. The bias conditions are the same as those used in Fig. 1.

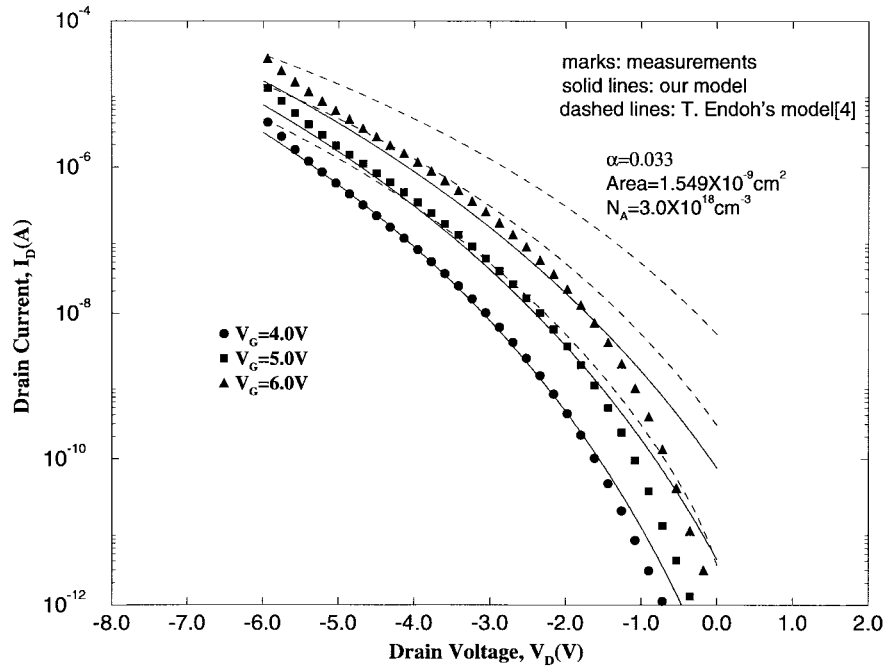


Fig. 6. Comparisons between measured results and calculated currents using our models. The bias conditions are the same as those used in Fig. 2.

where P is the tunneling rate per second per cubic centimeter. The band-to-band tunneling current I_d can be obtained by integrating the tunneling rate through the whole depletion region as

$$\begin{aligned}
 I_d &= A_{rea} \int_0^{Z_1} dz \cdot q \cdot P(E_{si}) \\
 &= A_{rea} \int_{E_1}^{E_2} dE_{si} \frac{dz}{dE_{si}} \cdot q \cdot P(E_{si}) \quad (10)
 \end{aligned}$$

where $Z_1 = W \sqrt{1 - E_{gp}/V_{bend}}$; $Z_2 = W \sqrt{E_{gp}/V_{bend}}$; $E_1 = E_{si}(Z = Z_2)$; $E_2 = E_{si}(Z = W)$ and A_{rea} is the effective area of band-to-band tunneling near the gate-to-drain overlap region in the direction of x and y , as shown in Fig. 4.

C. Comparisons Between Simulated and Measured Results

The band-to-band tunneling characteristics are calculated for various bias conditions. In Fig. 5, the gate bias is swept from 0 to 7.0 V and the drain is biased at -3.5 , -4.0 , -4.5 , -5.0 ,

TABLE I
THE PARAMETERS USED FOR MODEL CALCULATIONS

device's parameter values		fitting parameter values	
$T_{ox}(\text{\AA})$	100	α	0.033
$V_{FB}(V)$	-1.10	$A_{rea}(cm^2)$	1.549×10^{-9}
$N_A(cm^{-3})$	3.0×10^{18}		
$E_g(eV)$	1.12		

-5.5 V. It is obvious that the simulated and measured currents are quite matched for various bias conditions. The calculated band-to-band currents using our model agree well with the measured results over a wide range of band-to-band tunneling current from 10^{-12} to 10^{-5} A. In Fig. 6, the drain bias is swept from 0 to -6.0 V while the gate electrode is biased at 4.0, 5.0, 6.0 V. Because the avalanche breakdown voltage of the drain-substrate junction is about 6.5 V, the maximum value of measured drain bias is -6.0 V to avoid the involvement of the avalanche effect. It is clearly seen that the calculated results match well with the measured curves. The dotted lines shown in Figs. 5 and 6 are calculated by using the model in [4]. Note that the parameters used for calculations are identical between our model and conventional one [4] except that α is set to zero for the conventional model [4], which are listed in Table I. The band-to-band tunneling currents calculated by our quasi-2-D model almost match with experimental results for any bias condition, in which both the vertical and lateral electric fields are significant and can't be neglected. Therefore, our developed band-to-band tunneling current model can be used to predict the subbreakdown current and the band-to-band tunneling current induced hot-carrier injection for flash memory programming.

IV. CONCLUSION

A new accurate quasi-2-D band-to-band tunneling current model considering the effective tunneling barrier lowering has been proposed, which is effective in predicting the subbreakdown current under various gate and drain biases. The calculated currents match well with the measured results over a wide range of current from 10^{-12} to 10^{-5} A. Therefore, our proposed model can be used to simulate the programming of p-channel flash memory, based on hot-carrier band-to-band tunneling.

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