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# Transport phenomena and performance of a plate methanol steam micro-reformer with serpentine flow field design

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## **ABSTRACT**

A numerical investigation of the transport phenomena and performance of a plate methanol steam micro-reformer with serpentine flow field as a function of wall temperature, fuel ratio and Reynolds number are presented. The fuel Reynolds number and  $H<sub>2</sub>O/CH<sub>3</sub>OH$  molar ratio (S/C) that influence the transport phenomena and methanol conversion are explored in detail. In addition, the effects of various wall temperatures on the plates that heat the channel are also investigated. The predictions show that conduction through the wall plays a significant effect on the temperature distribution and must be considered in the modeling. The predictions also indicate that a higher wall temperature enhances the chemical reaction rate which, in turn, significantly increases the methanol conversion. The methanol conversion is also improved by decreasing the Reynolds number or increasing the S/C molar ratio. When the serpentine flow field of the channel is heated either through top plate  $(Y = 1)$  or the bottom plate  $(Y=0)$ , we observe a higher degree of methanol conversion for the case with top plate heating. This is due to the stronger chemical reaction for the case with top plate heating.

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## 1. Introduction

Fuel cells are widely regarded as the most promising energy systems for transport, stationary and portable power sources due to their properties of high efficiency and high environment compatibility. Many analyses, models and numerical simulations have been developed to study proton exchange membrane fuel cell (PEMFCs) which generate electrical energy from hydrogen [\[1–4\].](#page-9-0) In the portable electronics market, fuel cells promise to provide higher power density and longer durability than batteries [\[5\].](#page-9-0) Safety issues, storage problems, and size or portability considerations make pure hydrogen feeding relatively difficult for electronic equipment applications. The combination of a methanol reformer with a proton exchange membrane fuel cell (PEMFC) overcomes the high risk involved in carrying a large quantity of hydrogen, and is thus a promising choice for miniaturized portable electronic systems.

The vast literature devoted to the plate methanol steam microreformer has been reviewed on several occasions [\[6–15\]](#page-9-0). As the steam reforming reaction is an endothermic reaction, several researchers have used an electrical heater to supply heat flux to the plate methanol steam micro-reformers using micro-channels patterned on the plates. This has successfully produced hydrogen to supply the fuel cell [\[6–10\].](#page-9-0) A stack of alternate combustion and reforming chambers are separated by plates. Both sides of each plate are coated with a combustion catalyst and a reforming catalyst. The heat from the combustion reaction is used to drive the reforming reaction. As a result, several studies have successfully used catalytic burners to supply thermal energy to the entire micro-reformer which can generate high yields of hydrogen [\[11–15\].](#page-9-0)

Computational simulation and modeling are used extensively in research and industrial applications to obtain a better understanding of the fundamental processes and to optimize designs before building prototypes for engineering applications. Therefore, more theoretical modeling for study of the methanol reformer is in progress. Suh et al. [\[16,17\]](#page-9-0) employed a cylindrical mathematical model of a packed bed reformer to investigate heat and mass transport phenomena in a methanol reformer. The results showed good agreement between theoretical and experimental results, and also found that the internally heated reformer could improve the methanol conversion. The effects of the methanol steam reforming rate in both a packed bed reformer and a wall coated reformer were examined by Karim et al. [\[18,19\].](#page-9-0) The results showed that a wall coated reformer had better heat and mass transfer limitations and higher catalyst activity than a packed bed reformer. Pan and collaborators [\[20,21\]](#page-9-0) developed the numerical models for a plate-fin methanol steam reformer and a bench-scale methanol

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## Nomenclature



autothermal reformer. Their numerical model accurately predicted the methanol conversion and the gas distribution. A plate microreformer model and a radial flow packed bed reformer model were developed by Pattekar and Kothare, respectively [\[22,23\]](#page-9-0). The results demonstrated that a radial flow reformer had better hydrogen production rates and lower pressure drops than the micro-channels of a plate reformer. A cylindrical reformer model to simulate the conversion and temperature distribution in the methanol reformer was developed by Cao et al. [\[24\]](#page-9-0). Their results showed that the appropriate insulation thickness could reduce the heat losses and achieve a small volume and a high power density. Stamps and Gatzke [\[25\]](#page-9-0) developed a reformer with a model PEMFC to study various design and operating parameters on system performance. The simplest model to describe the methanol reformer conversion and the heat and mass transport phenomena was presented by Yoon et al. [\[26\]](#page-9-0). The results showed that appropriate reactor geometry can improve the reactant gas transport and the efficiency of thermal management. More theoretical modeling of steam reforming coupled with catalytic combustion for the plate reactors is currently in progress. A plate reactor model that combines a steam reformer and catalytic combustor was recently examined by several researchers [\[27–29\]](#page-9-0). The systems are fed by hydrocarbons which convert the hydrogen and generate heat. An appropriate numerical model of the micro-reformer with combustor has been developed to explore the heat and mass transport phenomena and conversion efficiency.

In recent years, different types of flow field designs for plate methanol micro-reformers have been used to achieve more efficient methanol conversion. Kundu et al. [\[30\]](#page-10-0) used different flow configurations, including serpentine and parallel flow fields, to improve plate methanol reformer performance. Several researchers have studied plate steam reformers with a parallel flow field which is attractive due to its simplicity [\[31–33\].](#page-10-0) Chen et al. [\[31\]](#page-10-0) used a plate micro-reformer model to analyze the transport phenomena in a methanol steam reformer. A numerical single channel model to analyze various height and width ratios on the micro-reformer performance and reactant gas transport characteristics was presented by Hsueh et al. [\[32\].](#page-10-0) Kim and Kwon [\[33\]](#page-10-0) proposed a onedimensional methanol steam reformer model to study the inner transport phenomena. Park et al. [\[34\]](#page-10-0) presented numerical simulations to study the reforming rate and heat and mass transfer phenomena in a methanol micro-reformer.

The literature cited above has shown that micro-reformer performance can be enhanced by suitable thermo-fluid parameters. However, there has been a limited amount of work investigating the effects of the different flow field designs on thermo-fluid parameters, especially for the serpentine flow field. Therefore, the objective of this paper is to establish a three-dimensional serpentine flow field model of the plate methanol steam micro-reformer to investigate its transport phenomena and methanol conversion efficiency.

## 2. Analysis

In this study, a three-dimensional computational model of heat and mass transfer in a micro-reformer with a serpentine flow field is proposed. The serpentine flow field has eight turns. A schematic illustration of the coordinate system is shown in [Fig. 1](#page-2-0). The channel of the serpentine flow field consists of a flow channel and a catalyst

<span id="page-2-0"></span>

Fig. 1. Schematic diagram of the present study.

layer. In order to simplify the analysis, assumptions are made as follows:

- (1) the flow is steady state;
- (2) the inlet fuel is an ideal gas;
- (3) the flow is laminar and incompressible;
- (4) the catalyst layer is isotropic;
- (5) the chemical reaction occurs only in the catalyst layer;
- (6) thermal radiation and conduction in the gas phase are negligible compared to convection.

According to the descriptions and assumptions above, the basic transport equations for the three-dimensional plate methanol steam micro-reformer are as follows: Continuity equation:

$$
\rho \left( \frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} + \frac{\partial w}{\partial z} \right) = 0 \tag{1}
$$

X-momentum equation:

$$
\varepsilon \rho \left( u \frac{\partial u}{\partial x} + v \frac{\partial u}{\partial y} + w \frac{\partial u}{\partial z} \right) = -\varepsilon \frac{\partial p}{\partial x} + \varepsilon \mu \left( \frac{\partial^2 u}{\partial x^2} + \frac{\partial^2 u}{\partial y^2} + \frac{\partial^2 u}{\partial z^2} \right) + S_u
$$
\n(2)

Y-momentum equation:

$$
\varepsilon \rho \left( u \frac{\partial v}{\partial x} + v \frac{\partial v}{\partial y} + w \frac{\partial v}{\partial z} \right) = -\varepsilon \frac{\partial p}{\partial y} + \varepsilon \mu \left( \frac{\partial^2 v}{\partial x^2} + \frac{\partial^2 v}{\partial y^2} + \frac{\partial^2 v}{\partial z^2} \right) + S_v
$$
\n(3)

Z-momentum equation:

$$
\varepsilon \rho \left( u \frac{\partial w}{\partial x} + v \frac{\partial w}{\partial y} + w \frac{\partial w}{\partial z} \right) = -\varepsilon \frac{\partial p}{\partial z} + \varepsilon \mu \left( \frac{\partial^2 w}{\partial x^2} + \frac{\partial^2 w}{\partial y^2} + \frac{\partial^2 w}{\partial z^2} \right) + S_w
$$
\n(4)

In the momentum equations,  $\varepsilon$  is the porosity of the medium.  $S_u$ ,  $S_v$  and  $S_w$  are corrected terms for the reactant gas flow in the porous material of the catalyst layer of the micro-reformer. Therefore,  $S_u$ ,  $S_v$  and  $S_w$  are different in each computation domain due to the difference in pressure when fluids pass through a porous medium. So,  $S_u$ ,  $S_v$  and  $S_w$  in the catalyst layer are:

$$
S_u = -\frac{\mu u}{k_p} - \frac{\beta u \rho}{2} \sqrt{u^2 + v^2 + w^2}
$$
 (5)

$$
S_{\nu} = -\frac{\mu \nu}{k_p} - \frac{\beta \nu \rho}{2} \sqrt{u^2 + v^2 + w^2}
$$
 (6)

$$
S_w = -\frac{\mu w}{k_p} - \frac{\beta w \rho}{2} \sqrt{u^2 + v^2 + w^2}
$$
 (7)

where  $k_p$  is the permeability and  $\beta$  is the inertial loss coefficient in each component direction. In the analysis, the Ergun equations [\[35\]](#page-10-0) are used to model  $k_p$  and  $\beta$  as below:

$$
k_p = \frac{D_p^2 \varepsilon^3}{150(1 - \varepsilon)^2} \tag{8}
$$

$$
\beta = \frac{3.5(1 - \varepsilon)}{D_p \varepsilon^3} \tag{9}
$$

where  $D_n$  is the diameter of the catalyst particles.

The dynamic viscosity in the gas mixture is calculated from an expression based on kinetic theory as [\[36\]](#page-10-0)

$$
\mu_{\text{mix}} = \sum_{i=1}^{5} \frac{M_i \mu_i}{\sum_{j=1}^{5} M_j \phi_{ij}} \tag{10}
$$

where

$$
\phi_{ij} = \sum_{i} \frac{\left[1 + \left(\frac{\mu_{i}}{\mu_{j}}\right)^{\frac{1}{2}} \left(\frac{M_{w,j}}{M_{w,i}}\right)^{\frac{1}{4}}\right]^{2}}{\left[8\left(1 + \frac{M_{w,i}}{M_{w,j}}\right)\right]^{\frac{1}{2}}}
$$
\n(11)

<span id="page-3-0"></span>Species equation:

$$
\left(u\frac{\partial m_i}{\partial x} + v\frac{\partial m_i}{\partial y} + w\frac{\partial m_i}{\partial z}\right) = D_{\text{eff}}\left(\frac{\partial^2 m_i}{\partial x^2} + \frac{\partial^2 m_i}{\partial y^2} + \frac{\partial^2 m_i}{\partial z^2}\right) + \varepsilon S_c
$$
\n(12)

In the species equation,  $m_i$  denotes the mass fraction of the *i*th species; the calculations have included  $CH<sub>3</sub>OH$ ,  $H<sub>2</sub>O$ ,  $H<sub>2</sub>$ ,  $CO<sub>2</sub>$  and CO. In Eq. (12),  $D_{\text{eff}}$  is the effective diffusion coefficient based on the Stefan–Maxwell equations [\[36\]](#page-10-0). Eq. (13) is employed to describe the influence of the porosity on the diffusion coefficient

$$
D_{\text{eff}} = D_k \varepsilon^{\tau} \tag{13}
$$

The diffusion coefficient  $D_k$  for the methanol steam micro-reformer was derived from the Stefan–Maxwell equations which were used to calculate the mean effective binary diffusivity [\[19\]](#page-9-0). In this work, the porosity  $\varepsilon$  is expressed as 0.38 and 1.00, in the catalyst layer and the flow channel, respectively.  $S_c$  represents the source terms due to the chemical reaction in the catalyst layer. Therefore,  $S_c$  is zero in the flow channel. Furthermore,  $S_c$  differs according to the reactant gases in the catalyst layer.

$$
S_c = M_{w,i}(R_{SR} + R_{\text{rWGS}})(\lambda_i'' - \lambda_i') \tag{14}
$$

where  $\lambda'_i$  and  $\lambda''_i$  are the stoichiometric coefficient for reactant i and product i, respectively, in the reaction.

According to the chemical kinetics of Purnama et al. [\[37\]](#page-10-0), the steam reforming reaction is much faster than the decomposition and water–gas-shift reaction. Therefore, only the steam reforming reaction, Eq. (15), and the reverse water–gas shift reaction, Eq. (16), are considered in this study.

$$
CH3OH + H2O \stackrel{k_1}{\rightarrow} CO2 + 3H2
$$
 (15)

$$
CO_2 + H_2 \frac{k_2}{k_{-2}} CO + H_2 O \tag{16}
$$

In this study, the model for methanol steam reforming is that used by Hsueh and collaborators [\[31,32\],](#page-10-0) and the Arrhenius equation is used to calculate the concentration of reactant gases generated by the chemical reaction.

$$
R_{\rm SR} = k_1 C_{\rm CH_3OH}^{0.6} C_{\rm H_2O}^{0.4} \exp\left(-\frac{E_a}{RT}\right) \tag{17}
$$

$$
R_{\rm rWGS} = k_2 C_{\rm CO_2} C_{\rm H_2} \exp\left(-\frac{E_a}{RT}\right) - k_{-2} C_{\rm CO} C_{\rm H_2O} \exp\left(-\frac{E_a}{RT}\right) \tag{18}
$$

where the steam reforming reaction is a non-reversible reaction and the reverse water–gas-shift reaction is reversible. The constants  $k_1$ and  $k<sub>2</sub>$  are the forward rate constants for the steam reforming reaction and the reverse water–gas-shift reaction, respectively. The constant  $k_{-2}$  is the backward rate constant for the water–gas-shift reaction.

To calculate the local temperature, the energy equations must be solved.

Energy equation:

$$
\rho c_p \left( u \frac{\partial T}{\partial x} + v \frac{\partial T}{\partial y} + w \frac{\partial T}{\partial z} \right) = k_{eff} \left( \frac{\partial^2 T}{\partial x^2} + \frac{\partial^2 T}{\partial y^2} + \frac{\partial^2 T}{\partial z^2} \right) + \varepsilon S_t \tag{19}
$$

In the energy equation, the effective thermal conductivity is given by

$$
k_{\text{eff}} = \varepsilon k_f + (1 - \varepsilon) k_s \tag{20}
$$

where  $k_f$  is the fluid phase thermal conductivity,  $k_s$  the solid medium thermal conductivity and  $\varepsilon$  the porosity of the medium.

The source term  $S_t$  in the energy equation due to the chemical reactions is determined by

$$
S_t = -(\Delta H_{SR} R_{SR} + \Delta H_{rWGS} R_{rWGS})
$$
\n(21)

#### Table 1





where  $\Delta H_{SR}$  is the enthalpy of reaction of the steam reforming reaction, and  $\Delta H_{\text{rwGS}}$  is the enthalpy of reaction of the reverse water– gas-shift reaction.

In the solid regions, the energy transport equation can be written as

$$
\frac{\partial^2 T}{\partial x^2} + \frac{\partial^2 T}{\partial y^2} + \frac{\partial^2 T}{\partial z^2} = 0
$$
\n(22)

The boundary conditions of the present computation include those at the inlet, outlet, wall, and the interfaces between the flow channel and the catalyst layer.

- (1) The boundary conditions for inlets at the flow channel and the catalyst layer. The inlet flow velocity is constant, the inlet gas composition is constant, and the inlet temperature is constant.
- (2) The boundary conditions for outlets at the flow channel and the catalyst layer. There is fully developed flow.
- (3) The boundary conditions for the interface between the solid wall and the insulated walls. The temperature gradients are zero.
- (4) The boundary conditions for the interface between the flow channel and the insulated walls. The velocities, temperature, temperature gradient, species concentration and species flux are zero.
- (5) The boundary conditions for the interface between the flow channel and solid wall. No slip and zero fluxes hold the velocities and the concentration gradients at zero.
- (6) The boundary conditions for the interface between the flow channel and the catalyst layer. The velocities, temperature, species concentration and species flux are continuous.
- (7) The boundary conditions for the interface between the heated wall and the catalyst layers. The velocities and the concentration gradient are assumed to be zero, and the temperature is assumed to be equal to the constant wall temperature.

[Fig. 1](#page-2-0) presents schematics of the three-dimensional plate methanol steam micro-reformer considered in this work. The inlet cross-section of the channel is 1 mm  $\times$  1 mm. The thicknesses of the catalyst layers are set at 50  $\mu$ m. The base conditions of the properties are as follows; the operating pressure is 1 atm and the inlet temperature is 120 °C. The inlet flow velocity is 1 m  $s^{-1}$  and the molar ratio of  $H<sub>2</sub>O/CH<sub>3</sub>OH$  is 1.3 in the first example. The physical properties of the channels are listed in Table 1.

## 3. Numerical method

The solution to the governing equations is obtained by employing a finite volume scheme with the model domain divided into a

Table 2 Temperature distributions ( $\degree$ C) for the various grid tests at different axial locations.

$I \times I \times K$							
	$\Omega$	0.152	0.303	0.455	0.606	0.758	0.909
$67 \times 16 \times 3$ $133 \times 26 \times 5$ $265 \times 36 \times 9$	120 120 120	176 174 169	202 201 200	218 217 216	224 224 223	227 227 226	228 228 228

number of cells and used as control volumes. The governing equations are numerically integrated over each of these computational cells or control volumes. The method exploits a collocated cell-centered variable arrangement with the local or cell-averaged values of the physical quantities evaluated and stored at each cell center.

A generalized form of the transport equation for mass, momentum, energy can be expressed in a conservative form as follows:

$$
\nabla \cdot (\rho \phi V) = \nabla \cdot (\Gamma \nabla \phi) + S_{\phi}
$$
 (23)

where  $\phi$  is a general dependent variable. V is the velocity vector,  $S_{\phi}$ is the source per unit volume and  $\rho$  is the density. With the discretization of the governing equations, the coupled finite-difference equations become

$$
a_P \phi_P = a_E \phi_E + a_W \phi_W + a_N \phi_N + a_S \phi_S + S_\phi \tag{24}
$$

where  $\phi_P$  is the value of  $\phi$  at the current point P,  $\phi_E$ , ...,  $\phi_S$  stand for the values of the grid points adjacent to the point P, and  $a_P$ , ...,  $a_S$ are known as the link coefficients. All equations were numerically solved using the commercial fluid dynamics program Fluent 6.1. The SIMPLE algorithm was employed to solve the convection–diffusion equations. The convergence criteria for the normalized residuals for each variable were restricted to be less than  $10^{-6}$ .

In order to study the effect of grid number on the numerical results, the grid independence was examined in preliminary test runs. For simplification of the analysis, three grid configurations were evaluated for the single channel of the plate methanol steam micro-reformer at a wall temperature of 230  $\degree$ C. The single channel length is 33 mm, and the cross-section of the channel is  $1$  mm  $\times$  1 mm. The thicknesses of the catalyst layer are set at 50  $\mu$ m. The inlet flow velocity is 1 m s<sup>-1</sup> and the molar ratio of H2O/CH3OH is 1.3. The physical properties of the channel are listed in [Table 1](#page-3-0). The numbers of grid lines in the  $x$ ,  $y$  and  $z$  directions were  $265 \times 36 \times 9$ ,  $133 \times 26 \times 5$ , and  $67 \times 16 \times 3$ . The influence of grid lines on the local temperatures is shown in Table 2. The deviations of local temperatures are 0.04–1.1% for grids  $67 \times 16 \times 3$  and  $133 \times 26 \times 5$ , and 0.4–2.8% for grids  $133 \times 26 \times 5$  and  $265 \times 36 \times 9$ . Therefore, Grid  $132 \times 25 \times 4$  was chosen for the simulation in the present study as a tradeoff between accuracy and CPU computation time.

In order to compare the numerical results and experimental data, a micro-reformer with parallel flow field is tested. Fig. 2 shows a comparison of the present prediction with previous experimental data. The solid symbols denote the experimental results of Park et al. [\[38\]](#page-10-0) and the curve is the present prediction. Only small discrepancies between the numerical results and the experimental data have been found. The numerical model accurately predicted the methanol conversion and the gas distributions. Hence, the proposed three-dimensional numerical model is adequate for analyzing the heat and mass transfer in a micro-reformer.

# 4. Results and discussion

In past studies of heat and mass transfer in a micro-reformer, simplified models without wall conduction effects have been developed. In this work, a detailed analysis of the three-dimensional modeling with wall conduction effects has been proposed



Fig. 2. Comparison of present prediction with the previous experimental data of Park et al. [\[38\]](#page-10-0). (a) Methanol conversion and (b) product gas composition.



Fig. 3. Comparisons of the predicted results with and without wall conduction effects for the temperature distributions and CH3OH mole fraction distributions along the centerline of the serpentine flow field  $(Y = 0.5)$ .

to examine the transport phenomena of heat and mass transfer in a micro-reformer with serpentine flow channels. Fig. 3 presents the effects of wall conduction on the local temperature distributions and CH<sub>3</sub>OH mole fraction distributions along the centerline of the serpentine flow field. In this plot, X denotes the dimensionless distance from the serpentine flow channel inlet to outlet. It is clearly seen that the temperature distribution is affected by the wall thermal conduction and shows a rapid and uniform development compared with that without a wall conduction effect. This implies that the effects of wall conduction on the thermal development in a micro-reformer are important. In addition, it is found in Fig. 3 that the effects of wall conduction on the methanol



Fig. 4. Effects of various wall temperatures on: (a) the temperature distributions and (b) the CH<sub>3</sub>OH mole fraction distributions along the centerline of the serpentine flow field  $(Y = 0.5)$ .

distribution are negligible, whereas their effect on the temperatures distribution is remarkable and cannot be neglected in the modeling. Therefore, wall conduction effects are taken into account in this work.

The wall temperature,  $T_w$ , on the heated wall is important. To this end, the effects of wall temperatures on the dimensionless temperature distributions along the centerline of the serpentine flow field at a Reynolds number of 41 and a  $H_2O$ /CH<sub>3</sub>OH molar ratio of 1.3 are examined in Fig. 4a. In Fig. 4a, the fluid temperature shows a significant increase when the fuel moves downstream. The fluid temperature increases along the serpentine flow channel due to the heated wall. The dimensionless temperature experiences a negligible variation when  $X > 0.3$ . This is due to thermal equilibrium for  $X > 0.3$ , in which the temperature has a more uniform distribution. In addition, it is found in Fig. 4a that the dimensionless temperature distributions are almost the same for various values of wall temperature. Fig. 4b shows the effects of the wall temperature  $T_w$  on the distribution of CH<sub>3</sub>OH and H<sub>2</sub> mole fraction along the centerline of the serpentine flow field. The  $CH<sub>3</sub>OH$  mole fraction gradually decreases along the flow directions due to the chemical reaction, whereas the  $H_2$  mole fractions increase when the flow moves downstream. It is also clearly seen that the  $CH<sub>3</sub>OH$  consumption and the  $H<sub>2</sub>$  mole fractions increase with an increase in the wall temperature. A lower  $CH<sub>3</sub>OH$  mole fraction and a higher  $H<sub>2</sub>$  mole fraction along the serpentine flow field represent a higher methanol conversion.

The velocity distributions along the centerline of the serpentine flow field for various values of wall temperature are presented in Fig. 5a. Significant variations are seen clearly at the turning points of the flow channel due to the change of velocity direction. The fuel flow velocity increases from the inlet to outlet, due to increases in



Fig. 5. Effects of various wall temperatures on: (a) the local velocity and (b) the local pressure along the centerline of the serpentine flow field  $(Y = 0.5)$ .

the temperature caused by significant density variation. The rise in fuel velocity becomes relatively insignificant since it is fully developed thermally at  $X > 0.3$ . The thermo-fluid parameters affect not only the methanol conversion efficiency, but also pressure loss (the difference between local and inlet pressures) in the serpentine flow channel. Large pressure drops in the channel mean that more pumping work is needed to pump the reactants. Thus, pressure loss is a significant issue. An exploration of pressure loss for various wall temperatures along the centerline of the serpentine flow field is presented in Fig. 5b. Local pressure loss increases along the serpentine flow field. A larger pressure loss occurs at higher wall temperatures. In addition, clearly observed variations in the pressure loss appear at the turning points, due to the velocity change. The pressure losses are higher for the serpentine flow field than for the parallel flow channel [\[32\]](#page-10-0), so reducing them is a priority for future use with the plate methanol micro-reformer.

The effects of the Reynolds number (Re) on the local dimensionless temperature,  $CH<sub>3</sub>OH$  and  $H<sub>2</sub>$  mole fraction distributions at  $T_w = 230$  °C are analyzed in [Fig. 6](#page-6-0). A careful examination of [Fig. 6](#page-6-0)a reveals that the dimensionless temperature increases as the fluid moves downstream due to the heated wall. In addition, the dimensionless temperature rise is lower for a higher Re. This is due to the fact that a higher Reynolds number significantly increases the heat leaving the flow channel, which decreases the temperature rise. In [Fig. 6](#page-6-0)b, the  $CH<sub>3</sub>OH$  mole fraction clearly decreases along the serpentine flow field. This is due to the temperature rise caused by the strong chemical reaction, which in turn increases the methanol consumption. Also, the methanol mole fraction distributions decrease with a lower Reynolds number. The results also show that a higher  $H_2$  mole fraction is noted for the lower Reynolds number. This is due to the longer gas resident

<span id="page-6-0"></span>

Fig. 6. Effects of Reynolds numbers at  $T_w = 230$  °C on: (a) the temperature distributions and (b) the  $CH<sub>3</sub>OH$  mole fraction distributions along the center line of the serpentine flow field  $(Y = 0.5)$ .

time which results in a better methanol conversion and a higher  $H_2$ production.

Fig. 7 presents the effects of the  $H_2O/CH_3OH$  (S/C) molar ratio on the CH<sub>3</sub>OH, H<sub>2</sub> and CO mole fraction distributions at  $T_w$  = 230 °C. Fig. 7a shows that more efficient methanol conversion is noted at a higher  $H<sub>2</sub>O/CH<sub>3</sub>OH$  molar ratio. It is also found that a higher molar ratio of  $H<sub>2</sub>O/CH<sub>3</sub>OH$  causes the  $H<sub>2</sub>$  mole fraction to fall. The CO mole fraction distributions are also shown in Fig. 7b. The CO concentration decreases with an increase in  $H<sub>2</sub>O/CH<sub>3</sub>OH$  molar ratio. This is because the higher  $H<sub>2</sub>O$  concentration would enhance the water–gas-shift reaction, which in turn reduces the CO concentration. However, the higher  $H_2O/CH_3OH$  molar ratio also reduces the  $H<sub>2</sub>$  concentration at the channel outlet.

[Figs. 8 and 9](#page-7-0) present the  $CH<sub>3</sub>OH$ ,  $H<sub>2</sub>$  and CO mole fraction distributions for  $T_w = 230$  °C along the middle cross-section of the serpentine flow field  $(Y = 0.5)$  and the interface between flow channel and catalyst layer  $(Y = 0.95)$ , respectively. An overall inspection of [Figs. 8 and 9](#page-7-0) reveals that the mole fraction of  $CH<sub>3</sub>OH$ decreases along both the middle cross-section of the channel and the interface between the flow channel and catalyst layer. A comparison between [Figs. 8 and 9](#page-7-0) show that the methanol consumption is higher at the interface compared with the middle crosssection. This is because the interface has a higher temperature and chemical reaction rate. It is evident that the mass fraction of  $H<sub>2</sub>$  and CO increases along the middle cross-section of the serpentine flow field and along the interface between the flow channel and catalyst layer. This confirms that a methanol micro-reformer with a higher  $CH<sub>3</sub>OH$  consumption indicates a higher  $H<sub>2</sub>$  and CO concentration, as would be expected.

[Fig. 10](#page-9-0) presents the effects of the Reynolds number Re on the methanol conversion  $\eta$  and H<sub>2</sub> production rate Q<sub>H2</sub> for various



Fig. 7. Effects of H<sub>2</sub>O/CH<sub>3</sub>OH molar ratio (S/C) at  $T_w = 230$  °C on: (a) the temperature distributions and (b) the CH3OH mole fraction and CO mole fraction distributions along the center line of the serpentine flow field ( $Y = 0.5$ ).

values of wall temperature  $T_w$ . In [Fig. 10,](#page-9-0) the value of  $Q_{H2}$  increases and  $\eta$  decreases with an increase in Re. This can be attributed to the fact that a lower Re implies an increase in the fuel residence time and temperature distributions in the micro-reformer, which in turn improves the methanol conversion. As for the  $Q_{H2}$ , a higher  $Q_{H2}$  is found for the case with a higher Re due to the higher inlet flow rate. In addition,  $\eta$  and  $Q_{H2}$  increase with an increase in wall temperature  $T_w$ . Besides, it is important to note that a higher Reynolds number will not necessarily provide a better  $H_2$  production rate. When the methanol conversion is too small, a higher Reynolds number may provide a lower  $H_2$  production rate.

The effects of various heated plates of the channel on the temperature distributions and  $CH<sub>3</sub>OH$  mole fraction distributions are examined in [Fig. 11.](#page-9-0) In [Fig. 11](#page-9-0), two cases of heated plates were tested, top heated plate  $(Y = 1)$  or bottom heated plate  $(Y = 0)$ . The effects of top versus bottom heated plates of the channel on temperature distributions along the center line of the serpentine flow field are shown in [Fig. 11](#page-9-0). A top heated plate has a higher temperature distribution than a bottom heated plate. This is because a bottom heated plate has a smaller temperature rise due to fuel convection effects, while a top heated plate with the heated position near the catalyst layer has a more significant temperature rise. For a bottom heated plate, it is also apparent that marked temperature variations occur at the turning points of the flow channel due to the change of the heat flow direction. The top heated plate clearly has a larger methanol consumption than the bottom heated plate. These phenomena make it obvious that the stronger chemical reaction for the case with the top heated plate is due to a higher temperature distribution.

<span id="page-7-0"></span>

Fig. 8. Local distributions of: (a) CH<sub>3</sub>OH mole fraction, (b) H<sub>2</sub> mole fraction and (c) CO mole fraction along the cross-section of Y = 0.5 at  $T_w$  = 230 °C.



Fig. 9. Local distributions of: (a) CH<sub>3</sub>OH mole fraction, (b) H<sub>2</sub> mole fraction and (c) CO mole fraction along the interface between the flow channel and catalyst layer (Y = 0.95) at  $T_w = 230^\circ$ .

<span id="page-9-0"></span>

Fig. 10. Effects of wall temperature and inlet fuel Reynolds number on the methanol conversion and  $H_2$  production rate.



Fig. 11. Effects of heating different channel plates on the temperature and CH<sub>3</sub>OH mole fraction distributions along the center line of the serpentine flow field  $(Y = 0.5)$ .

### 5. Conclusions

A three-dimensional heat and mass transfer of micro-reformer with the serpentine flow field numerical model is developed to examine the effects of various values of wall temperature, fuel ratio and Reynolds number on the plate methanol steam micro-reformer performance and local transport phenomena. The conclusions drawn from the analysis are:

- 1. The predictions with and without wall conduction effects show that this has a significant impact on the thermal development and a negligible effect on the chemical reactions.
- 2. The prediction shows that the methanol conversion increases with an increase in wall temperature  $T_w$ . The methanol conversion for  $T_w$  = 260 °C could be improved by 53% relative to that of  $T_w = 200$  °C.
- 3. A reduced Reynolds number for the reactant gas in the channel would raise the reactant gas residence time and yield more uniform temperature distributions, which in turn increases the reaction time and improves the methanol conversion efficiency.
- 4. The CO concentration would be reduced from 0.27% to 0.21% at  $T_w$  = 230 °C as the H<sub>2</sub>O/CH<sub>3</sub>OH molar ratio value is increased from 1.0 to 1.6.
- 5. By comparing the predicted results with top or bottom heated condition, it indicates that the methanol conversion for the condition with the top heated plate could be improved by 4.4%.

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