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On the Singular Behavior of the Solution of $v''(x) + x \sin v(x) = 0$

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In this paper we study the singular behavior as $a \uparrow \pi$, about the solution v(x; a) of the initial value problem $v''(x) + x \sin v(x) = 0$, v'(0) = 0, v(0) = a. We also illustrate its application to the large deformation of a heavy cantilever by its own weight. © 1992 Academic Press, Inc.

1. INTRODUCTION

In this paper we are concerned with the singular behavior of the solutions of the following initial value problem:

$$v''(x) + x \sin v(x) = 0,$$

 $v'(0) = 0,$ (I)_a
 $v(0) = a, \quad a \in \mathbb{R}.$

We denote the solution of $(I)_a$ by v(x; a). The qualitative behavior of the solutions v(x; a) is important to the studies of the following mathematical model (1.1) which describes the large deformations of a heavy cantilever by its own weight (see [1] or [2]):

$$v''(x) + x \sin v(x) = 0,$$

 $v'(0) = 0, \quad v(K) = \pi - \alpha, \ 0 \le \alpha \le \pi.$
(1.1)

In [2] the authors studied the two-point boundary value problem (1.1) by using the shooting method. From the uniqueness of the solution of the initial value problem $(I)_a$, it follows that

$$v(x; \pi) \equiv \pi, \qquad v(x; -\pi) \equiv -\pi, \qquad v(x; 0) \equiv 0;$$

$$v(x; 2\pi + a) = 2\pi + v(x; a), \qquad (1.2)$$

$$v(x; 2\pi - a) = 2\pi - v(x; a),$$

and it suffices to consider the problem (1)_a only for the case $0 < a < \pi$. We note that from [2] for all $0 < a < \pi$, v(x; a) is oscillatory over $[0, \infty)$ and $-\pi < v(x; a) < \pi$ for all $x \ge 0$. We introduce

$$\Delta(x;a) = \frac{dv}{da}(x;a), \qquad \phi(x) = \Delta(x;0).$$

Then differentiating $(I)_a$ with respect to a yields

$$\Delta''(x) + x(\cos v(x; a)) \Delta(x) = 0,$$

$$\Delta'(0) = 0,$$

$$\Delta(0) = 1.$$

(1.3)

Setting a = 0 in (1.3) yields

$$\phi''(x) + x\phi(x) = 0, \qquad \phi'(0) = 0, \qquad \phi(0) = 1.$$
 (1.4)

Let $y_n(a)$, $z_n(a)$ be the *n*th zeros of v(x; a) and v'(x; a), respectively, for n = 1, 2, ..., with $0 = z_1 < y_1 < z_2 < \cdots < y_n < z_{n+1} < y_{n+1} < \cdots$ and λ_n, γ_n be *n*th zero of $\phi(x)$ and $\phi'(x)$, respectively, for n = 1, 2, ... Then in [2] we have shown the following result.

THEOREM 1.1. Let $0 < a < \pi$; then $\Delta(x; a)$ has an infinite number of isolated zeros $\alpha_n(a)$ and $\Delta'(x; a)$ satisfies the following:

(i) If $0 < a < \pi/2$, then $\Delta'(x; a)$ has an infinite number of isolated zeros $\beta_n(a)$, $0 = \beta_1 < \beta_2 < \cdots < \beta_n < \cdots$. Furthermore, $\beta_1 = z_1 = 0 < y_1 < \alpha_1 < z_2 < \beta_2 < y_2 < \alpha_2 < \cdots < y_n < \alpha_n < z_{n+1} < \beta_{n+1} < y_{n+1} < \cdots$.

(ii) If $\pi/2 \le a < \pi$ then $\Delta'(x; a)$ has an infinite number of isolated zeros $\beta_n(a)$, $0 = \beta_0 < \beta_1 < \beta_2 < \cdots < \beta_n < \cdots$. Furthermore, $\beta_0 = z_1 = 0 < \beta_1 < y_1 < \alpha_1 < z_2 < \beta_2 < y_2 < \cdots < y_n < \alpha_n < z_{n+1} < \beta_{n+1} < y_{n+1} < \cdots$.

(iii) $\lim_{a\to 0^+} y_n(a) = \lambda_n$, $\lim_{a\to 0^+} z_n(a) = \gamma_n$, and $\lim_{a\to \pi^-} y_n(a) = \infty$, for n = 1, 2, ..., moreover,

$$\frac{dy_n}{da} > 0, \qquad \frac{dz_n}{da} > 0, \qquad for \quad n = 1, 2, \dots$$

We introduce the following Liapunov function:

$$V(x) = (1 - \cos v(x; a)) + \frac{1}{2} \frac{(v'(x; a))^2}{x}.$$
 (1.5)

It is easy to verify that

$$V'(x) = -\frac{1}{2} \left[\frac{v'(x;a)}{x} \right]^2 \le 0, \quad \text{for all} \quad x \ge 0.$$
 (1.6)

Then we have

$$1 - \cos v(0) > 1 - \cos v(z_1) > 1 - \cos v(z_2) > \cdots,$$
(1.7)

and it follows that $|v(x; a)| \leq a$ for all $x \geq 0$. That is, $\{|v(z_n(a); a)|\}$ is a monotone decreasing sequence; moreover from [3] we have

THEOREM 1.2. Given $a \in (0, \pi)$, we have

- (i) $v(z_{2n}(a); a)$ monotonically increases to zero as $n \to \infty$;
- (ii) $v(z_{2n+1}(a); a)$ monotonically decreases to zero as $n \to \infty$.

Consequently Theorem 1.2 says that for any given a, $0 < a < \pi$, the solution v(x; a) satisfies $\lim_{x \to \infty} v(x; a) = 0$.

2. MAIN RESULTS

In this section we illustrate the singular behavior of the solution v(x; a) as $a \to \pi^-$.

LEMMA 2.1. Let $h(a) = v^2(z_n(a); a)$. Then h(a) is strictly increasing on $(0, \pi)$.

Proof. We have

$$h'(a) = 2v(z_n(a); a) \left[v'(z_n(a); a) \frac{dz_n}{da} + \Delta(z_n(a); a) \right]$$

= 2v(z_n(a); a) \Delta(z_n(a); a).

From Theorem 1.1, it is easy to verify h'(a) > 0 for any $0 < a < \pi$.

In the following, we state and prove our main result.

THEOREM 2.2. For each n = 1, 2, ..., we have

$$\lim_{a \to \pi} v(z_n(a); a) = (-1)^{n+1} \pi, \quad \text{for all} \quad n = 1, 2, 3, \dots$$

Proof. Let $\varepsilon = \pi - a$, denote $v(x; \varepsilon)$ to be the solution of the initial value problem

$$v''(x) + x \sin v = 0$$

$$v'(0) = 0$$

$$v(0) = \pi - \varepsilon \qquad 0 < \varepsilon < \pi,$$

(2.1)

and let $y_n(\varepsilon)$, $z_n(\varepsilon)$ be the *n*th zero of $v(x;\varepsilon)$ and $v'(x;\varepsilon)$, respectively, for n = 1, 2, 3, ... It is equivalent to show $\lim_{\varepsilon \to 0} v(z_n(\varepsilon);\varepsilon) = (-1)^{n+1}\pi$, and we prove it by mathematical induction. We have $v(z_1(a);a) = a$, and Theorem 2.2 holds trivially for n = 1.

Step 1. For n=2 we prove $\lim_{\epsilon \to 0} v(z_2(\epsilon); \epsilon) = -\pi$ or $\lim_{a \to \pi} v(z_2(a); a) = -\pi$. Multiplying by v'(x) on both sides of (2.1) and integrating the resulting identity from a to b yields

$$\frac{1}{2} (v'(b;\varepsilon))^2 - \frac{1}{2} (v'(a;\varepsilon))^2 = b \cos v(b;\varepsilon) - a \cos v(a;\varepsilon) - \int_a^b \cos v(x;\varepsilon) \, dx.$$
(2.2)

For any $\delta > 0$ and sufficiently small $\varepsilon > 0$ with $\delta > \varepsilon > 0$, we define $y(\varepsilon, \delta)$ and $y^*(\varepsilon, \delta)$ to be the first real numbers satisfying $v(y(\varepsilon, \delta); \varepsilon) = \pi - \delta$ and $v(y^*(\varepsilon, \delta); \varepsilon) = \pi - \delta/2$, respectively. Obvious $y^*(\varepsilon, \delta) < y(\varepsilon, \delta)$. Since $v(x; \varepsilon)|_{\varepsilon=0} \equiv \pi$ for all $x \ge 0$, and from the continuous dependence on initial data, we have

$$\lim_{\varepsilon \to 0} y^*(\varepsilon; \delta) = +\infty.$$
 (2.3)

Setting a = 0, $b = y(\varepsilon; \delta)$ in (2.2) yields

$$\frac{1}{2} [v'(y(\varepsilon;\delta);\varepsilon)]^2 = y(\varepsilon,\delta) \cos v(y(\varepsilon;\delta);\varepsilon) - \int_0^{y(\varepsilon;\delta)} \cos v(x;\varepsilon) dx$$
$$= \int_0^{y(\varepsilon;\delta)} [\cos v(y(\varepsilon;\delta);\varepsilon) - \cos v(x;\varepsilon)] dx$$
$$> \int_0^{y^*(\varepsilon;\delta)} [\cos(\pi-\delta) - \cos v(x;\varepsilon)] dx$$
$$\ge y^*(\varepsilon;\delta) [\cos(\pi-\delta) - \cos(\pi-\delta/2)].$$

We have $[\cos(\pi - \delta) - \cos(\pi - \delta/2)] > 0$ and from (2.3)

$$\lim_{\varepsilon \to 0} v'(y(\varepsilon; \delta); \varepsilon) = -\infty.$$
(2.4)

Consequently that $v''(x) = -x \sin v < 0$ on $(0, y_1(\varepsilon))$ implies that

$$\lim_{\varepsilon \to 0} v'(y_1(\varepsilon); \varepsilon) = -\infty.$$
(2.5)

from the identity

$$\int_{y(\varepsilon;\delta)}^{y_1(\varepsilon)} v'(x;\varepsilon) \, dx = -(\pi-\delta),$$

and the concavity of v on $(y(\varepsilon; \delta), y_1(\varepsilon))$, we have

$$\frac{\delta - \pi}{v'(y_1(\varepsilon);\varepsilon)} < y_1(\varepsilon) - y(\varepsilon;\delta) < \frac{\delta - \pi}{v'(y(\varepsilon;\delta);\varepsilon)}.$$
(2.6)

Then from (2.4), (2.5), and (2.6) it follows that

$$\lim_{\varepsilon \to 0} y_1(\varepsilon) - y(\varepsilon; \delta) = 0.$$
 (2.7)

In the following, we establish that

$$\lim_{\varepsilon \to 0} \frac{[v'(y_1(\varepsilon);\varepsilon)]^2}{y_1(\varepsilon)} = 4.$$
(2.8)

Setting a = 0, $b = y_1(\varepsilon)$ in (2.2) yields

$$\frac{1}{2} (v'(y_1(\varepsilon);\varepsilon))^2 = y_1(\varepsilon) - \int_0^{y_1(\varepsilon)} \cos v(x;\varepsilon) dx$$
$$= y_1(\varepsilon) - \int_0^{y(\varepsilon;\delta)} \cos v(x;\varepsilon) dx - \int_{y(\varepsilon;\delta)}^{y_1(\varepsilon)} \cos v(x;\varepsilon) dx.$$

It is easy to verify that

$$y_{1}(\varepsilon) - y(\varepsilon; \delta) \cos(\pi - \delta) - \int_{y(\varepsilon; \delta)}^{y_{1}(\varepsilon)} \cos v(x; \varepsilon) dx$$

$$\leq \frac{1}{2} (v'(y_{1}(\varepsilon); \varepsilon))^{2}$$

$$\leq y_{1}(\varepsilon) - y(\varepsilon; \delta) \cos(\pi - \varepsilon) - \int_{y(\varepsilon; \delta)}^{y_{1}(\varepsilon)} \cos v(x; \varepsilon) dx.$$

From (2.7) and let $\varepsilon \to 0$ in the above inequality, we have

$$1 - \cos(\pi - \delta) \leq \lim_{\varepsilon \to 0} \frac{(v'(y_1(\varepsilon); \varepsilon))^2}{2y_1(\varepsilon)} \leq 1 + 1 = 2.$$

$$(2.9)$$

Since $\delta > 0$ is arbitrary, (2.8) follows directly from (2.9). We note that from (1.5) and (1.6), we have

$$V(0) > V(y(\varepsilon; \delta)) > V(y_1(\varepsilon))$$

or

$$1 - \cos(\pi - \varepsilon) \ge [1 - \cos(\pi - \delta)] + \frac{1}{2} \frac{(v'(y(\varepsilon; \delta); \varepsilon))^2}{y(\varepsilon; \delta)}$$
$$\ge \frac{(v'(y_1(\varepsilon); \varepsilon))^2}{2y_1(\varepsilon)}.$$
(2.10)

From (2.8), (2.10) we obtain

$$\lim_{\varepsilon \to 0} \frac{[v'(y(\varepsilon; \delta); \varepsilon)]^2}{y(\varepsilon; \delta)} = 2(1 + \cos(\pi - \delta)).$$
(2.11)

Then from (2.7), (2.8), and (2.9) we have

$$\lim_{\varepsilon \to 0} \frac{(v'(y_1(\varepsilon); \varepsilon))^2}{(v'(y(\varepsilon; \delta); \varepsilon))^2} = \lim_{\varepsilon \to 0} \frac{(v'(y_1(\varepsilon); \varepsilon))^2/y_1(\varepsilon)}{(v'(y(\varepsilon; \delta); \varepsilon))^2/y(\varepsilon; \delta)} \frac{y_1(\varepsilon)}{y(\varepsilon; \delta)}$$
$$= \frac{2}{1 - \cos(\pi - \delta)}.$$

Hence there exists a constant $\hat{C}_1 = \hat{C}_1(\delta) > 0$, such that for $\varepsilon > 0$ sufficiently small

$$|v'(y_1(\varepsilon);\varepsilon)| < \hat{C}_1 |v'(y_1(\varepsilon;\delta);\varepsilon)|.$$
(2.12)

From (2.6) and (2.12), we have

$$0 < y_1(\varepsilon) - y(\varepsilon; \delta) < \frac{\hat{C}_1(\pi - \delta)}{|v'(y_1(\varepsilon); \varepsilon)|} = \frac{C}{|v'(y_1(\varepsilon); \varepsilon)|}, \qquad (2.13)$$

where $C = C(\delta) = \hat{C}_1(\pi - \delta) > 0$ independent of ε .

Now we are in a position to show that $\lim_{\varepsilon \to 0} v(z_2(\varepsilon); \varepsilon) = -\pi$. Suppose this does not hold, then there exists $\delta^* > 0$ such that $v(z_2(\varepsilon); \varepsilon) > -\pi + 2\delta^*$ for all $\varepsilon > 0$. Let $u(x; \varepsilon) = v(x + y_1(\varepsilon); \varepsilon)$ and $w(x; \varepsilon) = -v(y_1(\varepsilon) - x; \varepsilon)$, then $u(x; \varepsilon)$ and $w(x; \varepsilon)$ satisfy the following:

$$u''(x) + y_1(\varepsilon) \sin u = -x \sin u,$$

$$u(0; \varepsilon) = 0, \quad u'(0; \varepsilon) = v'(y_1(\varepsilon); \varepsilon), \quad \text{for all} \quad x \ge 0,$$
(2.14)*

and

$$w''(x) + y_1(\varepsilon) \sin w = x \sin w,$$

$$w(0; \varepsilon) = 0, \qquad w'(0; \varepsilon) = v'(y_1(\varepsilon); \varepsilon), \quad \text{for all } x \ge 0.$$
(2.15)*

Let $\eta = xA(\varepsilon)$, where $A(\varepsilon) = -v'(y_1(\varepsilon); \varepsilon)$, $u(x) = u(\eta/A(\varepsilon)) = \phi(\eta)$ and $w(x) = w(\eta/A(\varepsilon)) = \psi(\eta)$. Then $\phi(\eta)$ and $\psi(\eta)$ satisfy the following:

$$\phi''(\eta) + y_1(\varepsilon) A(\varepsilon)^{-2} \sin \phi = -\eta A(\varepsilon)^{-3} \sin \phi, \quad \eta \ge 0,$$

$$\phi(0; \varepsilon) = 0, \quad \phi'(0; \varepsilon) = -1.$$
(2.14)

and

$$\psi''(\eta) + y_1(\varepsilon) A(\varepsilon)^{-2} \sin \psi = -\eta A(\varepsilon)^{-3} \sin \psi, \quad \eta \ge 0,$$

$$\psi(0; \varepsilon) = 0, \quad \psi'(0; \varepsilon) = -1.$$
(2.15)

If we choose $\delta = \delta^* > 0$, then from (2.13) there exists a constant $C = C(\delta^*)$ independent of ε , such that $0 < y_1(\varepsilon) - y(\varepsilon; \delta^*) < CA(\varepsilon)^{-1}$, provided $\varepsilon > 0$ is sufficiently small. Choose M > C; from (2.5), (2.8), and the continuous dependence on parameter ε , it follows that for all $\varepsilon > 0$ sufficiently small

$$|\phi(\eta) - \psi(\eta)| \leq \delta^*/4$$
, for $0 \leq \eta \leq M$.

In particular, let $\eta = (y_1(\varepsilon) - y(\varepsilon, \delta^*)) A(\varepsilon)$; then

$$w(y_1(\varepsilon) - y(\varepsilon; \delta^*)) + \delta^*/4 > u(y_1(\varepsilon) - y(\varepsilon; \delta^*)) > -\pi + 2\delta^*$$

or

$$-(\pi-\delta^*)+\delta^*/4>-\pi+2\delta^*.$$

This is a desired contradiction and we complete the proof for the case n=2.

Step 2. We now assume inductively that

$$\lim_{\varepsilon \to 0} v(z_k(\varepsilon); \varepsilon) = (-1)^{k+1} \pi$$
(2.16)

and

$$\lim_{\varepsilon \to 0} \frac{(v'(y_{k-1}(\varepsilon);\varepsilon))^2}{y_{k-1}(\varepsilon)} = 4,$$
(2.17)

for all k = 3, 4, 5, ..., n-1. We show that (2.16), (2.17) hold for k = n. For simplicity, we may assume that n is an odd number. We have that (2.16)

holds for k = 2, 3, 4, ..., n - 1. Given any $\delta > 0$, there exists $y_1^r(\varepsilon; \delta)$, $y_k^l(\varepsilon; \delta)$ and $y_k^r(\varepsilon, \delta)$, satisfying

$$0 < y_1^{r}(\varepsilon; \delta) < y_1(\varepsilon), \quad \text{with} \quad v(y_1(\varepsilon; \delta); \varepsilon) = \pi - \delta;$$

$$y_{k-1}(\varepsilon) < y_k^{l}(\varepsilon; \delta) < y_k^{r}(\varepsilon; \delta) < y_k(\varepsilon),$$

with
$$v(y_k^{l}(\varepsilon; \delta); \varepsilon) = v(y_k^{r}(\varepsilon; \delta); \varepsilon) = (-1)^{k+1}(\pi - \delta)$$

for k = 2, 3, ..., n - 1, provided ε is sufficiently small, (see Fig. 1). We claim that

$$\lim_{\epsilon \to 0} y_k^1(\epsilon; \delta) - y_{k-1}(\epsilon) = 0, \quad \text{for} \quad k = 2, 3, ..., n-1, \quad (2.18)$$

$$\lim_{\epsilon \to 0} y_k(\epsilon) - y_k^r(\epsilon; \delta) = 0, \qquad k = 1, 2, 3, ..., n - 1.$$
 (2.19)

Let $V(x) = (1 - \cos v(x; \varepsilon)) + (v'(x; \varepsilon))^2/2x$. Then from the fact that $V'(x) \leq 0$, we have $V(y_{k-1}(\varepsilon)) > V(y_k^1(\varepsilon; \delta)) > V(z_k(\varepsilon))$, for k = 2, 3, ..., n-1. Consequently from (2.16), (2.17), we have $\lim_{\varepsilon \to 0} V(y_k^1(\varepsilon; \delta)) = 2$, or

$$\lim_{\varepsilon \to 0} \frac{(v'(y_k^1(\varepsilon; \delta); \varepsilon))^2}{y_k^1(\varepsilon; \delta)} = 2(1 + \cos(\delta - \pi)).$$
(2.20)



FIG. 1. The graph for the solution $v(x, \varepsilon)$ of (2.1).

From (2.20) we obtain

$$\lim_{\varepsilon \to 0} |v'(y_k^1(\varepsilon; \delta); \varepsilon)| = \infty.$$
(2.21)

Consider the following identity:

$$\int_{y_{k-1}(\varepsilon)}^{y_k^1(\varepsilon;\,\delta)} v'(x;\,\varepsilon)\,dx = (-1)^{k+1}(\pi-\delta).$$

Then we have

$$v'(y_{k-1}(\varepsilon))(y_k^1(\varepsilon;\delta) - y_{k-1}(\varepsilon)) < \delta - \pi < v'(y_k^1(\varepsilon;\delta))(y_k^1(\varepsilon;\delta) - y_{k-1}(\varepsilon))$$

when k is even

or

$$v'(y_k^1(\varepsilon;\delta))(y_k^1(\varepsilon;\delta) - y_{k-1}(\varepsilon)) < \pi - \delta < v'(y_{k-1}(\varepsilon))(y_k^1(\varepsilon;\delta) - y_{k-1}(\varepsilon))$$

when k is odd.

In both cases we have

$$0 < y_k^1(\varepsilon; \delta) - y_{k-1}(\varepsilon) < \frac{\pi - \delta}{|v'(y_k^1(\varepsilon; \delta))|},$$
(2.22)

Hence (2.18) follows directly from (2.21), (2.22). Similarly if we prove

$$\lim_{\varepsilon \to 0} |v'(y_k^{\mathsf{r}}(\varepsilon; \delta); \varepsilon)| = \infty, \qquad (2.23)$$

then (2.19) holds. Setting $a = y_k^1(\varepsilon; \delta)$, $b = y_k^r(\varepsilon; \delta)$ in (2.2) yields

$$\frac{1}{2} (v'(y_k^{\mathsf{r}}(\varepsilon;\delta);\varepsilon))^2 - \frac{1}{2} (v'(y_k^{\mathsf{l}}(\varepsilon;\delta);\varepsilon))^2$$
$$= (y_k^{\mathsf{r}}(\varepsilon;\delta) - y_k^{\mathsf{l}}(\varepsilon;\delta)) \cos(\pi - \delta) - \int_{y_k^{\mathsf{l}}(\varepsilon;\delta)}^{y_k^{\mathsf{l}}(\varepsilon;\delta)} \cos v(x;\varepsilon) \, dx \ge 0.$$

That is,

$$|v'(y_k^{\mathrm{r}}(\varepsilon;\delta);\varepsilon)| \ge |v'(y_k^{\mathrm{l}}(\varepsilon;\delta);\delta)|.$$

Then (2.23) follows directly from (2.21).

We are now in a position to show that (2.17) holds for k = n. We set $a = 0, b = y_{n-1}(\varepsilon)$ in (2.2) to obtain

$$\frac{1}{2} (v'(y_{n-1}(\varepsilon);\varepsilon))^2 = y_{n-1}(\varepsilon) - \int_0^{y_{n-1}(\varepsilon)} \cos v(x;\varepsilon) dx$$
$$= y_{n-1}(\varepsilon) - \int_0^{y_1^r(\varepsilon;\delta)} \cos v(x;\varepsilon) dx$$
$$- \sum_{k=2}^{n-1} \int_{y_k^l(\varepsilon;\delta)}^{y_k^r(\varepsilon;\delta)} \cos v(x;\varepsilon) dx$$
$$- \sum_{k=2}^{n-1} \int_{y_k^{l-1}(\varepsilon;\delta)}^{y_k^l(\varepsilon;\delta)} \cos v(x;\varepsilon) dx - \int_{y_{n-1}^r(\varepsilon;\delta)}^{y_{n-1}(\varepsilon)} \cos v(x;\varepsilon) dx.$$

It is easy to verify the following inequality

$$1 - \frac{y_1^{\mathsf{r}}(\varepsilon;\delta)}{y_{n-1}(\varepsilon)} \cos\left(\pi - \delta\right) - \frac{\cos(\pi - \delta)}{y_{n-1}(\varepsilon)} \sum_{k=2}^{n-1} \left(y_k^{\mathsf{r}}(\varepsilon;\delta) - y_k^{\mathsf{l}}(\varepsilon;\delta)\right)$$
$$- \frac{1}{y_{n-1}(\varepsilon)} \sum_{k=2}^{n-1} \left(y_k^{\mathsf{l}}(\varepsilon;\delta) - y_{k-1}^{\mathsf{r}}(\varepsilon;\delta)\right) - \frac{1}{y_{n-1}(\varepsilon)} \left(y_{n-1}(\varepsilon) - y_{n-1}^{\mathsf{r}}(\varepsilon;\delta)\right)$$
$$\leq \frac{1}{2} \frac{\left(v'(y_{n-1}(\varepsilon);\varepsilon)\right)^2}{y_{n-1}(\varepsilon)}$$
$$\leq 1 + \frac{y_1^{\mathsf{r}}(\varepsilon;\delta)}{y_{n-1}(\varepsilon)} + \frac{1}{y_{n-1}(\varepsilon)} \sum_{k=2}^{n-1} \left(y_k^{\mathsf{r}}(\varepsilon;\delta) - y_k^{\mathsf{l}}(\varepsilon;\delta)\right)$$
$$+ \frac{1}{y_{n-1}(\varepsilon)} \sum_{k=2}^{n-1} \left(y_k^{\mathsf{l}}(\varepsilon;\delta) - y_{k-1}^{\mathsf{r}}(\varepsilon;\delta)\right) + \frac{1}{y_{n-1}(\varepsilon)} \left(y_{n-1}(\varepsilon) - y_{n-1}^{\mathsf{r}}(\varepsilon;\delta)\right).$$

Since $\lim_{\epsilon \to 0} y_{n-1}(\epsilon) = \infty$ and $\delta > 0$ is arbitrary, (2.18) and (2.19) imply

$$\lim_{\varepsilon \to 0} \frac{(v'(y_{n-1}(\varepsilon))^2}{y_{n-1}(\varepsilon)} = 4.$$
 (2.24)

Hence we establish (2.17) for k = n.

Using the same argument as we did in Step 1 yields

$$0 < y_{n-1}(\varepsilon) - y_{n-1}(\varepsilon; \delta) < \frac{C}{|v'(y_{n-1}(\varepsilon))|}$$
(2.25)

for some $C = C(\delta) > 0$ and for all $\varepsilon > 0$ sufficiently small. Since *n* is odd, we show that

$$\lim_{\varepsilon \to 0} v(z_n(\varepsilon); \varepsilon) = \pi.$$
 (2.26)

Suppose (2.26) does not hold. Then there exists a $\delta^* > 0$, such that

$$v(z_n(\varepsilon); \varepsilon) < \pi - 2\delta^*$$
 for all $\varepsilon > 0.$ (2.27)

Let $u(x;\varepsilon) = v(x+y_{n-1}(\varepsilon);\varepsilon)$, $w(x;\varepsilon) = -v(y_{n-1}(\varepsilon)-x;\varepsilon)$. From (2.24), (2.17) and the arguments for the case n = 2, we obtain

$$\pi - 2\delta^* > u(y_{n-1}(\varepsilon) - y_{n-1}^r(\varepsilon; \delta^*); \varepsilon)$$

$$> w(y_{n-1}(\varepsilon) - y_{n-1}^r(\varepsilon; \delta^*); \varepsilon) - \delta^*/4$$

$$= -(-\pi + \delta^*) - \delta^*/4$$

$$= \pi - 5\delta^*/4.$$

This is a desired contradiction. Thus we complete the proof of Theorem 2.2.

3. THE APPLICATION

In [1, 2] the authors discussed a mathematical model describing the deformation of a cantilever by its own weight. It is assumed that a cantilever of uniform cross-section, uniform density ρ , and total length L is held fixed at an angle α at one end, say the origin, and is free at the other end. Let s' be the arc length from the origin, and $\theta = \theta(s')$ be the local angle of inclination. Then we have the governing equation

$$EI \frac{d^2\theta}{ds'^2} = \rho(L - s') \sin \theta,$$

$$\theta(0) = \alpha, \qquad \frac{d\theta}{ds'}(L) = 0,$$
(3.1)

where EI is the flexural rigidity of the material. Let s = s'/L, then the governing equation becomes

$$\frac{d^2\theta}{ds^2} = K^3(1-s)\sin\psi, \qquad 0 \le s \le 1, \ K > 0,$$

$$\theta'(1) = 0, \qquad \theta(0) = \alpha, \qquad -\pi \le \alpha \le \pi,$$

(P)_{\alpha}

where $K = (\rho L^3/\text{EI})^{1/3}$ represents the importance of density and length relative to that of flexural rigidity. Let s = x, $v(x) = \theta(1 - x/K) - \pi$; then we reformulate our equation as the following:

$$v''(x) + x \sin v(x) = 0, \quad ' = d/dx,$$

 $v'(0) = 0, \quad v(K) = \alpha - \pi.$ (3.2)

The vertical case $\alpha = \pi$ was completely analyzed in [2].

We note that from (1.2) we have $v(K; a) = \pi - \alpha$ if and only if $v(K; -a) = \alpha - \pi$. For simplicity, instead of (3.2), we study the multiplicities of the solutions of the following boundary value problem:

$$v''(x) + x \sin v = 0$$

 $v'(0) = 0, \quad v(K) = \pi - \alpha, \quad \text{for} \quad 0 < \alpha < \pi.$
(3.3)

To solve (3.3) by the shooting method, we consider the following initial value problem:

$$v''(x) + x \sin v = 0,$$

 $v'(0) = 0, \quad v(0) = a, \quad \text{for} \quad -\pi < a < \pi.$
(3.4)

THEOREM 3.1. Given $\alpha \in (0, \pi)$,

(i) For each $n = 0, 1, 2, ..., there exists a unique <math>a_{2n+1} = a_{2n+1}(\alpha)$, $a_{2n+1} \in (\pi - \alpha, \pi)$, satisfying $v(z_{2n+1}(a_{2n+1}); a_{2n+1}) = \pi - \alpha$; moreover, $a_1 = \pi - \alpha < a_3 < a_5 < \cdots < a_{2n+1} < \cdots < \pi$, and $\lim_{n \to \infty} a_{2n+1} = 0$.



a - axis

FIG. 2. The graph of the functions $Y_{2n+1}^{u}(a)$ and $Y_{2n+1}^{l}(a)$ for fixed $0 < a < \pi$, n = 0, 1, 2, ...

(ii) Given any $n \ge 1$, for each $a \in (a_{2n-1}, a_{2n+1})$, the equation $v(x; a) = \pi - \alpha$, has exactly 2n - 1 isolated zeros $\{y_1^u, y_{2m+1}^l, y_{2m+1}^u\}_{m=1}^{n-1}$, where $y_1^u = y_1^u(a; \alpha)$, $y_{2m+1}^l = y_{2m+1}^l(a; \alpha)$, $y_{2m+1}^u = y_{2m+1}^u(a; \alpha)$, satisfying $0 = z_1(a) < y_1^u < y_3^l < z_3(a) < y_3^u < \cdots < y_{2n-1}^l < z_{2n-1}(a) < y_{2n-1}^u$; moreover, for $a = a_{2n-1}$, we have $y_{2n-1}^l(a_{2n-1}, \alpha) = y_{2n-1}^u(a_{2n-1}, \alpha) = z_{2n-1}(a_{2n-1})$.

(iii) For each n = 1, 2, ..., as a function of $a, y_{2n+1}^{1}(a)$ attains global minimum at a point $\eta_{2n+1} \in (\pi - \alpha, \pi)$, $a_{2n+1} < \eta_{2n+1}$, satisfying $y_{2n+1}^{1}(\eta_{2n+1}) = \alpha_{2n}(\eta_{2n+1})$, where $\alpha_{2n}(a)$ is the 2nth zero of $\Delta(x; a)$, and $\lim_{a \to \pi^{-}} y_{2n+1}^{1}(a) = +\infty$. On the other hand, $y_{2n+1}^{u}(a)$ is strictly increasing on $[a_{2n+1}, \pi)$ and $\lim_{a \to \pi^{-}} y_{2n+1}^{u}(a) = +\infty$. (See Fig. 2.)

For analogous results, we have

(i)* For each n = 1, 2, ..., there exists a unique $a_{2n} = a_{2n}(\alpha)$, $a_{2n} \in (-\pi, -\pi + \alpha)$, satisfying $v(z_{2n}(a_{2n}); a_{2n}) = \pi - \alpha$. Moreover, $0 > a_2 > a_4 > \cdots > a_{2n} > \cdots > -\pi$, and $\lim_{n \to \infty} a_{2n} = -\pi$.

(ii)* Given any $n \ge 1$, for each $a \in (a_{2n}, a_{2n+2})$, the equation $v(x; a) = \pi - \alpha$ has exactly 2n isolated zeros $\{y_{2m}^1, y_{2m}^u\}_{m=1}^n$, where $y_{2m}^1 = y_{2m}^1(a; \alpha)$,



a — axis

FIG. 3. The graph of the functions $Y_{2n}^{u}(a)$ and $Y_{2n}^{1}(a)$ for fixed $0 < \alpha < \pi$, n = 1, 2, ...

 $y_{2m}^{u} = y_{2m}^{u}(a; \alpha), \quad satisfying \quad y_{2}^{1} < z_{2}(a) < y_{2}^{u} < \cdots < y_{2n}^{1} < z_{2n}(a) < y_{2n}^{u},$ moreover, for $a = a_{2n}$, we have $y_{2n}^{1}(a_{2n}, \alpha) = y_{2n}^{u}(a_{2n}, \alpha) = z_{2n}(a_{2n}).$

(iii)* For each $n = 1, 2, ..., y_{2n}^1(a)$, defined on $(-\pi, a_{2n}]$, attains a global minimum at a point $\eta_{2n} \in (-\pi, \alpha - \pi)$ with $a_{2n} > \eta_{2n}$, satisfying $y_{2n}^1(\eta_{2n}) = \alpha_{2n-1}(\eta_{2n})$, where $\alpha_{2n-1}(a)$ is the (2n-1)th zero of $\Delta(x; a)$ and $\lim_{a \to -\pi^+} y_{2n}^1(a) = +\infty$. On the other hand, $y_{2n}^u(a)$ is strictly increasing on $(-\pi, a_{2n}]$ and $\lim_{a \to -\pi^+} y_{2n}^u(a) = +\infty$. (See Fig. 3.)

Proof. From Lemma 2.1 and Theorem 2.2, it is easy to show that there exists a unique a_{2n+1} , depending on α , $a_{2n+1} \in (\pi - \alpha, \pi)$, satisfying $v(z_{2n+1}(a_{2n+1}); a_{2n+1}) = \pi - \alpha$, and $a_1 = \pi - \alpha < a_3 < a_5 < \cdots < a_{2n+1} < \pi$. We claim $\lim_{n \to \infty} a_{2n+1} = \pi$. If not, then $\lim_{n \to \infty} a_{2n+1} = \pi - \delta_0$ for some $\delta_0 > 0$. Choose $a = \pi - \delta_0/2$. Then from Lemma 3.2 for any $n = 1, 2, ..., v(z_{2n+1}(a); a) > v(z_{2n+1}(a_{2n+1}); a_{2n+1}) = \pi - \alpha$. This is a desired contradiction to Lemma 3.1. Thus we complete the proof for part (i). Part (ii) follows directly from Lemma 3.2 and the oscillatory behavior of the solution v(x; a) and (3.3). We have the relation

$$v(y_{2n+1}^{1}(a;\alpha);a) = \pi - \alpha.$$
(3.5)

Differentiating (3.5) with respect to a yields

$$\frac{dy_{2n+1}^{1}(a;\alpha)}{da} = \frac{-\Delta(y_{2n+1}^{1}(a;\alpha);a)}{v'(y_{2n+1}^{1}(a;\alpha);a)}.$$
(3.6)

Since $y_{2n+1}^1(a_{2n+1}, \alpha) = z_{2n+1}(a_{2n+1})$, we have

$$\frac{dy_{2n+1}^{i}(a;\alpha)}{da}\Big|_{a=a_{2n+1}}=-\infty.$$

However, $y_{2n}(a) \leq y_{2n+1}^1(a, \alpha)$ and $\lim_{a \to \pi} y_{2n}(a) = +\infty$; this shows $\lim_{a \to \pi} y_{2n+1}^1(a) = +\infty$ and the existence of a global minimum η_{2n+1} of $y_{2n+1}^1(a, \alpha)$. From (3.6) we have

$$0 = \frac{dy_{2n+1}^{i}(\eta_{2n+1};\alpha)}{da} = -\frac{\Delta(y_{2n+1}^{i}(\eta_{2n+1};\alpha);\eta_{2n+1})}{v'(y_{2n+1}^{1}(\eta_{2n+1};\alpha);\eta_{2n+1})},$$
(3.7)

and $y_{2n+1}^1(\eta_{2n+1}; \alpha) = \alpha_{2n}(\eta_{2n+1})$ follows directly from (3.7). On the other hand, we have

$$\frac{dy_{2n+1}^{u}(a;\alpha)}{da} = \frac{-\Delta(y_{2n+1}^{u}(a;\alpha);a)}{v'(y_{2n+1}^{u}(a;\alpha);a)}$$

From the relation $\alpha_{2n}(a) < z_{2n+1}(a) < y_{2n+1}^u(a) < y_{2n+1}(a)$, n = 1, 2, ..., it

follows that $dy_{2n+1}^u/da > 0$ for all $a \in (\alpha_{2n+1}, \pi)$. The analogous results for (i)*, (ii)*, and (iii)* can be proved similarly.

Remark 1. For each n = 2, 3, ..., if every extremum of the function $y_n^1(a; \alpha)$ is a local minimum, then η_n is the unique local minimum and $y_n^1(a; \alpha)$ is strictly increasing (decreasing) for $a \ge \eta_n$ ($a \le \eta_n$) provide *n* is odd (even). Differentiating the identity

$$v(y_n^1(a;\alpha);a) = \pi - \alpha,$$

twice with respect to a and setting $dy_n^1/da = 0$ yields

$$\frac{d^2 y_n^l(a,\alpha)}{da^2} = \frac{-(d\Delta/da)(y_n^l(a,\alpha);a)}{v'(y^l(a,\alpha);a)},$$
(3.8)

From Theorem 3.1 (iii) and (iii)*, if $dy_n^1/da = 0$ then $y_n^1(a, \alpha) = \alpha_{n-1}(a)$, and (3.8) becomes

$$\frac{d^2 y_n^1(a,\alpha)}{da^2} = \frac{-(d\Delta/da)(\alpha_{n-1}(a),a)}{v'(\alpha_{n-1}(a),a)}.$$
(3.9)

Since $\Delta(\alpha_{n-1}(a), a) = 0$ for all $a \in (0, \pi)$, it follows that

$$\frac{d\alpha_{n-1}}{da} = \frac{-(d\Delta/da)(\alpha_{n-1}(a), a)}{\Delta'(\alpha_{n-1}(a), a)}.$$
(3.10)

From (3.9), (3.10), and Theorem 2.1(i), (ii), $d^2y_n^1(a, \alpha)/da^2 > 0$ provide $d\alpha_{n-1}/da > 0$, for all $a \in (0, \pi)$, n = 2, 3, ... Let $w(x; a) = (d\Delta/da)(x; a)$, then w(x; a) satisfies

$$w''(x) + xw \cos v = x \Delta^2 \sin v, \qquad w(0) = 0, \qquad w'(0) = 0.$$
 (3.11)

We recall that

$$v''(x) + x \sin v = 0, \quad v(0) = a, \quad v'(0) = 0.$$
 (3.12)

$$\Delta''(x) + x \Delta \cos v = 0, \qquad \Delta(0) = 1, \qquad \Delta'(0) = 0. \tag{3.13}$$

We conjecture that the following hold:

(i) For $0 < a < \pi$, w(x; a) and w'(x; a) are oscillatory over $[0, \pi)$ with zeros $p_n = p_n(a)$, $q_n = q_n(a)$, respectively, for n = 1, 2, ..., where $p_1 = q_1 = 0$.

(ii) For $0 < a < \pi/2$, we have

$$0 = p_1 = q_1 = z_1 = \beta_1 < y_1 < q_2 < \alpha_1 < z_2 < p_2 < \beta_2 < y_2 < q_3 < \cdots$$

$$< \cdots < y_n < q_{n+1} < \alpha_n < z_{n+1} < p_{n+1} < \beta_{n+1} < y_{n+1} < \cdots.$$
(3.14)

For $\pi/2 \leq a < \pi$, we have

$$0 = p_1 = q_1 = z_1 = \beta_0 < \beta_1 < y_1 < q_2 < \alpha_1 < z_2 < p_2 < \beta_2 < y_2 < q_3 < \cdots$$

$$< \cdots < y_n < q_{n+1} < \alpha_n < z_{n+1} < p_{n+1} < \beta_{n+1} < y_{n+1} < \cdots.$$
(3.15)

From (3.10), (3.14), and (3.15) it is easy to verify that $\alpha_n(a)$ is strictly increasing on $(0, \pi)$. In Fig. 4 we plot a graph for the functions $K = y_n^u(a; \alpha)$ and $K = y_n^1(a; \alpha)$ for $0 < \alpha < \pi$. From the figure there follow the bifurcation phenomena of problem $(P)_{\alpha}$, $0 < \alpha < \pi$, or (4.2) as the parameter K varies. It is interesting to note that when $\alpha = 0$ the problem $(P)_0$ has a unique solution [2] for any K while our results show that given any α , $0 < \alpha < \pi$, and any positive integer n, there exists K such that $(P)_{\alpha}$ has n distinct solutions.

Remark 2. We note that for any n = 1, 2, ...,

$$\lim_{\alpha \to 0} a_n(\alpha) = (-1)^{n+1} \pi.$$
 (3.16)

It follows directly from $\pi - \alpha = a_1(\alpha) < a_3(\alpha) < \cdots < a_{2n+1}(\alpha) < \cdots < \pi$, and $-\pi + \alpha > a_2(\alpha) > a_4(\alpha) > \cdots > a_{2n}(\alpha) > \cdots > -\pi$. (3.16) indicates that the bifurcation phenomena will disappear as $\alpha = 0$.



FIG. 4. The graph of the functions $K = Y_n^{i}(a; \alpha)$ and $K = Y_n^{i}(a; \alpha)$ for $0 < \alpha < \pi$, which are the bifurcation pictures for the boundary value problem (3.2).

Remark 3. As $\alpha \rightarrow \pi$ we have

- (i) $\lim_{\alpha \to \pi} a_n(\alpha) = 0$, for all n = 1, 2, ...
- (ii) $\lim_{\alpha \to \pi} y_n^1(a_n(\alpha); \alpha) = \lim_{\alpha \to \pi} y_n^u(a_n(\alpha); \alpha) = \gamma_n.$
- (iii) For arbitrary $\mu > 0$, we have

$$\lim_{\alpha \to \pi} y_n^{l}(a; \alpha) = y_{n-1}(a) \text{ uniformly for all } a \in [\mu, \pi) \text{ if } n \text{ is odd.}$$

$$\lim_{\alpha \to \pi} y_n^{u}(a; \alpha) = y_n(a) \text{ uniformly for all } a \in [\mu, \pi) \text{ if } n \text{ is odd.}$$

$$\lim_{\alpha \to \pi} y_n^{l}(a; \alpha) = y_{n-1}(a) \text{ uniformly for all } a \in (-\pi, -\mu] \text{ if } n \text{ is even.}$$

$$\lim_{\alpha \to \pi} y_n^{u}(a; \alpha) = y_n(a) \text{ uniformly for all } a \in (-\pi, -\mu] \text{ if } n \text{ is even.}$$

To prove (i) we shall only consider the case n is odd; the argument is similar for the case n is even. We have the relation

$$v(z_n(\alpha_n(\alpha)); a_n(\alpha)) = \pi - \alpha. \tag{3.17}$$

Differentiating (3.17) with respect to α yields

$$v'(z_n; a_n(\alpha)) \frac{dz_n}{da} \frac{da_n(\alpha)}{d\alpha} + \Delta(z_n; a_n(\alpha)) \frac{da_n(\alpha)}{d\alpha} = -1.$$

From Theorem 1.1 $\Delta(z_n; a_n(\alpha))$ is positive for odd *n* and $v'(z_n; a_n(\alpha)) = 0$, then we have $da_n(\alpha)/d\alpha < 0$. If $\lim_{\alpha \to \pi} a_n(\alpha) \neq 0$, say $A = \lim_{\alpha \to \pi} a_n(\alpha) > 0$, by Lemma 2.1, we have

$$0 < v^{2}(z_{n}(A); A) < v^{2}(z_{n}(a_{n}(\alpha)); a_{n}(\alpha)) = (\pi - \alpha)^{2}.$$
(3.18)

Let $\alpha \to \pi$ in (3.18), then this leads to a contradiction, $v(z_n(A); A) = 0$. Part (ii) follows directly from Theorem 1.1 (iii) and Theorem 3.1 (ii), (ii)*. For part (iii), we consider only the first case, *n* an odd number. Given $\mu > 0$, from (i) there exists $\delta_1 = \delta_1(\mu) > 0$, such that $a_n(\alpha) < \mu < \pi$, provided $|\pi - \alpha| < \delta_1$. Hence $y_n^1(\alpha; \alpha)$ is well-defined for $\alpha \in [\mu, \pi)$ and $|\pi - \alpha| < \delta_1$. Consider the identity

$$\int_{y_{n-1}(a)}^{y_n^1(a;\alpha)} v'(x;a) \, dx = \pi - \alpha.$$

We have

$$v'(y_n^1(a;\alpha),a)(y_n^1(a;\alpha)-y_{n-1}(a)) < \pi - \alpha$$

or

$$0 < y_n^1(a; \alpha) - y_{n-1}(a) < \frac{\pi - \alpha}{|v'(y_n^1(a; \alpha), a)|}.$$
(3.19)

From (2.21) we have

$$\lim_{a\to\pi}v'(y_n^{\mathbf{l}}(a;\alpha),a)=+\infty.$$

Then

$$M = \max_{a \in [\mu,\pi)} \frac{1}{|v'(y_n^1(a;\alpha),a)|} < \infty.$$

Hence given $\mu > 0$ for any $\varepsilon > 0$, choose $\delta = \min{\{\varepsilon/M, \delta_1\}}$; then

$$|y_n^1(a; \alpha) - y_{n-1}(a)| < \varepsilon,$$
 for all $a \in [\mu, \pi)$,

provided $|\pi - \alpha| < \delta$. Hence the first case of part (iii) holds. By similar arguments it is easy to verify the other cases of part (iii). (See Fig. 4.)

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