國 立 交 通 大 學

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碩 士 論 文

 在負偏壓溫度不穩定回復時臨界電壓改變量分佈 之統計特性和模式及其時間演繹

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Statistical Characterization and Modeling of ΔV_{th} **Distribution in NBTI Recovery and Its Time Evolution**

> 研 究 生 :謝泓達 指導教授 :汪大暉 博士

中華民國 一○一 年 七 月

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研 究 生: 謝泓達 Student : Hung-Da, Hsieh 指導教授 : 汪大暉 博士 Advisor : Dr. Tahui Wang

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學生:謝泓達指導教授:汪大暉 博士

國立交通大學 電子工程學系 電子研究所

摘要

WWW.

 在本篇論文中,為了了解負偏壓溫度回復時的臨界電壓改變量分佈特 性,我們萃取單一電荷散逸導致臨界電壓改變量,還有單一電荷散逸時 間。

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 我們考慮了活化能(Activation Energy)和缺陷能量(Trap Energy)分 佈來模擬負偏壓溫度(NBT)操作後的回復(Recovery)效應之散逸時間 (De-trapping Time)散佈。此外,也研究在經過負偏壓溫度操作後回復時, 單一電荷散逸導致臨界電壓改變量的散佈。我們把這兩個散佈結合到電荷 散逸模型(Trapped Charge Emission Model)。由上述兩個實驗參數來進 行蒙地卡羅模擬經過負偏壓溫度操作後回復的臨界電壓散佈,此結果可精 準符合實驗數據。

Statistical Characterization and Modeling of ΔV_{th} Distribution in NBTI Recovery and Its **Time Evolution**

Student: Hung-Da HsiehAdvisor: Dr. Tahui Wang

Department of Electronics Engineering &

Institute of Electronics National Chiao Tung University **Abstract**

In this thesis, single charge de-trapping induced threshold voltage shift (Δv_{th}) and single charge de-trapping time (τ_e) are extracted to analyze the characteristic of the threshold voltage (ΔV_{th}) dispersion in NBTI recovery.

Activation energy (E_a) and trap energy (E_T) distributions are investigated to simulate de-trapping time (τ_e) dispersion during negative bias temperature instability (NBTI) recovery. Furthermore, single charge de-trapping induced Δv_{th} dispersion is investigated during NBTI recovery. A statistical model combining the trapped charge emission model with the two dispersions is developed. According to the characterization of the single charge phenomenon, we proposed a Monte Carlo simulation to simulate the NBTI recovery threshold voltage dispersion. Our model can fit the measured ΔV_{th} dispersion during NBTI recovery very well.

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2012.7

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Chapter 1 Introduction

For downscaling metal-oxide-semiconductor field-effect-transistor (MOSFET) devices, thickness of gate dielectric is required to be smaller in the progressive technology node. Below the physical thickness of 16 \AA , the gate leakage current exceeds the specifications (1 A/cm) **[1]**. Recently, the high permittivity (high- κ) dielectrics have been proposed to solve this problem due to its higher physical thickness while keeping the same EOT **[2]**. Although significant reductions in the gate leakage have been achieved, two major issues remain for CMOS circuit applications: 1) the high threshold voltages **[3]** and 2) polysilicon (poly-Si) depletion. Consequently, sub-45nm CMOS technology scaling will likely utilize high- κ /metal gate stacks. And the $SiO₂$ interfacial layer offers better device performance and quality of the oxide/Si interface than high- κ dielectrics directly placed on the Si substrate.

THE REAL PROPERTY.

A high- κ /metal gate stack is required for scaled MOSFETs, but one of its most serious problems is in V_{th} control. The BTI degradation of high- κ /metal gate devices can be seen to be much more severe. NBTI is nowadays the most critical device degradation mechanism and became a limiting factor in scaling of modern CMOS technologies (Fig.1.1) **[4]**. PMOS Negative Bias Temperature Instability (NBTI) degrades threshold voltage and drive current, raising an important concern for analog and digital circuits. However, NBTI recovery is an important consideration particularly for typical CMOS operation where applied gate bias switches between "high" and "low" voltages repeatedly **[5]**. Trapped charge would emit when the

applied gate bias is "low" voltage which is known as the recovery mode. The NBTI recovery phenomenon leads to a fast reduction of the stress induced electrical device parameter degradation after end of stress. The NBTI recovery phenomenon endangers proper lifetime estimation. In the past, many experimental and theoretical attempts have been made in order to explain the logarithm time dependence of V_{th} recovery over many decades of time (Fig. 1.2). From these studies, two main NBTI mechanisms have emerged **[6]-[9]**.

In Chapter 2, stair-like drain current traces are observed in the small size devices during NBTI Recovery **[10]-[12]**. In the recovery phase, the tapped charges emit from the high-k layer by thermal assisted tunneling. Single charge de-trapping induced Δv_{th} amplitude and its de-trapping time are extracted from the measurement data.

In Chapter 3, it is found that the de-trapping time $(\tau_{e,i})$ distributions are Gaussian-like distributions. *E^a* distribution can be extracted from the de-trapping time distributions. We assume that the trap energy distribution is uniform. Based on the distributions of E_a and E_T , we simulate the de-trapping time distributions. The de-trapping time distributions are related to the fluctuation of de-trapping hole number. According to the probability distribution of *Δvth*, each single charge de-trapping induced *Δvth* amplitude is simulated. Consequently, staircase-like *Vth* trace during NBTI recovery can be demonstrated a by Monte Carlo simulation. The mean and variance of ΔV_{th} are acquired during NBTI recovery. We also derive the ΔV_{th} distribution by closed form method. Finally, we give a conclusion in Chapter 4.

Fig. 1.1 NBTI/PBTI induced V_{th} drift versus stress time in high- κ /metal gate MOSFETs **[4]**.

Fig. 1.2 Recovery data after three selected stress times obtained at different stress voltages **[16]**.

Chapter 2

Single Charge Effect in NBTI Recovery

2.1 Preface

 It is commonly accepted that mechanism of BTI is due to charge trapping in the high- κ layer **[13]**. BTI is characterized by stressing transistor at high temperature and electric field, periodically interrupting the stress to monitor threshold voltage or drain current ("SMS" technique). The other way to investigate BTI effect is to measure post-stress recovery behavior **[14]**. **THEFT**

To identify the charge de-trapping mechanism, a direct measurement of single charge effect is demonstrated, in which the stair-like threshold voltage trace is measured. The Δv_{th} probability distribution is found by extracting the threshold voltage amplitude caused by single charge emission. We also extract single charge de-trapping time, and it is found that the first three individual de-trapping time distribution are Gaussian-like distributions. The occurrence of total de-trapping time distribution is a uniform distribution. According to the tunneling front model, this implies the trap is a uniform distribution in high- κ layer spatially.

2.2 Device Details and Measurement setup

 The device structure we used in this thesis is illustrated in Fig. 2.1. The devices are pMOSFETs with a metal gate (TiN) and a bi-layer gate dielectric stack consisting of HfON and interfacial SiON layer (IL). The gate length and width are 30nm and 80nm respectively. Its equivalent oxide thickness (EOT) is 7.8 Å .

The Measurement setup is shown in Fig. 2.2. A two channel Agilent B1500 pulse generator connects to drain and gate electrode and it changes the gate bias and drain bias simultaneously. The pulse waveform is shown in Fig. 2.3(a). First, we measure the devices initial state to extract fresh V_{th} ω_{I_d} =500nA. Then, the devices are subjected to a negative gate bias ($Vg=1.8V$) stress during 100s, and recover at fresh *Vth* during 1000s. In the recovery phase, the trapping holes will emit by thermally assisted tunneling (TAT), as shown in Fig. 2.3(b).

2.3 *Δvth* **Probability Distribution in NBTI Recovery**

The recovery ΔV_{th} exhibits a stair-like trace, as shown in Fig. 2.4. Each jump in Fig. 2.4 is believed due to a single trapped hole emission from the high- κ layer. The Δv_{th} amplitude is extracted, as shown in Fig. 2.4. The probability distribution of Δv_{th} amplitude exhibits an exponential function.

$$
f(\Delta v_{th}) = \frac{1}{\sigma} \exp\left(-\frac{\Delta v_{th}}{\sigma}\right)^{2/3} \sum_{k=1}^{\infty} \exp\left(-\frac{(\Delta v_{th})^2}{\sigma}\right)
$$
 Eq. (2.1)

The cumulative probability distribution, as illustrated in Fig. 2.5, can be expressed below,

$$
f(\Delta v_{th}) = \exp\left(-\frac{\Delta v_{th}}{\sigma}\right)
$$
 Eq. (2.2)

The Δv_{th} distribution is attributed to a current path percolation effect due to random dopants in substrate (Fig. 2.6) **[15]**. According to the Eq. (2.2), the slope of the distribution is $-1/\sigma$, as shown in Fig. 2.5. The physical meaning of σ is the average Δv_{th} induced by single charge de-trapping. Average of each stair-like ΔV_{th} trace is

shown in Fig. 2.7. The ΔV_{th} versus recovery time follows logarithmic dependence [8], **[16]**. This phenomenon is observed in bigger-size devices.

2.4 Single Charge De-trapping Time Distribution

We measure the emission times of the first de-trapping hole (τ_{el}) , the second de-trapping hole (τ_{e2}) and the third de-trapping hole (τ_{e3}), respectively (Fig. 2.4). The first three individual de-trapping time distributions (τ_{e1} , τ_{e2} , and τ_{e3}) are shown in Fig. 2.8. It is found that each de-trapping time distribution is a Gaussian-like distribution and their μ_i and σ_i are shown in Table I.

$$
f(\log(\tau_{e,i})) = \frac{1}{\sqrt{2\pi}\sigma_i} \exp\left[\frac{(\log(\tau_{e,i}) - \mu_i)^2}{2\sigma_i^2 S} \right]
$$
 Eq. (2.3)

The de-trapping time $(\tau_{e, i})$ distribution broadens with an increasing the index i. The occurrence of total τ_e distribution, as shown in Fig. 2.9, is a uniform distribution. According to the tunneling front model, the traps in the high- κ are uniform distribution in gate-to-substrate direction.

$$
\tau_e \propto \exp(\alpha x) \tag{2.4}
$$

2.5 ΔV_{th} versus Number of De-trapping Hole in NBTI **Recovery**

 ΔV_{th} distribution is measured at different recovery times. The total ΔV_{th} and number of de-trapping hole is extracted. The measurement results at the recovery time of 0.1(s), 10(s), and 1000(s) are shown in Fig. 2.10. Each data point represents a device. The slope of straight line is the average of single hole de-trapping induced ΔV_{th} . As shown in Fig. 2.10., distributions of the de-trapping hole number and ΔV_{th} broaden with an increasing recovery time.

Fig. 2.1 In this thesis, we characterize NBTI recovery in high- κ /metal gate pMOSFETs with gate length of 30nm and gate width of 80nm.

Fig. 2.2 The instrument setting is used for the fast transient measurement.

(a)

Fig. 2.3 (a) Waveforms applied to gate and drain during stress and measurement. (b) Schematic illustration of hole de-trapping formation during NBTI recovery.

Fig. 2.4 We can extract Δv_{th} and τ_e from the stair-like threshold voltage trace in NBTI recovery.

Fig. 2.5 Complementary cumulative probability distribution of single de-trapping hole induced Δv_{th} in NBTI recovery.

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Fig. 2.9 The occurrence of total τ_e distribution is uniform.

Fig. 2.10 ΔV_{th} versus Number of De-trapping Hole

Table I μ_i and σ_i of the first three de-trapping time distribution

Chapter 3

Simulation of ΔV_{th} Dispersion during NBTI **Recovery**

3.1 Preface

Physical origins of ΔV_{th} distribution during NBTI recovery in pMOSFETs are explored and characterized. Two factors are found to affect a NBTI ΔV_{th} distribution. One is the dispersion of single charge de-trapping induced Δv_{th} . The other one is the dispersion of de-trapping time. In the previous chapter, it is found that the first three individual de-trapping time distributions are Gaussian-like distributions. We attribute de-trapping time distribution to activation energy (E_a) and trap energy (E_T) distribution in the trapped charge emission model **[10]**. A statistical model combining the trapped charge emission model with the two dispersions is developed. Our simulation result fits the measured ΔV_{th} dispersion during NBTI recovery very well. Furthermore, we derive the ΔV_{th} dispersion using closed form method, and its result fit the measured data very well.

3.2 Extraction E_a **from** $\tau_{e,i}$ **Distributions**

Single charge phenomenon is measured in the previous chapter. The individual de-trapping time distribution is a Gaussian-like distribution. From the measured $\tau_{e,i}$ distributions, *E^a* is extracted using trapped charge emission model. The formula of trapped charge emission model is expressed below,

$$
\tau_{e,i} = \tau_0 \exp(\frac{E_a}{kT}) \exp(\alpha_{IL} T_{IL}) \exp(\alpha_k x_i)
$$
 Eq. (3.1)

where

$$
\tau_0 = [N_{\nu}(1 - f_{\nu})\sigma V_{th}]^{-1}
$$
 Eq. (3.1-a)

$$
\alpha_{IL} = \frac{2\sqrt{2m_{IL}^*q(E_T + \phi_B)}}{\hbar} \quad ; \quad \alpha_k = \frac{2\sqrt{2m_k^*qE_T}}{\hbar}
$$
\nEq. (3.1-b)\n
$$
m_{IL}^* = 0.41 m_0 [17] \quad ; \ m_k^* = 0.18 m_0 [18]
$$

Eq. (3.1) reveals the nature of tunneling for trapped hole emission time, $\tau_{e, i}$ *.* N_v is the effective density-of-state in the Si valence band, and $N_v(1-f_v)$ is the amount of available states in Si substrate for out-tunneling holes from high-k traps. σ_0 represents the trap cross-section. Other variables have their usual definitions **[10]**. Energy band diagram is shown in Fig. 3.1. Our estimation is that the trap energy range is about 0.8eV by a 1-D Poisson simulation. E_T ranges between 0 and 0.8 eV from the Si valence band. In the formula definition, E_T ranges between 2.7 and 3.5 eV in our **TATITUDING** devices.

The flow of E_a extraction is shown in Fig. 3.2. First, we generate $\tau_{e,i}$ from the measurement data randomly, and we also generate x_i randomly. E_T is generated by a uniform distribution. As a result, *E^a* distributions can be calculated, as shown in Fig. 3.3. These E_a distributions are almost the same, and its standard deviation of $E_a(\sigma(E_a))$ is shown in Table II. The extracted *E^a* distributions are Gaussian-like distributions.

3.3 Monte Carlo Simulation of ΔV_{th} Distribution

A Monte Carlo simulation employing the trapped charge emission model and taking account the E_a and E_T distributions is developed to simulate a ΔV_{th} distribution and its recovery time evolution. Simulation flow is shown in Fig. 3.4. First, we assume E_T is a uniform distribution, and Ea is extracted from $\tau_{e,i}$ distributions. A statistical model combining the trapped charge emission model with the *E^a* and *E^T* distribution is developed. For reproducing the de-trapping time distribution, we would calculate its average (peak) value first. As illustrated in Fig. 2.8, the trap is a uniform distribution in gate-to-substrate direction. μ_i of each de-trapping time distribution is shown in Table I. It is found that $\mu_{t+1} - \mu_t$ equal about 1. The mean value of $E_T(\mu(E_T))$ equals to 3.1eV. As expressed in Eq. (3.2), we can get the $(x_{i+1} - \Delta x_i)$. Its value is about 3\AA .

$$
\log(\tau_{i+1}/\tau_i) = \frac{1}{\ln(10)} \frac{2\sqrt{2m_k^*q}}{\hbar} \sqrt{E_T(x_{i+1} - x_i) - E_T(x_{i+1} - x_i)}
$$
Eq. (3.2)

MC simulation result is illustrated in Fig. 3.5. Its value is shown in Table III. Simulation result of total τ_e distribution is shown in Fig. 3.6. They all fit the experimental data well.

In our Monte Carlo simulation, an *E^a* is randomly selected according to extracted E_a and E_T according to a uniform distribution. The single hole de-trapping induced Δv_{th} dispersion and its de-trapping time distribution are characterized. Consequently, we can demonstrate the stair-like ΔV_{th} trace during NBTI recovery by MC simulation, as shown in Fig. 3.7. The slope of the dash line corresponds to an average ΔV_{th} caused

by a trapped hole emission. Fig. 3.8 shows probability distribution of NBTI recovery induced ΔV_{th} at 0.1s, 10s, and 1000s from measurement and Monte Carlo simulation. The ΔV_{th} distribution broadens with an increasing recovery time. The simulation results fit the measurement date well, especially in distribution tail. We extract the mean value and the variance of the ΔV_{th} distribution at different recovery times. The mean value of the ΔV_{th} follows a logarithm dependence on recovery time in five decades (Fig. 3.9). The variance of the ΔV_{th} distribution also increases with recovery time (Fig. 3.10). We can't simulate the ΔV_{th} distribution and its recovery time evolution exactly when we use the fixed E_a and E_f (Percolation Effect). Taking account the E_a and E_T distributions is essential to character the ΔV_{th} distribution during NBTI recovery.

3.4 Closed Form Derivation

In this section, we use another method to derive the ΔV_{th} distribution during NBTI recovery. Flow chart is shown in Fig. 3.11. \vec{E}_a is a Gaussian distribution with μ =0.2(eV) and σ =0.08(eV). The probability of E_a is expressed below,

$$
f(E_a) = \frac{1}{\sqrt{2\pi}\sigma} \exp[\frac{-(E_a - \mu)^2}{2\sigma^2}]
$$
 Eq. (3.4)

 E_T is a uniform distribution, and it ranges between 2.7 and 3.5 (eV).

$$
f(E_T) = const.
$$
 Eq. (3.5)

Taking the probability of E_a and E_T into Eq. (3.6), we can get the probability distribution of individual de-trapping time ($\tau_{e,m}$).

$$
\tau_{e,m}(E_a, E_T, x_m) = \tau_0 \exp(\frac{E_a}{kT}) \exp(\alpha_{IL} T_{IL}) \exp(\alpha_k x_m)
$$
 Eq. (3.6)

With the SRH model, as expressed in Eq. (3.7), we derive the probability distribution of individual de-trapping time (τ_m^*).

$$
f(\tau_{e,m}) = \frac{1}{\tau_{e,m}} \exp(-\frac{\tau_m^*}{\tau_{e,m}})
$$
 Eq. (3.7)

Then, we would calculate the de-trapping probability of each τ_m^* at different recovery time. $P(A_m)$ stands for de-trapping probability of m-th de-trapping hole at a $\mathcal{L} = \frac{1}{\sqrt{2}}$ certain recovery time(t).

$$
P(A_n) = \int_{-\infty}^{t} f(\tau_n^*) d\tau^* = \left[\begin{array}{ccc} \frac{1}{2} & \frac{1}{2} \\ \frac{1}{2} & \frac{1}{2} \\ \frac{1}{2} & \frac{1}{2} \end{array}\right]
$$
 Eq. (3.8)

 $P(B_m)$ stands for the absence of the m-th de-trapping hole, can be expressed below, **THULLER**

$$
P(B_m) = 1 - P(A_m)
$$
 Eq. (3.9)

 The probability of total amount of m de-trapping holes at a recovery time (t) can be expressed below,

$$
P(Num = 0) = P_0 = P(B_1) * P(B_2) * P(B_3) * ... * P(B_n) = P(\bigcap_{m=1}^{M} B_m)
$$
 Eq. (3.10)

$$
P(Num. = 1) = P_1 = P(A1) * P(B2) * P(B3) * \dots + P(B1) * P(A2) * P(B3) * \dots + \dots
$$

$$
P(Num. = 1) = P_1 = P(A1) * P(B2) * P(B3) * \dots + P(B1) * P(A2) * P(B3) * \dots + \dots
$$

=
$$
\sum_{i=1}^{M} \frac{P(\bigcap_{m=1}^{M} B_m)}{P(B_i)} * P(A_i)
$$
 Eq. (3.11)

$$
P(Num. = 2) = P_2 = \sum_{\substack{i,j=1 \ (j>i)}}^{M} \frac{P(\bigcap_{m=1}^{M} B_m)}{P(B_i)P(B_j)} * P(A_i)P(A_j)
$$
 Eq. (3.12)

$$
P(Num. = 3) = P_3 = \sum_{i,j,k=1}^{M} \frac{P(\bigcap_{m=1}^{M} B_m)}{P(B_i)P(B_j)P(B_k)} * P(A_i)P(A_k)
$$
 Eq. (3.13)
And so on.

M is the maximum de-trapping number of our measured devices.

As aforementioned, the trap number distribution is calculated and its average and variance are shown in Table IV. μ_A and σ_A are 3.28 (mV). μ_A is the value we measure in Fig. 2.5. Consequently, we can derive the mean and variance of ΔV_{th} at different recovery time, as expressed below,

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$$
\sigma^2 = \frac{\sum_{i=1}^n x_i^2}{n} - \mu^2
$$
 Eq. (3.14)

$$
\sigma_1^2 = \frac{\sum_{i=n_0+1}^{n_0+n_1} x_i^2}{n_1} - \mu_1^2 \qquad \qquad \sum_{i=n_0+1}^{n_0+n_1} x_i^2 = n_1 (\sigma_1^2 + \mu_1^2) \qquad \qquad \text{Eq. (3.15)}
$$

$$
\sigma_2^2 = \frac{\sum\limits_{i=n_0+n_1+1}^{n_1} x_i^2}{n_2} - \mu_2^2
$$

And so on.

$$
\sigma_{total}^{2} = \frac{\sum_{i=1}^{N} x_{i}^{2}}{N} - \mu_{total}^{2} = \frac{\sum_{i=n_{0}}^{n_{0}+n_{1}} x_{i}^{2} + \sum_{i=n_{0}+n_{1}}^{n_{0}+n_{1}+n_{2}} x_{i}^{2} + \dots}{N} - \mu_{total}^{2}
$$
\n
$$
= \frac{n_{1}(\sigma_{1}^{2} + \mu_{1}^{2}) + n_{2}(\sigma_{2}^{2} + \mu_{2}^{2}) + n_{3}(\sigma_{3}^{2} + \mu_{3}^{2}) + \dots}{N} - (P_{1}\mu_{1} + P_{2}\mu_{2} + \dots + P_{M}\mu_{M})^{2}
$$
\n
$$
= \sum_{i=1}^{M} P_{i}(\sigma_{i}^{2} + \mu_{i}^{2}) - (\sum_{i=1}^{M} P_{i}\mu_{i})^{2}
$$
\nEq. (3.17)

$$
\mu_{total} = P_1 \mu_1 + P_2 \mu_2 + ... + P_M \mu_M = \sum_{i=1}^{M} P_i \mu_i
$$
 Eq. (3.18)

where $N = n_0 + n_1 + n_2 + n_3 + \cdots$ **;** $P_i = \frac{n_i}{N}, i=0,1,2,3...$ $P_i = \frac{n}{i}$ *N* $=$

M is the maximum de-trapping number of our measured devices.

The closed form and Monte Carlo simulation result is shown in Fig. 3.12. The mean and variance of ΔV_{th} fit the results of Monte Carlo simulation very well.

Fig. 3.1 Schematic representation of band diagram and trap positions in recovery phase.

Fig. 3.2 Flow of *E^a* extraction

Fig. 3.3 Extracted E_a from $\tau_{e,i}$ Distribution

Fig. 3.4 Flow chart of Monte Carlo simulation.

Fig. 3.5 The probability distribution of trapped charge de-trapping time (τ_e) in NBTI recovery. Monte Carlo simulation result fits the measurement data well.

Fig. 3.6 MC simulation result of total τ_e distribution.

Fig. 3.7 Simulation of stair-like ΔV_{th} fluctuation.

Fig. 3.8 Probability distribution of NBTI recovery induced ΔV_{th} at 0.1s, 10s, and 1000s from measurement and Monte Carlo simulation.

Fig. 3.9 The mean value of the ΔV_{th} distribution versus NBTI recovery time from measurement and Monte Carlo simulation. Percolation has δ -function like *E^a* and *E^T* distribution.

Fig. 3.10 The variance of the ΔV_{th} distribution versus NBTI recovery time from measurement and Monte Carlo simulation. Percolation has δ -function like *E^a* and *E^T* distribution.

Fig. 3.11 Flow chart of closed form method.

(a)

Fig. 3.12 The (a) mean value and (b) variance of the ΔV_{th} distribution versus NBTI recovery time from Monte Carlo simulation and closed form method.

μ_i	$i=1$	$i=2$	$i=3$				
Experiment	-1.019	0.036	1.064				
MC Simulation	-0.982	0.029	1.035				
WWW. SALE							
σ_i	$i=1$	$i=2$	$i=3$				
Experiment	1.070	1.181	1.285				
MC Simulation	1.098	1.146	1.196				

Table III μ and σ of the first three de-trapping time distribution from measurement and Monte Carlo simulation.

Number				.	M
μ	$\mu_1 = \mu_A$	$\mu_2 = 2\mu_A$	$\mu_3 = 3\mu_A$.	$\mu_M = M\mu_A$
σ^2	$\sigma_1^2 = \sigma_A^2$	$\sigma_2^2 = 2\sigma_A^2 \sqrt{\sigma_3^2} = 3\sigma_A^2$.	$\sigma_M^2 = M \sigma_A^2$
Probability				.	$P_{\rm M}$

Table IV , ², and probability of each de-trapping hole number.

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E

Chapter 4 Conclusion

Single charge emission from traps in the high- κ gate dielectric is observed in small-size devices. Physics of single charge de-trapping induced ΔV_{th} in NBTI recovery is characterizing in this work. Two factors are found to affect a NBTI ΔV_{th} distribution. One is the dispersion of single charge de-trapping induced Δv_{th} . The other one is the dispersion of de-trapping time. The cumulative probability distribution of Δv_{th} amplitude follows the exponential distribution. The individual de-trapping time has Gaussian-like distribution. The occurrence distribution of total de-trapping time is a uniform distribution, which implies the trap is a uniform distribution in high- κ layer spatially.

From the measured $\tau_{e,i}$ distributions, E_a is extracted using trapped charge emission model. The extracted E_a is a Gaussian-like distribution. According to the Δv_{th} amplitude distribution and the de-trapping time obtained from the experiment, a Monte Carlo simulation employing the trapped charge emission model and taking account the E_a and E_T distributions is developed to simulate a ΔV_{th} distribution and its recovery time evolution. We extract the mean value and the variance of the ΔV_{th} distribution at different recovery times. The mean value of the ΔV_{th} follows a logarithm dependence on recovery time. The ΔV_{th} distribution broadens with an increasing recovery time. Furthermore, we also calculate the mean value and the variance of the ΔV_{th} distribution at different recovery times using closed form method. These two simulation results fit the experimental data well.

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