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碩士論文

整數階與分數階變革式心搏系統的渾沌及其同步與 已批准 反控制

Chaos, Its Synchronization and Anticontrol of Integral and Fractional Order Modified Heartbeat Systems

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中華民國九十五年六月

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摘 要

Summer.

本篇論文探討有關整數階與分數階變革式心搏系統的渾沌、渾沌同步及其反控 制。藉由相圖、龐卡萊圖及分歧圖等數值模擬結果,可以得知週期及渾沌的現象。 另外經由參數激發法使得兩個非耦合的整數階或分數階變革式心搏系統達到同步。 並透過外加常數項或非線性項獲得分數階變革式心搏系統的反控制。最後,成功的 利用參數激發的方式,讓整數階與分數階的變革式心搏系統達到渾沌化。

Chaos, Its Synchronization and Anticontrol of Integral and Fractional Order Modified Heartbeat Systems

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Abstract

Chaos, chaos synchronization and anticontrol for integral and fractional modified heartbeat systems are studied in the thesis. By applying numerical results, phase portrait, Poincaré maps and bifurcation diagrams, a variety of the phenomena of the periodic and chaotic motion can be presented. The synchronizations of two uncoupled integral and fractional order chaotic modified heartbeat systems are accomplished by parameter excited synchronization. Anticontrol of chaos of fractional order modified heartbeat systems can be obtained by addition of a constant term and addition of a nonlinear term. Chaotization of integral and fractional order modified heartbeat systems by parameter excited method, replacement of the parameter of first system by the function of the chaotic state variables of a second chaotic system, is successfully accomplished.

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Chapter 1 Introduction

Chaos is desirable under certain circumstances and chaotic phenomena are quite useful in many applications such as fluid mixing [1], human brain [2], and heart beat regulation [3], etc. But like many other theories in science, the theory of chaos did not get enough attention for almost two decades, ever since Edward Lorenz reported his discovery of the theory in 1963. However, the invention of high-speed computers in the late 1970s and early 1980s brought about a major revolution in nonlinear dynamical theories in general, and chaos theory in particular that is why chaos theory developed in the natural sciences during the 1970s, and the social sciences during the 1980s.

Chaos synchronization [4-12] is a very important topic in the nonlinear [13-15] science and it has been developed extensively. Recently many scientists in various fields have been attracted to investigate chaos synchronization due to its application in a variety of fields such as secure communications, chemical, physical, and biological systems, neural networks, etc. So various synchronization schemes, such as variable structure control [16], parameters adaptive control [17-24], observer based control [25, 26], active control [27-33], nonlinear control [34, 35] and so on have been successfully applied to the chaos synchronization.

Fractional calculus is a 300-year-old mathematical topic. Although it has a long history, the applications of fractional calculus to physics and engineering are just a recent focus of interest [36]. In recent years, many scholars have devoted themselves to study the applications of the fractional order system to physics and engineering such as viscoelastic systems [37], dielectric polarization, and electromagnetic waves. More recently, there is a new trend to investigate the control [38] and dynamics [39-46] of the fractional order dynamical systems [47-49]. In [37] it has been shown that nonlinear chaotic systems still have chaotic behavior when their models become fractional. In [47], chaos control was investigated for fractional chaotic systems by the "back-stepping" method of nonlinear control design. In [48] and [49], it was found that chaos exists in a fractional order Chen system with order less than 3. Linear feedback control of chaos in

this system was also studied. In [41], chaos synchronization of fractional order chaotic systems was studied. The existence and uniqueness of solutions of initial value problems for fractional order differential equations have been studied in the literature [50-53].

In addition, anticontrol [54-60] of chaos have received great attention for many research activities in recent years and it is an interesting, new and challenging phenomenon [61-63]. As a reverse process of suppressing or eliminating chaotic behaviors in order to reduce the complexity of an individual system or a coupled system, anticontrol of chaos aims at creating or enhancing the system complexity for some special applications. More precisely, anticontrolling chaos is to generate some chaotic behaviors from a given system, which is non-chaotic or even is stable originally. By fully exploiting the intrinsic nonlinearity, this "control" technique provides another dimension for feedback systems design. Its potential applications can be easily found in many fields, including typically physics, biology, engineering, and medical as well as social sciences.

In this thesis, integral and fractional order modified heartbeat systems, i.e., van der Pol systems, are investigated. Chapter 2 presents a fractional derivative and the model of modified Van der Pol system. In Chapter 3, the dynamics of integral and fractional order modified van der Pol systems are investigated and the numerical results of periodic and chaotic phenomena are presented. The synchronizations of two uncoupled integral and fractional order chaotic modified van der Pol systems are achieved by parameter excited synchronization in Chapter 4. Anticontrol of chaos of fractional order modified van der Pol systems by addition of a constant term and addition of k|x|sinx term are discussed in Chapter 5. The parameter excited chaotization of integral and fractional order modified van der Pol systems is studied in Chapter 6. Finally, the conclusions are draw.

Chapter 2

A Fractional Derivative and a Modified Heartbeat System

2.1 A fractional derivative and its approximation

There are several definitions of fractional derivatives. The commonly used definition for a general fractional derivative is the Riemann-Liouville definition [64], which is given by

$$\frac{d^{q} f(t)}{dt^{q}} = \frac{1}{\Gamma(n-q)} \frac{d^{n}}{dt^{n}} \int_{0}^{t} \frac{f(\tau)}{(t-\tau)^{q-n+1}} d\tau$$
(2.1)

where $\Gamma(\cdot)$ is the gamma function and n is an integer such that n-1 < q < n. This definition is different from the usual intuitive definition of derivative. Fortunately, the basic engineering tool for analyzing linear systems, the Laplace transform, is still applicable and works as one would expect:

$$L\left\{\frac{d^{q}f(t)}{dt^{q}}\right\} = s^{q}L\left\{f(t)\right\} - \sum_{k=0}^{n-1} s^{k} \left[\frac{d^{q-1-k}f(t)}{dt^{q-1-k}}\right]_{t=0}, \text{ for all } q, \qquad (2.2)$$

where n is an integer such that n-1 < q < n. Upon considering the initial conditions to be zero, this formula reduces to the more expected form

$$L\left\{\frac{d^{q}f(t)}{dt^{q}}\right\} = s^{q}L\left\{f(t)\right\}$$
(2.3)

An efficient method is to approximate fractional operators by using standard integer order operators. In [65-69], an effective algorithm is developed to approximate fractional order transfer functions. Basically the idea is to approximate the system behavior based on frequency domain arguments. By utilizing frequency domain techniques based on Bode diagrams, one can obtain a linear approximation of the fractional order integrator, the order of which depends on the desired bandwidth and discrepancy between the actual and the approximate magnitude Bode diagrams. In Table 1 (See Appendix) of [70], approximations for $\frac{1}{s^q}$ with $q = 0.1 \sim 0.9$ in steps 0.1 are given, with errors of approximately 2 dB. These approximations are used in the following simulations.

2.2 A modified heartbeat system and the corresponding fractional order system

Firstly, a heartbeat system, i.e. a van der Pol oscillator[71-73], driven by a periodic force is considered. The equation of motion can be written as:

$$\ddot{x} + \varphi x + a \dot{x} (x^2 - 1) - b \sin \omega t = 0$$
(2.4)

In Eq. (2.4), the linear term stands for a conservative harmonic force which determines the intrinsic oscillation frequency. The self-sustaining mechanism which is responsible for the perpetual oscillation rests on the nonlinear term. Energy exchange with the external agent depends on the magnitude of displacement $|\mathbf{x}|$ and on the sign of velocity \dot{x} . During a complete cycle of oscillation, the energy is dissipated if displacement $\mathbf{x}(t)$ is large than one, and that energy is fed-in if $|\mathbf{x}| < 1$. The time-dependent term stands for the external driving force with amplitude b and frequency ω .Eq. (2.4) can be rewritten as two first order equations:

$$\begin{cases} \dot{x} = y \\ \dot{y} = -\varphi x + a(1 - x^2)y + b\sin \omega t \end{cases}$$
(2.5)

The modified heartbeat system, i.e. the modified van der Pol system, and its fractional order system studied in this paper are

$$\begin{cases} \frac{d^{\alpha}x}{dt^{\alpha}} = y \\ \frac{d^{\beta}y}{dt^{\beta}} = -x + a(1 - x^{2})y + bz \\ \dot{z} = w \\ \dot{w} = -cz - dz^{3} \end{cases}$$
(2.6)

where α, β are integer numbers and fractional numbers respectively.

System (2.6) can be separated into two parts:

$$\begin{cases} \frac{d^{\alpha}x}{dt^{\alpha}} = y \\ \frac{d^{\beta}y}{dt^{\beta}} = -x + a(1 - x^{2})y + bz \end{cases}$$
(2.7)

and

$$\begin{cases} \dot{z} = w \\ \dot{w} = -cz - dz^3 \end{cases}$$
(2.8)

In Eq. (2.5) changing the integral order derivatives to the fractional order derivations and replaceing $\sin \omega t$ by z which is the periodic time function solution of the nonlinear oscillator (2.8), we obtain system (2.7). In Eq. (2.8) if d = 0, z is a sinusoidal function of time. Now d \neq 0, z is a periodic motion of time but not a sinusoidal function of time. As a result, system (2.7) can be considered as a nonautonomous system with two states, while system (2.6) consisting of Eq. (2.7) and Eq. (2.8) can be considered as an autonomous system with four states. When $\alpha = \beta = 1$, Eq.(2.6) is the modified van der Pol system.



Chapter 3

Chaos in a Modified Heartbeat System and in Its Fractional Order Systems

Chaos in a modified van der Pol system and in its fractional order systems is studied in this chapter. It is found that chaos exists both in the system and in the fractional order systems with total order from 1.8 down to 0.8 much less than the number of states of the system, two. By phase portraits, Poincaré maps and bifurcation diagrams, the chaotic behaviors of integral and fractional order modified van der Pol systems are presented.

In this chapter, phase portraits and bifurcation diagrams are studied for system (2.6) for $\alpha + \beta \le 2$ and three parameters *a*, *c*, *d* are chosen as a = 5, c = 0.01, d = 0.001. A time step of 0.001 is used.

SUMMER

3.1 Chaos in a integral order modified heartbeat system

Let $\alpha = \beta = 1$. Fig. 3.1 shows the bifurcation diagram of the 2 order system. It is shown that chaos exists when $b \in [0, 1.0]$. Fig. 3.2 is the phase portrait of chaotic motion with b = 1.0. Fig. 3.3 ~ 3.5 are phase portraits of periodic motions with b = 1.1, 1.5, 3, respectively.

3.2 Chaos in a fractional order modified heartbeat system

Case 1 Let $\alpha = 0.9$, $\beta = 0.9$.

Fig. 3.6 shows the bifurcation diagram of the 1.8 order system. It is shown that chaos exists when $b \in [0, 9.7]$. Fig. 3.7 is the phase portrait of chaotic motion with b = 9.7. Fig. 3.8 ~ 3.11 are phase portraits of periodic motions with b = 9.8, 14, 23, 40 respectively.

Case 2 Let $\alpha = 0.9$, $\beta = 0.8$.

Fig. 3.12 shows the bifurcation diagram of the 1.7 order system. It is shown that chaos exists when $b \in [0, 9.8]$. Fig. 3.13 is the phase portrait of chaotic motion with b = 9.8. Fig. 3.14 ~ 3.17 are phase portraits of periodic motions

with b = 9.9, 12, 25, 30, respectively.

Case 3 Let $\alpha = 0.8$, $\beta = 0.9$.

Fig. 3.18 shows the bifurcation diagram of the 1.7 order system. It is shown that chaos exists when $b \in [0, 10.1]$. Fig. 3.19 is the phase portrait of chaotic motion with b = 10.1. Fig. 3.20 ~ 3.24 are phase portraits of periodic motions with b = 10.2, 15, 23, 30, 45 respectively.

Case 4 Let $\alpha = 0.8$, $\beta = 0.8$.

Fig. 3.25 shows the bifurcation diagram of the 1.6 order system. It is shown that chaos exists when $b \in [0, 10.1]$. Fig. 3.26 is the phase portrait of chaotic motion with b = 10.1. Fig. 3.27 ~ 3.31 are phase portraits of periodic motions with b = 10.2, 14.5, 20, 35, 40 respectively.

Case 5 Let $\alpha = 0.7$, $\beta = 0.7$.

Fig. 3.32 shows the bifurcation diagram of the 1.4 order system. It is shown that chaos exists when $b \in [0, 7.9]$. Fig. 3.33 is the phase portrait of chaotic motion with b = 7.9. Fig. 3.34 ~ 3.37 are phase portraits of periodic motions with b = 8.0, 15, 35, 40 respectively.

Case 6 Let $\alpha = 0.6$, $\beta = 0.6$.

Fig. 3.38 shows the bifurcation diagram of the 1.2 order system. It is shown that chaos exists when $b \in [0, 6.0]$. Fig. 3.39 is the phase portrait of chaotic motion with b = 6.0. Fig. 3.40 ~ 3.44 are phase portraits of periodic motions with b = 6.1, 6.5, 9.5, 20, 45 respectively.

Case 7 Let $\alpha = 0.5$, $\beta = 0.5$.

Fig. 3.45 shows the bifurcation diagram of the 1.0 order system. It is shown that chaos exists when $b \in [0, 4.2]$. Fig. 3.46 is the phase portrait of chaotic motion with b = 2. Fig. 3.47 ~ 3.49 are phase portraits of periodic motions with b = 6.0, 10, 15 respectively.

Case 8 Let
$$\alpha = 0.4$$
, $\beta = 0.4$.

Fig. 3.50 shows the bifurcation diagram of the 0.8 order system. It is shown that chaos exists when $b \in [0, 1.8]$. Fig. 3.51 is the phase portrait of chaotic motion with b = 0.7. Fig. 3.52 \checkmark Fig. 3.53 are phase portraits of periodic

motions with b = 0.7, 1.3 respectively. When we tried to reduce the total order to 0.6, the phase portraits become periodic motions, as shown in Fig. 3.54 and Fig. 3.55, for any *b* value.

Chaos in modified van der Pol system and in its fractional order systems is studied in this chapter. It is found that the range of the chaos in the system gradually decreases as the total order number $\alpha + \beta$ decreases. Nine cases for $0.8 \le (\alpha + \beta) \le 2.0$ are studied. The lowest total order for chaos existence in the system is found to be 0.8.





Fig. 3.2 The phase portrait for $\alpha = \beta = 1, b=1.0$.



Fig. 3.4 The phase portrait for $\alpha = \beta = 1, b=1.5$.





Fig. 3.7 The phase portrait for $\alpha = 0.9$, $\beta = 0.9$, b = 9.7.



Fig. 3.9 The phase portrait for $\alpha = 0.9$, $\beta = 0.9$, b = 14.



Fig. 3.11 The phase portrait for $\alpha = 0.9$, $\beta = 0.9$, b = 40.



Fig 3.13 The phase portrait for $\alpha = 0.9$, $\beta = 0.8$, b = 9.8.



Fig. 3.15 The phase portrait for $\alpha = 0.9$, $\beta = 0.8$, b = 12.



Fig. 3.17 The phase portrait for $\alpha = 0.9$, $\beta = 0.8$, b = 30.



Fig. 3.19 The phase portrait for $\alpha = 0.8$, $\beta = 0.9$, b = 10.1.



Fig. 3.21 The phase portrait for $\alpha = 0.8$, $\beta = 0.9$, b = 15.



Fig. 3.23 The phase portrait for $\alpha = 0.8$, $\beta = 0.9$, b = 30.





Fig. 3.26 The phase portrait for $\alpha = 0.8$, $\beta = 0.8$, b = 10.1.



Fig. 3.28 The phase portrait for $\alpha = 0.8$, $\beta = 0.8$, b = 14.5.


Fig. 3.30 The phase portrait for $\alpha = 0.8$, $\beta = 0.8$, b = 35.





Fig. 3.33 The phase portrait for $\alpha = 0.7$, $\beta = 0.7$, b = 7.9.



Fig. 3.35 The phase portrait for $\alpha = 0.7$, $\beta = 0.7$, b = 15.



Fig. 3.37 The phase portrait for $\alpha = 0.7$, $\beta = 0.7$, b = 40.



Fig. 3.39 The phase portrait for $\alpha = 0.6$, $\beta = 0.6$, b = 6.0.



Fig. 3.41 The phase portrait for $\alpha = 0.6$, $\beta = 0.6$, b = 6.5.



Fig. 3.43 The phase portrait for $\alpha = 0.6$, $\beta = 0.6$, b = 20.





Fig. 3.46 The phase portrait for $\alpha = 0.5$, $\beta = 0.5$, b = 2.



Fig. 3.48 The phase portrait for $\alpha = 0.5$, $\beta = 0.5$, b = 10.





Fig. 3.51 The phase portrait for $\alpha = 0.4$, $\beta = 0.4$, b = 0.7.



Fig. 3.53 The phase portrait for $\alpha = 0.4$, $\beta = 0.4$, b = 6.0



Fig. 3.55 The phase portrait for $\alpha = 0.3$, $\beta = 0.3$, b = 1.3.

Chapter 4

Parameter Excited Chaos Synchronizations of Integral and Fractional Order Modified Heartbeat Systems

In this chapter, the synchronizations of two uncoupled integral and fractional order chaotic modified heartbeat systems, i.e. modified van der Pol systems, are achieved by replacing the corresponding parameters of two systems by the same function of the chaotic state variables of a third chaotic system. It is named parameter excited synchronization which can be successfully obtained for very low total fractional order 0.2. Numerical simulations are illustrated by phase portraits, Poincaré maps and state error plots.

4.1 The scheme of parameter excited chaos synchronizations

In this section, both parameters *b* of two identical modified fractional order van der Pol systems with different initial conditions

$$\begin{cases} \frac{d^{a} x_{1}}{dt^{a}} = y_{1} \\ \frac{d^{\beta} y_{1}}{dt^{\beta}} = -x_{1} + a(1 - x_{1}^{2})y_{1} + bz_{1} \\ \dot{z}_{1} = w_{1} \\ \dot{w}_{1} = -cz_{1} - dz_{1}^{3} \end{cases}$$

$$\begin{cases} \frac{d^{a} x_{2}}{dt^{a}} = y_{2} \\ \frac{d^{\beta} y_{2}}{dt^{\beta}} = -x_{2} + a(1 - x_{2}^{2})y_{2} + bz_{2} \\ \dot{z}_{2} = w_{2} \\ \dot{w}_{2} = -cz_{2} - dz_{2}^{3} \end{cases}$$

$$(4.1)$$

are replaced by the same cosine function of state variables of a third system (2.6). We study the parameter excited chaos synchronization for various α , β . The parameters a, cof systems (4.1) and (4.2) are adjusted to achieve synchronization for different α and β . d = 0.001 is fixed. Errors $e_1 = x_2 - x_1$, $e_2 = y_2 - y_1$ are defined. If $e_1 \rightarrow 0$, $e_2 \rightarrow 0$ when $t \to \infty$, synchronization is obtained. A time step of 0.001 is used. Three cases are studied.

4.2 Parameter Excited Chaos Synchronizations by kcosx kcosy kcosxcosy

Then, the corresponding parameters b of two uncoupled identical integral and fractional order chaotic modified van der Pol systems will be replaced by *kcosx, kcosy, kcosx, sxcosy, where x, y* are the states of a third identical system (2.6).

4.2.1 Parameter Excited Chaos Synchronizations by *kcosx*

Let the parameter *b* of Eqs. (4.1) and (4.2) be replaced by *kcosx* where *x* is the state variable of system (2.6). Parameter *k* is chosen as k = 0.5.

- Case 1 Let $\alpha = \beta = 1$. Fig. 4.1 shows that the chaos exists when a = 2, c = 10. Fig. 4.2 shows the error dynamics for synchronization.
- Case 2 Let $\alpha = \beta = 0.9$. Fig. 4.3 shows that the chaos exists when a = 5, c = 5. Fig. 4.4 shows the error dynamics for synchronization.
- Case 3 Let $\alpha = \beta = 0.8$. Fig. 4.5 shows that the chaos exists when a = 5, c = 5. Fig. 4.6 shows the error dynamics for synchronization.
- Case 4 Let $\alpha = \beta = 0.7$. Fig. 4.7 shows that the chaos exists when a = 2.5, c = 35. Fig. 4.8 shows the error dynamics for synchronization.
- Case 5 Let $\alpha = \beta = 0.6$. Fig. 4.9 shows that the chaos exists when a = 2.5, c = 50. Fig. 4.10 shows the error dynamics for synchronization.
- Case 6 Let $\alpha = \beta = 0.5$. Fig. 4.11 shows that the chaos exists when a = 2, c = 35. Fig. 4.12 shows the error dynamics for synchronization.
- Case 7 Let $\alpha = \beta = 0.4$. Fig. 4.13 shows that the chaos exists when a = 2, c = 65. Fig. 4.14 shows the error dynamics for synchronization.
- Case 8 Let $\alpha = \beta = 0.3$. Fig. 4.15 shows that the chaos exists when a = 10, c = 65. Fig. 4.16 shows the error dynamics for synchronization.
- Case 9 Let $\alpha = \beta = 0.2$. Fig. 4.17 shows that the chaos exists when a = 20, c = 55. Fig. 4.18 shows the error dynamics for synchronization.
- Case 10 Let $\alpha = \beta = 0.1$. Fig. 4.19 shows that the chaos exists when a = 10, c = 0.5. Fig. 4.20 shows the error dynamics for synchronization.

4.2.2 Parameter Excited Chaos Synchronizations by kcosy

Let the parameter b of Eqs. (4.1) and (4.2) be replaced by *kcosy* where y is the state variable of system (2.6). Parameter k is chosen as k = 0.5.

- Case 1 Let $\alpha = \beta = 1$. Fig. 4.21 shows that the chaos exists when a = 2, c = 10. Fig. 4.22 shows the error dynamics for synchronization.
- Case 2 Let $\alpha = \beta = 0.9$. Fig. 4.23 shows that the chaos exists when a = 5, c = 5. Fig. 4.24 shows the error dynamics for synchronization.
- Case 3 Let $\alpha = \beta = 0.8$. Fig. 4.25 shows that the chaos exists when a = 5, c = 5. Fig. 4.26 shows the error dynamics for synchronization.
- Case 4 Let $\alpha = \beta = 0.7$. Fig. 4.27 shows that the chaos exists when a = 2.5, c = 35. Fig. 4.28 shows the error dynamics for synchronization.
- Case 5 Let $\alpha = \beta = 0.6$. Fig. 4.29 shows that the chaos exists when a = 2.5, c = 50. Fig. 4.30 shows the error dynamics for synchronization.
- Case 6 Let $\alpha = \beta = 0.5$. Fig. 4.31 shows that the chaos exists when a = 2, c = 35. Fig. 4.32 shows the error dynamics for synchronization.
- Case 7 Let $\alpha = \beta = 0.4$. Fig. 4.33 shows that the chaos exists when a = 2, c = 65. Fig. 4.34 shows the error dynamics for synchronization.
- Case 8 Let $\alpha = \beta = 0.3$. Fig. 4.35 shows that the chaos exists when a = 5, c = 5. Fig. 4.36 shows the error dynamics for synchronization.
- Case 9 Let $\alpha = \beta = 0.2$. Fig. 4.37 shows that the chaos exists when a = 20, c = 55. Fig. 4.38 shows the error dynamics for synchronization.
- Case 10 Let $\alpha = \beta = 0.1$. Fig. 4.39 shows that the chaos exists when a = 10, c = 1. Fig. 4.40 shows the error dynamics for synchronization.

4.2.3 Parameter Excited Chaos Synchronizations by *kcosxcosy*

Let the parameter *b* of Eqs. (4.1) and (4.2) be replaced by *kcosxcosy* where *x*, *y* are the state variables of system (6). Parameter *k* is chosen as k = 0.5.

- Case 1 Let $\alpha = \beta = 1$. Fig. 4.41 shows that the chaos exists when a = 2, c = 10. Fig. 4.42 shows the error dynamics for synchronization.
- Case 2 Let $\alpha = \beta = 0.9$. Fig. 4.43 shows that the chaos exists when a = 5, c = 5. Fig. 4.44 shows the error dynamics for synchronization.

- Case 3 Let $\alpha = \beta = 0.8$. Fig. 4.45 shows that the chaos exists when a = 10, c = 5. Fig. 4.46 shows the error dynamics for synchronization.
- Case 4 Let $\alpha = \beta = 0.7$. Fig. 4.47 shows that the chaos exists when a = 2.5, c = 35. Fig. 4.48 shows the error dynamics for synchronization.
- Case 5 Let $\alpha = \beta = 0.6$. Fig. 4.49 shows that the chaos exists when a = 2.5, c = 50. Fig. 4.50 shows the error dynamics for synchronization.
- Case 6 Let $\alpha = \beta = 0.5$. Fig. 4.51 shows that the chaos exists when a = 5.5, c = 45. Fig. 4.52 shows the error dynamics for synchronization.
- Case 7 Let $\alpha = \beta = 0.4$. Fig. 4.53 shows that the chaos exists when a = 2, c = 65. Fig. 4.54 shows the error dynamics for synchronization.
- Case 8 Let $\alpha = \beta = 0.3$. Fig. 4.55 shows that the chaos exists when a = 4, c = 35. Fig. 4.56 shows the error dynamics for synchronization.
- Case 9 Let $\alpha = \beta = 0.2$. Fig. 4.57 shows that the chaos exists when a = 2, c = 1. Fig. 4.58 shows the error dynamics for synchronization.
- Case 10 Let $\alpha = \beta = 0.1$. Fig. 4.59 shows that the chaos exists when a = 2, c = 10. Fig. 4.60 shows the error dynamics for synchronization.

In this paper, parameter excited chaos synchronizations of uncoupled integral and fractional order modified van der Pol systems are studied by means of phase portraits, Poincaré maps and error dynamics plots. It is found that this approach is very effective even for very low total fractional order 0.2.



Fig. 4.1 The phase portrait for replacing *b* by *kcosx* with $\alpha = \beta = 1$, a = 2, c = 10.



Fig. 4.2 Error dynamics for Case 1.



Fig. 4.3 The phase portrait for replacing *b* by *kcosx* with $\alpha = \beta = 0.9$, a = 5, c = 5.



Fig. 4.4 Error dynamics for Case 2.



Fig. 4.5 The phase portrait for replacing *b* by *kcosx* with $\alpha = \beta = 0.8$, a = 5, c = 5.



Fig. 4.6 Error dynamics for Case 3.



Fig. 4.7 The phase portrait for replacing *b* by *kcosx* with $\alpha = \beta = 0.7$, a = 2.5, c = 35.



Fig. 4.8 Error dynamics for Case 4.



Fig. 4.9 The phase portrait for replacing *b* by *kcosx* with $\alpha = \beta = 0.6$, a = 2.5, c = 50.



Fig. 4.10 Error dynamics for Case 5.



Fig. 4.11 The phase portrait for replacing *b* by *kcosx* with $\alpha = \beta = 0.5$, a = 2, c = 35.



Fig. 4.12 Error dynamics for Case 6.



Fig. 4.13 The phase portrait for replacing *b* by *kcosx* with $\alpha = \beta = 0.4$, a = 2, c = 65.



Fig. 4.14 Error dynamics for Case 7.



Fig. 4.15 The phase portrait for replacing *b* by *kcosx* with $\alpha = \beta = 0.3$, a = 10, c = 65.



Fig. 4.16 Error dynamics for Case 8.



Fig. 4.17 The phase portrait for replacing *b* by *kcosx* with $\alpha = \beta = 0.2$, a = 20, c = 55.



Fig. 4.18 Error dynamics for Case 9.



Fig. 4.19 The phase portrait for replacing *b* by *kcosx* with $\alpha = \beta = 0.1$, a = 10, c = 0.5.



Fig. 4.20 Error dynamics for Case 10.



Fig. 4.21 The phase portrait for replacing *b* by *kcosy* with $\alpha = \beta = 1$, a = 2, c = 10.



Fig. 4.22 Error dynamics for Case 1.



Fig. 4.23 The phase portrait for replacing *b* by *kcosy* with $\alpha = \beta = 0.9$, a = 5, c = 5.



Fig. 4.24 Error dynamics for Case 2.



Fig. 4.25 The phase portrait for replacing *b* by *kcosy* with $\alpha = \beta = 0.8$, a = 5, c = 5.



Fig. 4.26 Error dynamics for Case 3.



Fig. 4.27 The phase portrait for replacing *b* by *kcosy* with $\alpha = \beta = 0.7$, a = 2.5, c = 35.



Fig. 4.28 Error dynamics for Case 4.



Fig. 4.29 The phase portrait for replacing *b* by *kcosy* with $\alpha = \beta = 0.6$, a = 2.5, c = 50.



Fig. 4.30 Error dynamics for Case 5.



Fig. 4.31 The phase portrait for replacing *b* by *kcosy* with $\alpha = \beta = 0.5$, a = 2, c = 35.



Fig. 4.32 Error dynamics for Case 6.



Fig. 4.33 The phase portrait for replacing *b* by *kcosy* with $\alpha = \beta = 0.4$, a = 2, c = 65.



Fig. 4.34 Error dynamics for Case 7.


Fig. 4.35 The phase portrait for replacing *b* by *kcosy* with $\alpha = \beta = 0.3$, a = 10, c = 65.



Fig. 4.36 Error dynamics for Case 8.



Fig. 4.37 The phase portrait for replacing *b* by *kcosy* with $\alpha = \beta = 0.2$, a = 2, c = 1.



Fig. 4.38 Error dynamics for Case 9.



Fig. 4.39 The phase portrait for replacing *b* by *kcosy* with $\alpha = \beta = 0.1$, a = 10, c = 1.



Fig. 4.40 Error dynamics for Case 10.



Fig. 4.41 The phase portrait for replacing *b* by *kcosxcosy* with $\alpha = \beta = 1$, a = 2, c = 10.



Fig. 4.42 Error dynamics for Case 1.



Fig. 4.43 The phase portrait for replacing *b* by *kcosxcosy* with $\alpha = \beta = 0.9$, a = 5, c = 5.



Fig. 4.44 Error dynamics for Case 2.



Fig. 4.45 The phase portrait for replacing *b* by *kcosxcosy* with $\alpha = \beta = 0.8$, a = 10, c = 5.



Fig. 4.46 Error dynamics for Case 3.



Fig. 4.47 The phase portrait for replacing *b* by *kcosxcosy* with $\alpha = \beta = 0.7$, a = 2.5, c = 35.



Fig. 4.48 Error dynamics for Case 4.



Fig. 4.49 The phase portrait for replacing *b* by *kcosxcosy* with $\alpha = \beta = 0.6$, a = 2.5, c = 50.



Fig. 4.50 Error dynamics for Case 5.



Fig. 4.51 The phase portrait for replacing *b* by *kcosxcosy* with $\alpha = \beta = 0.5$, a = 5.5,



Fig. 4.52 Error dynamics for Case 6.



Fig. 4.53 The phase portrait for replacing *b* by *kcosxcosy* with $\alpha = \beta = 0.4$, a = 2, c = 65.



Fig. 4.54 Error dynamics for Case 7.



Fig. 4.55 The phase portrait for replacing *b* by *kcosxcosy* with $\alpha = \beta = 0.3$, a = 4, c = 35.



Fig. 4.56 Error dynamics for Case 8.



Fig. 4.57 The phase portrait for replacing *b* by *kcosxcosy* with $\alpha = \beta = 0.2$, a = 2, c = 1.



Fig. 4.58 Error dynamics for Case 9.



Fig. 4.59 The phase portrait for replacing *b* by *kcosxcosy* with $\alpha = \beta = 0.1$, a = 2, c = 10.



Fig. 4.60 Error dynamics for Case 10.

Chapter 5

Anticontrol of Chaos of the Fractional Order Modified Heartbeat Systems

Anticontrol of chaos of fractional order modified heartbeat systems is studied. Anticontrol of chaos is making a non-chaotic dynamical system chaotic, i.e., the regular behavior will be destroyed and replaced by chaotic behavior. Addition of a constant term and addition of k/x/sinx term where x is a state of the system are used to anticontrol the system effectively. By applying numerical results, phase portrait, Poincaré maps and bifurcation diagrams a variety of the phenomena of the chaotic motion can be presented.

5.1 Anticontrol of chaos by addition of a constant term

We add a constant term k in the second equation of system (2.6) and it become:

$$\begin{cases} \frac{d^{\alpha} x}{dt^{\alpha}} = y \\ \frac{d^{\beta} y}{dt^{\beta}} = -x + a(1 - x^{2})y + bz + k \end{cases}$$

$$(5.1)$$

$$\dot{z} = w$$

$$\dot{w} = -cz - dz^{3}$$

In system (5.1), the parameter *b* is adjusted to achieve periodic motion for different α and β when k = 0. *a*, *c*, *d* are fixed and they are chosen as a = 5, c = 0.01, d = 0.001.

- Case 1 Let $\alpha = \beta = 0.9$, b = 2.5, and $k \in [0.9,1.2]$. Fig 5.1(a) shows the bifurcation diagram of the 1.8 order system. Fig 5.1(b) ~ 5.1(d) are the phase portraits with k = 0, 1.05, 1.1.
- Case 2 Let $\alpha = \beta = 0.8$, b = 1.5, and $k \in [0.6,1.2]$. Fig 5.2(a) shows the bifurcation diagram of the 1.6 order system. Fig 5.2(b) ~ 5.2(d) are the phase portraits with k = 0, 0.87, 1.05.
- Case 3 Let $\alpha = \beta = 0.7$, b = 1.3, and $k \in [0.1,1.6]$. Fig 5.3(a) shows the bifurcation diagram of the 1.4 order system. Fig 5.3(b) ~ 5.3(d) are the phase portraits with k = 0, 1.5, 1.6.

- Case 4 Let $\alpha = \beta = 0.6$, b = 1, and $k \in [1.282, 1.312]$. Fig 5.4(a) shows the bifurcation diagram of the 1.2 order system. Fig 5.4(b) ~ 5.4(d) are the phase portraits with k = 0, 1.291, 1.298.
- Case 5 Let $\alpha = \beta = 0.5$, b = 1, and $k \in [1.32, 1.42]$. Fig 5.5(a) shows the bifurcation diagram of the 1.0 order system. Fig 5.5(b) ~ 5.5(d) are the phase portraits with k = 0, 1.35, 1.38.
- Case 6 Let $\alpha = \beta = 0.4$, b = 1.5, and $k \in [1.88,1.94]$. Fig 5.6(a) shows the bifurcation diagram of the 0.8 order system. Fig 5.6(b) ~ 5.6(d) are the phase portraits with k = 0, 1.915, 1.93.
- Case 7 Let $\alpha = \beta = 0.3$, b = 1.5, and $k \in [1.89, 1.98]$. Fig 5.7(a) shows the bifurcation diagram of the 0.6 order system. Fig 5.7(b) ~ 5.7(d) are the phase portraits with k = 0, 1.9, 1.94.

5.2 Anticontrol of chaos by addition of a nonlinear term

We add a non-linear term k/x/sinx in the second equation of system (2.6) and it become:

$$\begin{cases} \frac{d^{\alpha}x}{dt^{\alpha}} = y \\ \frac{d^{\beta}y}{dt^{\beta}} = -x + a(1-x^{2})y + bz + k \mid x \mid sinx \\ \dot{z} = w \\ \dot{w} = -cz - dz^{3} \end{cases}$$
(5.2)

In system (5.2), the parameter *b* is also adjusted to achieve periodic motion for different α and β when k = 0. *a*, c, *d* are fixed as 5.1.

- Case 1 Let $\alpha = \beta = 0.9$, b = 2.5, and $k \in [0.2, 0.4]$. Fig 5.8(a) shows the bifurcation diagram of the 1.8 order system. Fig 5.8(b) ~ 5.8(d) are the phase portraits with k = 0, 0.35, 0.4.
- Case 2 Let $\alpha = \beta = 0.8$, b = 2.5, and $k \in [1.0,1.3]$. Fig 5.9(a) shows the bifurcation diagram of the 1.6 order system. Fig 5.9(b) ~ 5.9(d) are the phase portraits with k = 0, 1.03, 1.2.
- Case 3 Let $\alpha = \beta = 0.7$, b = 0.5, and $k \in [1.19, 1.24]$. Fig 5.10(a) shows the bifurcation diagram of the 1.4 order system. Fig 5.10(b) ~ 5.10(d) are the phase portraits

with *k* = 0, 1.205, 1.215.

- Case 4 Let $\alpha = \beta = 0.6$, b = 1, and $k \in [0.8,1.3]$. Fig 5.11(a) shows the bifurcation diagram of the 1.2 order system. Fig 5.11(b) ~ 5.11(d) are the phase portraits with k = 0, 0.95, 1.05.
- Case 5 Let $\alpha = \beta = 0.5$, b = 0.5, and $k \in [0.1,0.7]$. Fig 5.12(a) shows the bifurcation diagram of the 1.0 order system. Fig 5.12(b) ~ 5.12(d) are the phase portraits with k = 0, 0.4, 0.6.
- Case 6 Let $\alpha = \beta = 0.4$, b = 1, and $k \in [1.16, 1.37]$. Fig 5.13(a) shows the bifurcation diagram of the 0.8 order system. Fig 5.13(b) ~ 5.13(d) are the phase portraits with k = 0, 1.25, 1.3.

Anticontrol of chaos in the fractional order systems of a modified van der Pol system are studied in the chapter. An efficient way to transform a non-chaotic dynamical system into a chaotic one is easily made by addition of a constant term or by addition of k/x/sinx term where x is a state variable of the system. It is found that chaos exists in the fractional order systems with order from 1.8 down to 0.6 for the addition of constant term, and from 1.8 down to 0.8 for the addition of k/x/sinx term.

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Fig 5.1 The bifurcation diagram for $\alpha = \beta = 0.9$ and phase portrait for $\alpha = \beta = 0.9$ with k = 0, k = 1.05, k = 1.1 respectively.



Fig 5.2 The bifurcation diagram for $\alpha = \beta = 0.8$ and phase portrait for $\alpha = \beta = 0.8$ with k = 0, 0.87, 1.05, respectively.



Fig 5.3 The bifurcation diagram for $\alpha = \beta = 0.7$ and phase portrait for $\alpha = \beta = 0.7$ with k = 0, 1.5, 1.6, respectively.



Fig 5.4 The bifurcation diagram for $\alpha = \beta = 0.6$ and phase portrait for $\alpha = \beta = 0.6$ with k = 0, 1.291, 1.298, respectively.



Fig 5.5 The bifurcation diagram for $\alpha = \beta = 0.5$ and phase portrait for $\alpha = \beta = 0.5$ with k = 0, 1.35, 1.38, respectively.



Fig 5.6 The bifurcation diagram for $\alpha = \beta = 0.4$ and phase portrait for $\alpha = \beta = 0.4$ with k = 0, 1.915, 1.93, respectively.



Fig 5.7 The bifurcation diagram for $\alpha = \beta = 0.3$ and phase portrait for $\alpha = \beta = 0.3$ with k = 0, 1.9, 1.94, respectively.



Fig 5.8 The bifurcation diagram for $\alpha = \beta = 0.9$ and phase portrait for $\alpha = \beta = 0.9$ with k = 0, 0.35, 0.4, respectively.







Fig 5.10 The bifurcation diagram for $\alpha = \beta = 0.7$ and The phase portrait for $\alpha = \beta = 0.7$ with k = 0, 1.205, 1.215, respectively.







Fig 5.12 The bifurcation diagram for $\alpha = \beta = 0.5$ and phase portrait for $\alpha = \beta = 0.5$ with k = 0, 0.4, 0.6, respectively.





Chapter 6

Parameter Excited Chaotization of Integral and Fractional Order Modified Heartbeat Systems

Chaotization of integral and fractional order modified heartbeat systems by replacing the parameter of modified heartbeat system by the function of the chaotic state variables of a second chaotic nano resonator system is studied. It is named parameter excited chaotization which can be successfully obtained for very low total fractional order 0.2. By phase portraits, Poincaré maps and bifurcation diagrams, the simulation results of fractional order modified van der Pol systems are presented.

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6.1 A nano resonator system and the chaotization scheme

Chaotization, i.e. anticontrol of chaos, is making a non-chaotic dynamical system chaotic. This means that the regular behavior will be destroyed and replaced by chaotic behavior.

Nano resonator system is a modified form of nonlinear damped Mathieu system which is obtained when the nano Mathieu oscillator has nonlinear time-dependent spring constant. The nonlinear damped Mathieu system is a nonautonomous system with two states x and y:

$$\begin{cases} \frac{dx}{dt} = y \\ \frac{dy}{dt} = -(e+f\sin\omega_1 t)x - (e+f\sin\omega_1 t)x^3 - gy + h\sin\omega_2 t \end{cases}$$
(6.1)

where *a*, *b*, *c*, *d* are constant parameters, and ω_1 , ω_2 are circular frequencies. Let $\omega_1 = \omega_2 = \omega$, and replace *sin* ωt by *z* which is the periodic time function solution of the nonlinear oscillator

$$\begin{cases} \frac{dz}{dt} = w \\ \frac{dw}{dt} = -iz - jz^3 \end{cases}$$
(6.2)

where i, j are constants. Then we have the modified nonlinear damped Mathieu system, i.e. the nano resonator system:

$$\begin{cases} \frac{dx}{dt} = y \\ \frac{dy}{dt} = -(e+fz)x - (e+fz)x^3 - gy + hz \\ \frac{dz}{dt} = w \\ \frac{dw}{dt} = -iz - jz^3 \end{cases}$$
(6.3)

It becomes an autonomous system with four states where e, f, g, h, i and j are constants of the system. In our numerical simulations, six parameters e = 0.2, f = 0.2, g = 0.4, h = 50, i = 1 and j = 0.3 are chosen to make the states x and y of system (6.3) be a chaotic [74].

Next, let the parameter *b* of system (2.6) be replaced by *ksiny* where *k* is a constant and *y* is a chaotic state variable of system (6.3). We study the parameter excited chaotization of system (2.6) for various α , β .

6.2 Chaotization by parameter excited method

The parameter k is adjusted to achieve chaos for different α and β . a = 5, c = 4, i = 0.0001 are fixed and $k \in [0,5]$.

6.2.1 Chaotization of a integral order modified heartbeat system

Let $\alpha = \beta = 1$. Fig 6.1(a) shows the bifurcation diagram of the 2 order system. Fig 6.1(b) is the nonchaotic phase portrait of the system (2.6) with b = 2.5. Fig 6.1(c) shows chaotic phase portrait when k = 4.

6.2.2 Chaotization of a fraction order modified van der Pol system

- Case 1 Let $\alpha = \beta = 0.9$. Fig 6.2(a) shows the bifurcation diagram of the 1.8 order system. Fig 6.2(b) is the nonchaotic phase portrait of the system (2.6) with b = 2.5. Fig 6.2(c) shows chaotic phase portrait when k = 4.
- Case 2 Let $\alpha = \beta = 0.8$. Fig 6.3(a) shows the bifurcation diagram of the 1.6 order system. Fig 6.3(b) is the nonchaotic phase portrait of the system (2.6) with b = 1.5. Fig 6.3(c) shows chaotic phase portrait when k = 3.
- Case 3 Let $\alpha = \beta = 0.7$. Fig 6.4(a) shows the bifurcation diagram of the 1.4 order system. Fig 6.4(b) is the nonchaotic phase portrait of the system (2.6) with b = 1.3. Fig 6.4(c) shows chaotic phase portrait when k = 3.5.
- Case 4 Let $\alpha = \beta = 0.6$. Fig 6.5(a) shows the bifurcation diagram of the 1.2 order system. Fig 6.5(b) is the nonchaotic phase portrait of the system (2.6) with b = 1. Fig 6.5(c) shows chaotic phase portrait when k = 3.5.
- Case 5 Let $\alpha = \beta = 0.5$. Fig 6.6(a) shows the bifurcation diagram of the 1.0 order system. Fig 6.6(b) is the nonchaotic phase portrait of the system (2.6) with b = 1. Fig 6.6(c) shows chaotic phase portrait when k = 3.
- Case 6 Let $\alpha = \beta = 0.4$. Fig 6.7(a) shows the bifurcation diagram of the 0.8 order system. Fig 6.7(b) is the nonchaotic phase portrait of the system (2.6) with b = 1.5. Fig 6.7(c) shows chaotic phase portrait when k = 2.5.

Case 7 Let $\alpha = \beta = 0.3$. Fig 6.8(a) shows the bifurcation diagram of the 0.6

order system. Fig 6.8(b) is the nonchaotic phase portrait of the system (2.6) with b = 1.5. Fig 6.8(c) shows chaotic phase portrait when k = 3.5.

- Case 8 Let $\alpha = \beta = 0.2$. Fig 6.9(a) shows the bifurcation diagram of the 0.4 order system. Fig 6.9(b) is the nonchaotic phase portrait of the system (2.6) with b = 5. Fig 6.9(c) shows chaotic phase portrait when k = 1.
- Case 9 Let $\alpha = \beta = 0.1$. Fig 6.10(a) shows the bifurcation diagram of the 0.2 order system. Fig 6.10(b) is the nonchaotic phase portrait of the system (2.6) with b = 8. Fig 6.10(c) shows chaotic phase portrait when k = 0.5.

Chaotizations in the integral and fractional order systems of a modified van der Pol systems are studied in the chapter. An efficient way to transform a non-chaotic dynamics of the system into a chaotic one is easily made by replacing a parameter of the system by *ksiny* where k is a adjustable constant and y is a chaotic state variable of a second system, a modified nano resonator system. It is found that chaos exists in the integral and fractional order systems with total order from 2 down to 0.2.



Fig 6.1(a) The bifurcation diagram for $\alpha = \beta = 1$. (b) The phase portrait for $\alpha = \beta = 1, b = 2.5$. (c) The phase portrait for $\alpha = \beta = 1, k = 4$.



Fig 6.2(a) The bifurcation diagram for $\alpha = \beta = 0.9$. (b) The phase portrait for $\alpha = \beta = 0.9$, b = 2.5. (c) The phase portrait for $\alpha = \beta = 0.9$, k = 4.



Fig 6.3(a) The bifurcation diagram for $\alpha = \beta = 0.8$. (b) The phase portrait for $\alpha = \beta = 0.8$, b = 1.5. (c) The phase portrait for $\alpha = \beta = 0.8$, k = 3.


Fig 6.4(a) The bifurcation diagram for $\alpha = \beta = 0.7$. (b) The phase portrait for $\alpha = \beta = 0.7$, b = 1.3. (c) The phase portrait for $\alpha = \beta = 0.7$, k = 3.5.



Fig 6.5(a) The bifurcation diagram for $\alpha = \beta = 0.6$. (b) The phase portrait for $\alpha = \beta = 0.6$, b = 1. (c) The phase portrait for $\alpha = \beta = 0.6$, k = 3.5.



Fig 6.6(a) The bifurcation diagram for $\alpha = \beta = 0.5$. (b) The phase portrait for $\alpha = \beta = 0.5$, b = 1. (c) The phase portrait for $\alpha = \beta = 0.5$, k = 3.



Fig 6.7(a) The bifurcation diagram for $\alpha = \beta = 0.4$. (b) The phase portrait for $\alpha = \beta = 0.4$, b = 1.5. (c) The phase portrait for $\alpha = \beta = 0.4$, k = 2.5.



Fig 6.8(a) The bifurcation diagram for $\alpha = \beta = 0.3$. (b) The phase portrait for $\alpha = \beta = 0.3$, b = 1.5. (c) The phase portrait for $\alpha = \beta = 0.3$, k = 3.5.



Fig 6.9(a) The bifurcation diagram for $\alpha = \beta = 0.2$. (b) The phase portrait for $\alpha = \beta = 0.2$, b = 5. (c) The phase portrait for $\alpha = \beta = 0.2$, k = 1.



Fig 6.10(a) The bifurcation diagram for $\alpha = \beta = 0.1$. (b) The phase portrait for $\alpha = \beta = 0.1$, b = 8. (c) The phase portrait for $\alpha = \beta = 0.1$, k = 0.5.

Chapter 7

Conclusions

Chaos is desirable under certain circumstances and chaotic phenomena are quite useful in many applications such as fluid mixing, human brain, and heart beat regulation, etc. Thus, this motivates us to investigate anticontrol of chaos. Recently many scientists in various fields have been attracted to investigate chaos synchronization due to its application in a variety of fields such as secure communications, chemical, physical, and biological systems, neural networks, etc. Many scholars have investigated the control and dynamics of the fractional order dynamical systems. The results have been shown that nonlinear chaotic systems still have chaotic behavior and chaos synchronization can be achieved when their models become fractional.

In this thesis, integral and fractional order modified van der Pol systems are investigated. Chapter 2 contains a fractional derivative and the model of modified van der Pol system.

In Chapter 3, the dynamics of integral and fractional order modified van der Pol systems are investigated. The chaotic phenomena exist for $0.8 \le (\alpha + \beta) \le 2.0$ and the range of the chaos gradually decreases as the total order number $\alpha + \beta$ decreases. The numerical results of periodic and chaotic phenomena are presented by phase portraits, Poincaré maps and bifurcation diagrams.

The synchronizations of two uncoupled integral and fractional order chaotic modified van der Pol systems are achieved by parameter excited synchronization in Chapter 4. It is found that this approach is very effective even for very low total fractional order 0.2.

Anticontrol of chaos of fractional order modified van der Pol systems by addition

of a constant term and addition of k|x|sinx term are discussed in Chapter 5. The results show that chaos exists in the fractional order systems with total order from 1.8 down to 0.6 for the addition of constant term, and from 1.8 down to 0.8 for the addition of k|x|sinx term.

Finally, the parameter excited chaotization of integral and fractional order modified van der Pol systems is studied in Chapter 6. It is found that chaotization can be successfully obtained in the integral and fractional order systems with total order from 2 down to 0.2.



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Appendix

Table 1. FRACTIONAL OPERATORS WITH APPROXIMATELY2 db ERROR FROM $\mathcal{O} = 10-2$ TO 102 rad/sec

$\frac{1}{100} \approx \frac{220.4s^4 + 5004s^3 + 503s^2 + 234.5s + 0.484}{5000000000000000000000000000000000000$
$s^{0.1}$ $s^3 + 359.8s^4 + 5742s^3 + 4247s^2 + 147.7s + 0.2099$
1 $60.95s^4 + 816.9s^3 + 582.8s^2 + 23.24s + 0.04934$
$\overline{s^{0.2}} \approx \overline{s^5 + 134s^4 + 956.5s^3 + 383.5s^2 + 8.953s + 0.01821}$
1 $23.76s^4 + 224.9s^3 + 129.1s^2 + 4.733s + 0.01052$
$\overline{s^{0.3}} \approx \overline{s^5 + 64.51s^4 + 252.2s^3 + 63.61s^2 + 1.104s + 0.002267}$
1 $25s^4 + 558.5s^3 + 664.2s^2 + 44.15s + 0.1562$
$\overline{s^{0.4}} \approx \overline{s^5 + 125.6s^4 + 840.6s^3 + 317.2s^2 + 7.428s + 0.02343}$
$1 \sim 15.97s^4 + 593.2s^3 + 1080s^2 + 135.4s + 1$
$\overline{s^{0.5}} \approx \overline{s^5 + 134.3s^4 + 1072s^3 + 543.4s^2 + 20.1s + 0.1259}$
$1 \qquad 8.579s^4 + 255.6s^3 + 405.3s^2 + 35.93s + 0.1696$
$\overline{s^{0.6}} \approx \overline{s^5 + 94.22s^4 + 472.9s^3 + 134.8s^2 + 2.639s + 0.009882}$
1 $4.406s^4 + 177.6s^3 + 209.6s^2 + 9.179s + 0.0145$
$\overline{s^{0.7}} \approx \overline{s^5 + 88.12s^4 + 279.2s^3 + 33.3s^2 + 1.927s + 0.0002276}$
1 $5.235s^3 + 1453s^2 + 5306s + 254.9$
$\overline{s^{0.8}} \approx \overline{s^4 + 658.1s^3 + 5700s^2 + 658.2s + 1}$
$1 1.766s^2 + 38.27s + 4.914$
$\frac{1}{s^{0.9}} \approx \frac{11000 + 50.278 + 10.911}{s^3 + 36.15s^2 + 7.789s + 0.01}$

Paper List

- [1] Ge Z.-M. and Zhang A.-R., "Chaos in a Modified Heartbeat System and in Its Fractional order Systems", Chaos, Solitons and Fractals, Dec. 2005, accepted and proofed.(SCI, Impact Factor:1.938)
- [2] Ge Z.-M. and Zhang A.-R., "Parameter Excited Chaos Synchronizations of Integral and Fractional Order Modified Heartbeat Systems", submitted to Chaos, Solitons and Fractals, 2006.
- [3] Ge Z.-M. and Zhang A.-R., "Anticontrol of Chaos of the Fractional Order Modified Heartbeat Systems", submitted to Chaos, Solitons and Fractals, 2006.
- [4] Ge Z.-M. and Zhang A.-R., "Parameter Excited Chaotization of Integral and Fractional Order Modified Heartbeat Systems", submitted to Chaos, Solitons and Fractals, 2006.

