Introduction

Chaotic systems have been a focal point of renewed interest for many researchers in the past few decades. Such nonlinear systems can occur in various natural and man-made systems, and are known to have great sensitivity to initial conditions [1].

Besides the theoretical interest in the analysis of such nonlinear systems, there is another dimension to that interest; namely, utilizing such systems for useful practical applications [2-13]. Many researchers have devoted themselves to finding new ways to control chaos more efficiently [14-17]. Chaotic phenomena are quite useful in many applications such as fluid mixing [18], human brain dynamics [19], and heart beat regulation [20], information processing, etc. Therefore, making a periodic dynamical system chaotic, or preserving chaos of a chaotic dynamical system, is very meaningful and worthy to be investigated [21,22].

Fractional calculus is a 300-year-old mathematical topic [23-26]. Although it has a long history, for many years it was not used in physics and engineering. However, during the last 10 years or so, fractional calculus starts to attract increasing attention of physicists and engineers from an application point of view [27,28]. It was found that many systems in interdisciplinary fields can be elegantly described with the help of fractional derivatives. Many systems are known to display fractional-order dynamics, such as viscoelastic systems [29], dielectric polarization [30], electrode–electrolyte polarization [31], electromagnetic waves [32], quantitative finance [33], and quantum evolution of complex systems [34].

There are few investigations about delay system with fractional order. In this proposal, chaos in new integral and fractional order double Ikeda delay systems is

studied. A double Ikeda delay system consists of two traditional Ikeda delay systems which are coupled together. Numerical simulations display the chaotic behaviors of the integral and fractional order delay systems by phase portraits, Poincaré maps and bifurcation diagrams.

Since the pioneering work was given by Pecorra and Carroll [35], chaos synchronization [36-41] has become an important topic in engineering science. Many effective control schemes have been developed in a variety of fields, such as parameters adaptive control [42-49], observer-based control [50, 51], variable structure control [52, 53], active control [54-58], anti-control [59-63], nonlinear control [64-66] and so on. The applications of chaos synchronization are implemented extensively in secure communications, chemical, physical, and biological systems and neural networks.

Recently, the concept of synchronization has been extended the scope to, such as generalized synchronization, lag and anticipated synchronization, phase synchronization and anti-synchronization. The basic synchronization called complete synchronization is that the state vectors of the first system x(t) is equal to the state vectors of the second system y(t): y(t) = x(t). The lag synchronization [67] is that the state vector of the second system y(t) delays that of driver system x: y(t) = x(t-T) with positive T. If T is negative, we have anticipated synchronization. Lag anti-synchronization [68] means y(t) = -x(t-T). When T is negative, we have anticipated anti-synchronization.

Time-delayed systems are ubiquitous in nature, technology, and society because of finite signal transmission times, switching speeds, and memory effects [69]. Therefore the study of chaos synchronization in these systems is of considerable practical significance. It is well known that dissipative systems with a nonlinear time-delayed feedback or "memory" can produce chaotic dynamics, and the

dimension of their chaotic attractors can be made arbitrarily large by increasing their delay time sufficiently. Thus time-delay systems are good candidates for secure communications based on chaos synchronization [70].

This thesis is organized as follows. Chapter 2 gives the dynamic equation of double Ikeda system. The fractional derivative and its approximation are introduced. The system under study is described both in its integer and fractional forms. Numerical simulation results are presented.

In Chapter 3, it is discovered that lag synchronization and lag anti-synchronization appear for two identical double Ikeda systems, without any control scheme or coupling terms, but with different initial conditions.

In Chapter 4, the chaotic behaviors of double Ikeda systems are obtained by replacing the original constant delay time by a function of chaotic state variable of a second chaotic double Ikeda system. Numerical simulations are illustrated by phase portraits. Phase portrait is expressed by numerical analysis.

In Chapter 5, the chaotic behaviors of double Ikeda systems are obtained by replacing the parameters by different chaotic state variables of a third chaotic double Ikeda system. The method is named parameter excited method for synchronization which will be successfully used for uncoupled synchronization. Numerical simulations are illustrated by phase portraits and time histories.

In Chapter 6, conclusions are drawn.

Chaos in Integral and Fractional Order Double Ikeda Systems

2.1 Preliminaries

is

In this Chapter, the chaotic behaviors in integral and fractional order double Ikeda systems are studied numerically by phase portraits, Poincaré maps and bifurcation diagrams. It is found that chaos exists for all systems with total orders of derivatives from 2 to 0.2.

2.2 Fractional Derivative and Its Approximation

The idea of fractional integrals and derivatives has been known since the development of the regular calculus, with the first reference probably being associated with Leibniz in 1695 [71].

Two commonly used definitions for the general fractional differintegral are the Grunwald definition and the Riemann-Liouville definition. The latter is given here

$$\frac{d^{q} f(t)}{dt^{q}} = \frac{1}{\Gamma(n-q)} \frac{d^{n}}{dt^{n}} \int_{0}^{t} \frac{f(\tau)}{(t-\tau)^{q-n+1}} d\tau$$
 (2.1)

where $n-1 \le q < n$ and $\Gamma(\cdot)$ is an Euler's gamma function.

The Laplace transformation of the Riemann-Liouville fractional derivative (2.1)

$$L\left\{\frac{d^{q} f(t)}{dt^{q}}\right\} = s^{q} L\left\{f(t)\right\} - \sum_{k=0}^{n-1} s^{k} \left[\frac{d^{q-1-k} f(t)}{dt^{q-1-k}}\right]_{t=0}, \text{ for } n-1 \le q < n$$
 (2.2)

By considering the initial conditions to be zero, this formula reduces to the more

expected and comforting form

$$L\left\{\frac{d^{q}f(t)}{dt^{q}}\right\} = s^{q}L\left\{f(t)\right\} \tag{2.3}$$

and the fractional integral of order q can be described as $F(s) = \frac{1}{s^q}$ in the frequency domain.

The standard definitions of the fractional differintegral do not allow direct implementation of the operator in time domain simulations of complicated systems with fractional elements. Using the standard integer order operators to approximate the fractional operators is an effective method to analyze such systems.

The approximation approach taken here is to approximate the system behavior in the frequency domain [72]. By utilizing frequency domain techniques based in Bode diagrams, one can obtain a linear approximation of a fractional order integrator. Thus an approximation of any desired accuracy over any frequency band can be achieved. Table I of Ref. [73] gives approximations for $\frac{1}{s^q}$ with $q = 0.1 \sim 0.9$ in steps of 0.1 with errors of approximately 2 dB from $\omega = 10^{-2}$ to 10^2 rad/s. These approximations will be used in the following numerical simulations.

2.3 Integral and Fractional Order Double Ikeda Systems

The traditional Ikeda system is

$$\dot{x} = -ax - b\sin x_{\tau} \tag{2.4}$$

where $x_{\tau} \equiv x(t-\tau)$ with positive τ . x is the phase lag of the electric field across the resonator, and $x(t-\tau)$ is the round trip time of the light in the resonator or feedback delay time in the coupled systems.

In this paper, we consider a new delay system which consists of two coupled Ikeda systems, called double Ikeda system:

$$\begin{cases} \dot{x} = -a_1 x - b_1 \sin x_{\tau_1} - c \sin y \\ \dot{y} = -a_2 y - b_2 \sin y_{\tau_2} \end{cases}$$
 (2.5)

where a_1 , a_2 , b_1 , b_2 and c are constants, and the delay times of x and y can be represented as τ_1 and τ_2 , respectively. The chaotic behaviors are also changed when the time delays are varied.

Now, consider a fractional order modified Double Ikeda system. Here, the conventional derivatives in Eq.(2.5) are replaced by the fractional derivatives as follows:

$$\begin{cases} \frac{d^{q_1}x}{dt^{q_1}} = -a_1x - b_1\sin x_{\tau_1} - c\sin y\\ \frac{d^{q_2}y}{dt^{q_2}} = -a_2y - b_2\sin y_{\tau_2} \end{cases}$$
 (2.6)

where q_1, q_2 are two fractional order numbers.

In this paper, we analyse and present simulation results of the chaotic dynamics produced from a new fractional Double Ikeda system as the fractional order derivatives q_1, q_2 in the state equations of Eq. (2.6) is varied from 1 to 0.1.

2.4 Simulation Results

We vary the derivative orders q_1, q_2 and the system parameter c. Other system parameters are fixed, which are given as: $a_1 = 1$, $a_2 = 15$, $b_1 = b_2 = 20$, $\tau_1 = \tau_2 = 2$. The numerical simulations are carried out by MATLAB with using the fractional operator in table I of [73] in the Simulink environment.

The phase portraits, Poincaré maps and bifurcation diagrams of the systems with total order of derivatives from 2~0.2 are showed in Fig.2.1~Fig.2.10. Chaos exists for all cases. The period 2 phase portraits in Fig.2.1(b), Fig.2.2(b), Fig.2.3(b) seems contradictory with the "four lines" bifurcations in Fig.2.1(d), Fig.2.2(d), Fig.2.3(d), respectively. Actually, the "four lines" are dotted lines, the corresponding two dotted

lines form a undotted line. Therefore period 2 motions also appear in Fig.2.1(d), Fig.2.2(d), Fig.2.3(d).



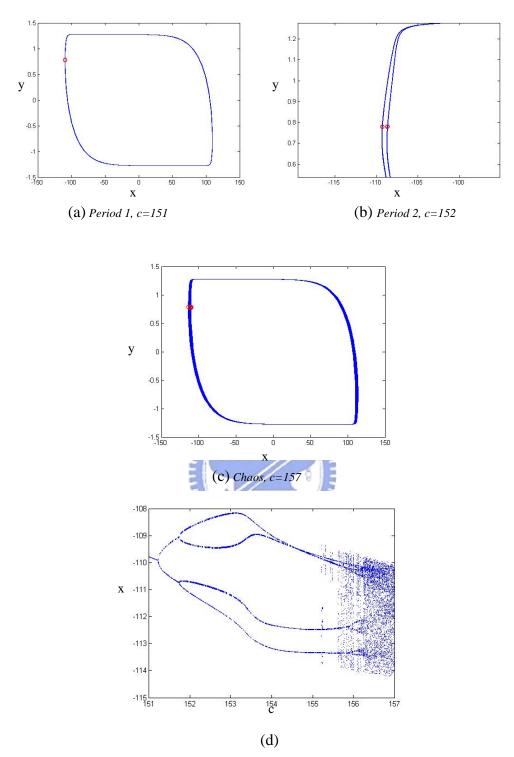


Fig. 2.1 The phase portraits and the bifurcation diagram for double Ikeda system with order $q_1 = 1$ and $q_2 = 1$. The first bifurcation point c=151.2. The second bifurcation point c=151.7. The third bifurcation point c=155.7.

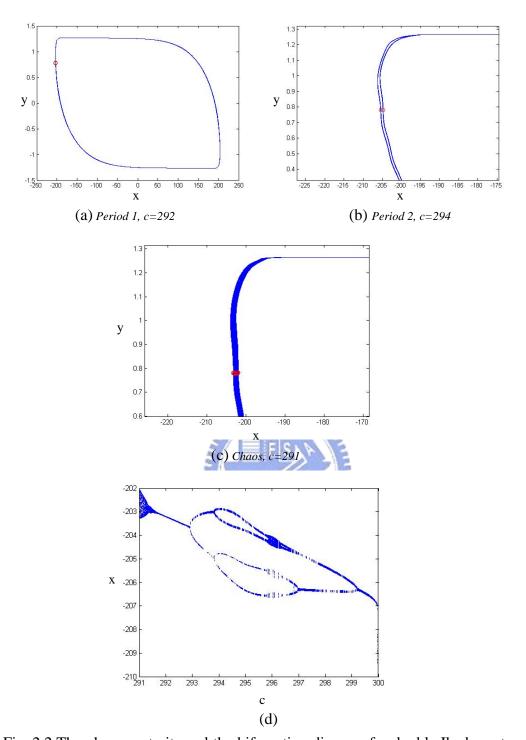


Fig. 2.2 The phase portraits and the bifurcation diagram for double Ikeda system with order $q_1 = 0.9$ and $q_2 = 0.9$. The first bifurcation point c=292.9. The second bifurcation point c=293.8. The third bifurcation point c=297.1. The fourth bifurcation point c=299.2.

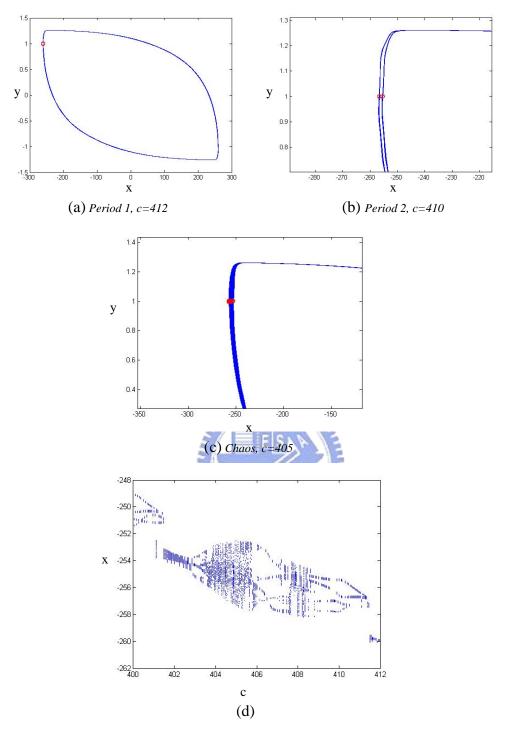


Fig. 2.3 The phase portraits and the bifurcation diagram for double Ikeda system with order $q_1 = 0.8$ and $q_2 = 0.8$. The first bifurcation point c=411.4.

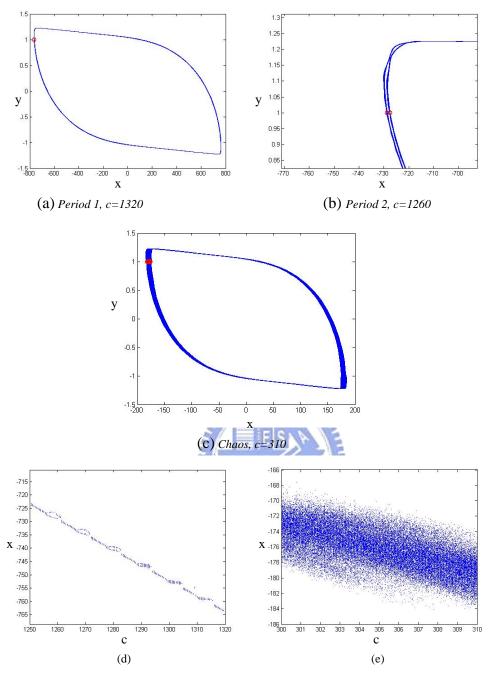


Fig. 2.4 The phase portraits and the bifurcation diagram for double Ikeda system with order $q_1=0.7$ and $q_2=0.7$. The first bifurcation point c=1256. The second bifurcation point c=512.

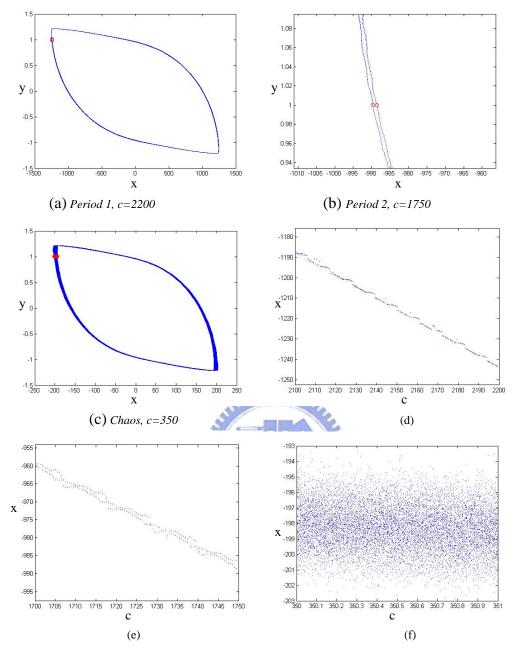


Fig. 2.5 The phase portraits and the bifurcation diagram for double Ikeda system with order $q_1 = 0.6$ and $q_2 = 0.6$. The first bifurcation point c=1821. The second bifurcation point c=437.

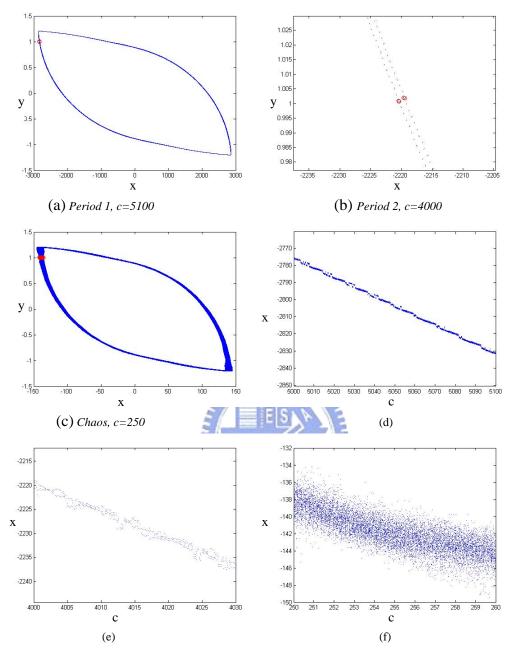


Fig. 2.6 The phase portraits and the bifurcation diagram for double Ikeda system with order $q_1=0.5$ and $q_2=0.5$. The first bifurcation point c=4316. The second bifurcation point c=334.

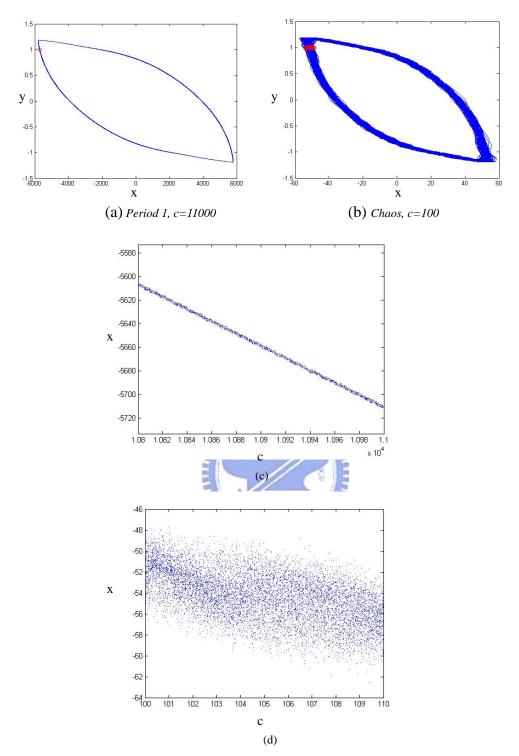


Fig. 2.7 The phase portraits and the bifurcation diagram for double Ikeda system with order $q_1 = 0.4$ and $q_2 = 0.4$. The first bifurcation point c=211.

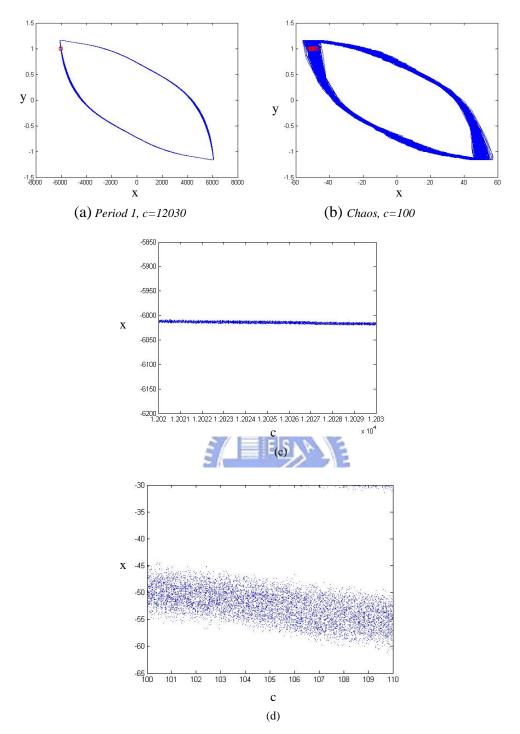


Fig. 2.8 The phase portraits and the bifurcation diagram for double Ikeda system with order $q_1 = 0.3$ and $q_2 = 0.3$. The first bifurcation point c=225.

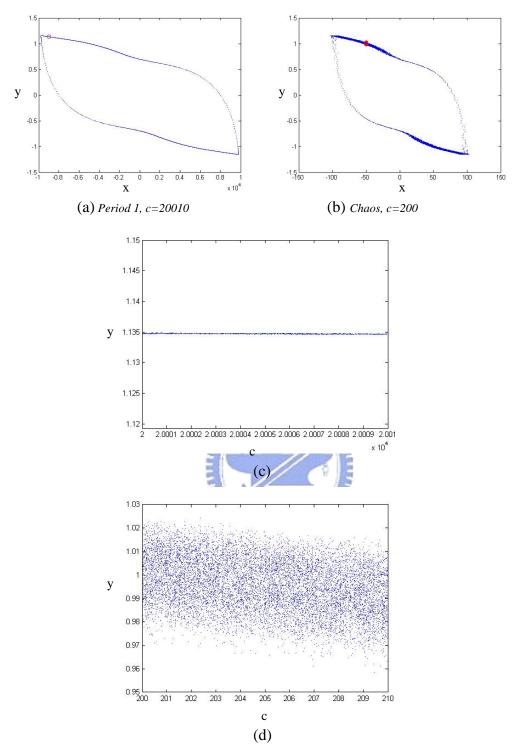


Fig. 2.9 The phase portraits and the bifurcation diagram for double Ikeda system with order $q_1 = 0.2$ and $q_2 = 0.2$. The first bifurcation point c=230.

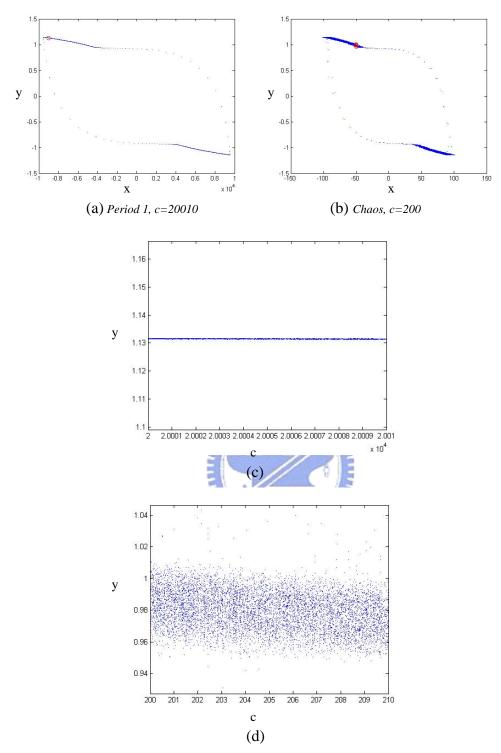


Fig. 2.10 The phase portraits and the bifurcation diagram for double Ikeda system with order $q_1 = 0.1$ and $q_2 = 0.1$. The first bifurcation point c=232.

Lag and Anticipated Synchronization and Anti-synchronization of Two Uncoupled Time-delayed Chaotic Systems

3.1 Preliminaries

Lag and anticipated synchronization and lag and anticipated anti-synchronization are newly discovered in two identical double Ikeda systems with different initial conditions without any control scheme and coupling terms. There are two situations for all possible initial conditions. They are the lag or anticipated synchronization, the lag or anticipated anti-synchronization.

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3.2 Lag or anticipated synchronization and lag or anticipated anti-synchronization

Consider the first time-delay chaotic system

$$\dot{x} = f(x, x_{\tau}, t) \tag{3.1}$$

and second time-delay chaotic system

$$\dot{y} = f(y, y_{\tau}, t) \tag{3.2}$$

where $x, y \in \mathbb{R}^n$ are *n*-dimensional state vectors, $x_\tau = x(t-\tau)$ are corresponding time-delay state vectors. The error are defined as e = x(t-T) - y(t). If the following conditions hold, the systems are in lag synchronization.

$$e_i = x_{iT_i} - y_i = 0, \quad i = 1, 2, ..., p \le n$$
 (3.3)

where x_i , y_i are the state vectors of the system, T_i is the time which x_i lag behind y_i . When T_i is negative, we have anticipated synchronization.

In the case of anti-synchronization, the states of the systems which have opposite signs, the error e = x(t-T) + y(t) will converge to zero. Therefore, we can say the lag anti-synchronization is achieved when the following conditions are satisfied:

$$e_i = x_{iT_i} + y_i = 0, \quad i = 1, 2, ..., p \le n$$
 (3.4)

where x_i , y_i are the state vectors of the system, T_i is the time which x_i lag behind y_i . When T_i is negative, we have anticipated anti-synchronization.

3.3 The lag or anticipated synchronization of two identical double Ikeda systems

We consider two double Ikeda systems which consist of two coupled Ikeda equations:

System A

$$\begin{cases} \dot{x}_1 = -a_1 x_1 - b_1 \sin x_{1\tau_1} - c \sin y_1 \\ \dot{y}_1 = -a_2 y_1 - b_2 \sin y_{1\tau_2} \end{cases}$$
(3.5)

and system B

$$\begin{cases} \dot{x}_2 = -a_1 x_2 - b_1 \sin x_{2\tau_1} - c \sin y_2 \\ \dot{y}_2 = -a_2 y_2 - b_2 \sin y_{2\tau_2} \end{cases}$$
(3.6)

with positive $a_{1,2}$ and $b_{1,2}$

This investigation is of considerable practical importance, as the equations of the class B lasers with feedback (typical representatives of class B are solid-state, semiconductor, and low pressure CO2 lasers [74]) can be reduced to an equation of

the Ikeda type [75].

The Ikeda model was introduced to describe the dynamics of an optical bistable resonator, plays an important role in electronics and physiological studies and is well-known for delay-induced chaotic behavior [76,77,78,79]. Physically x is the phase lag of the electric field across the resonator; a is the relaxation coefficient for the driving x and driven y dynamical variables; $b_{1,2}$ are the laser intensities injected into the driving and driven systems, respectively. $\tau_{1,2}$ are the feedback delay times in the coupled systems.

We keep the delay time fixed in 2 second ($\tau = 2$) and the parameters are shown as follows: $a_1 = 1$, $a_2 = 15$, $b_1 = 20$, $b_2 = 20$, c = 157. The system is chaotic in foregoing conditions as shown in Fig. 3.1. The numerical simulations are carried out by MATLAB, The initial conditions we choose are constant.

Fig. 3.2 shows the time histories of double Ikeda system with initial conditions as follows in Table 1. Fig. 3.2 (b) is the magnified diagram of Fig. 3.2 (a). After 180 sec, state variables of system A become in lag synchronization with that of the system B. Lag of x_1 to x_2 is 0.14868 sec. And lag of y_1 to y_2 is 0.14868 sec. The situation remain unchanged until 20000 sec. From Fig. 3.2, we can see the lag synchronization in Case 1~8 by Table 3.1. From Fig. 3.2 (b), state variables of system A is in lag synchronization with that of system B, or state variables of system B is in anticipated synchronization with that of system A.

Table 3.1 The initial condition of lag or anticipated synchronization.

Case	x_1	\mathcal{Y}_1	X_2	\mathcal{Y}_2
1	1	1	0.000000001	0.000000001
2	1	1	-0.000000001	0.000000001

3	1	-1	0.000000001	-0.000000001
4	1	-1	-0.000000001	-0.000000001
5	-1	1	0.000000001	0.000000001
6	-1	1	-0.000000001	0.000000001
7	-1	-1	0.000000001	-0.000000001
8	-1	-1	-0.000000001	-0.000000001

3.4 The lag or anticipated anti-synchronization of two identical double Ikeda systems

In this section, we change the initial conditions. The lag anti-synchronization is occurred. Fig.3.3 shows the time histories of double Ikeda system with initial conditions as follows in Table 3.2. Fig.3.3 (b) is the magnified diagram of Fig.3.3 (a). Fig.3.3 (c) is the time history of $-x_1$ (blue) and x_2 (red), $-y_1$ (blue) and y_2 (red) of double Ikeda systems. Lag of x_1 to x_2 is 0.14868 sec. And lag of y_1 to y_2 is 0.14868 sec. The situation remain unchanged until 20000 sec. From Fig. 3.3, we can see the lag synchronization in Case 9~16 by Table 2. From Fig.3.3 (b), state variables of system A is in lag anti-synchronization with that of system B, or state variables of system B is in anticipated anti-synchronization with that of system A.

Table 3.2 The initial condition of lag or anticipated anti-synchronization.

Case	x_1	y_1	X_2	y_2
9	1	1	0.000000001	-0.000000001
10	1	1	-0.000000001	-0.000000001
11	1	-1	0.000000001	0.000000001
12	1	-1	-0.000000001	0.000000001

13	-1	1	0.000000001	-0.000000001
14	-1	1	-0.000000001	-0.000000001
15	-1	-1	0.000000001	0.000000001
16	-1	-1	-0.000000001	0.000000001



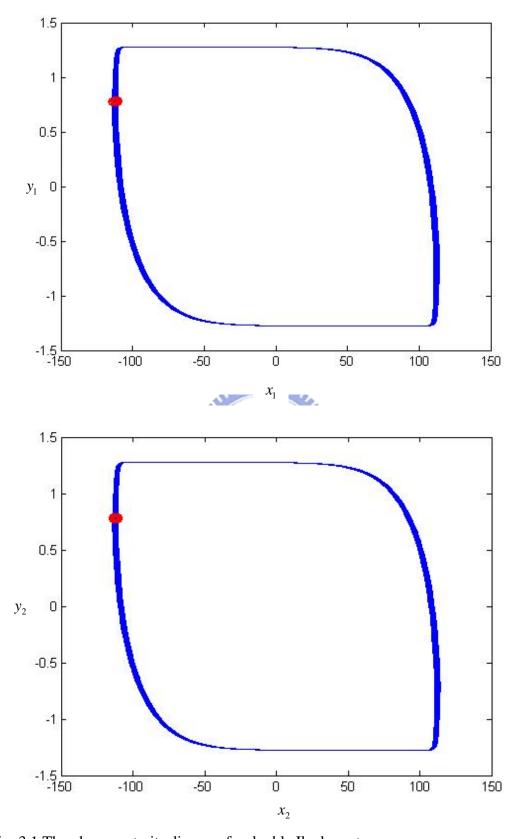


Fig. 3.1 The phase portraits diagram for double Ikeda system.

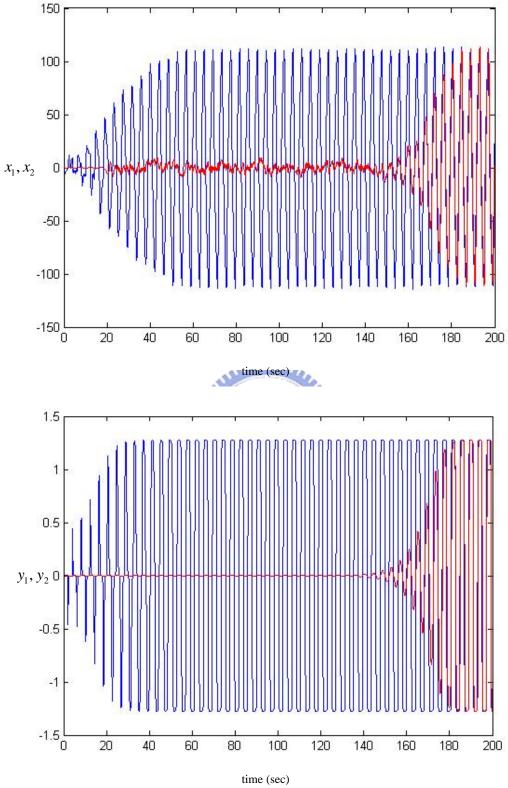


Fig.3.2 (a) The time histories of x_1 (blue) and x_2 (red), y_1 (blue) and y_2 (red) of double Ikeda systems with initial conditions in Table 3.1.

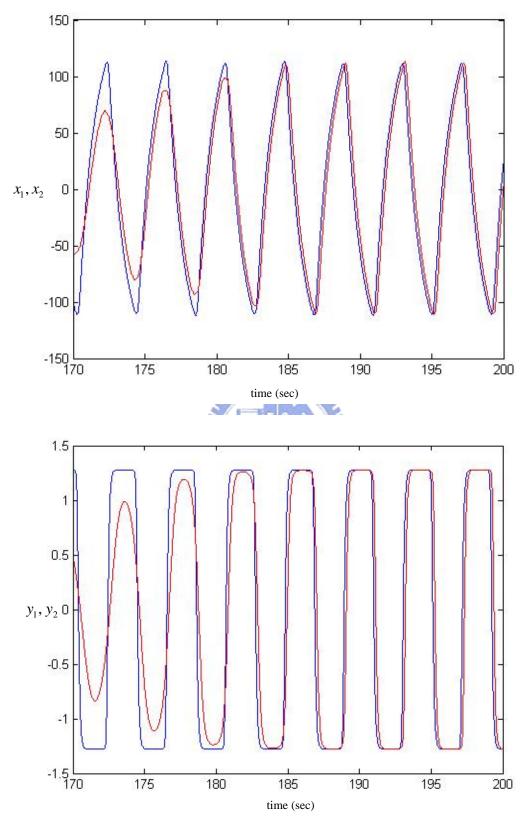


Fig.3.2 (b) The time histories of x_1 (blue) and x_2 (red), y_1 (blue) and y_2 (red) of double Ikeda systems with initial conditions in Table 3.1.

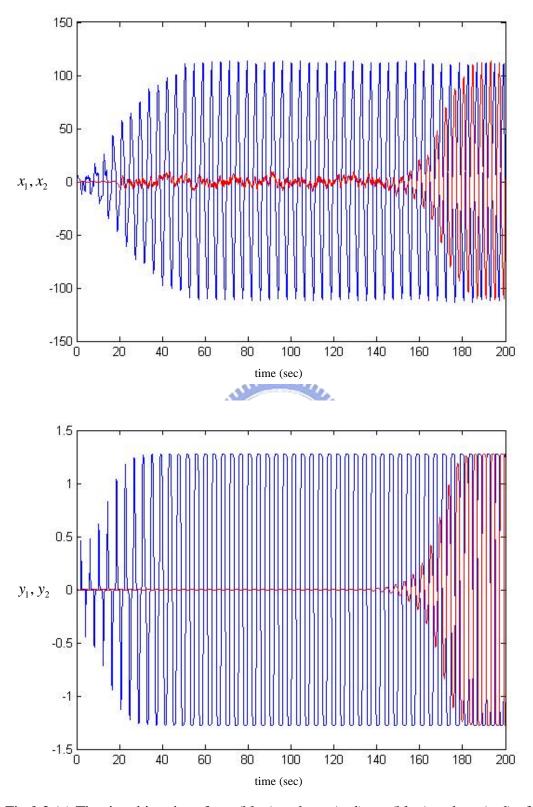


Fig.3.3 (a) The time histories of x_1 (blue) and x_2 (red), y_1 (blue) and y_2 (red) of double Ikeda systems with initial conditions in Table 3.2.

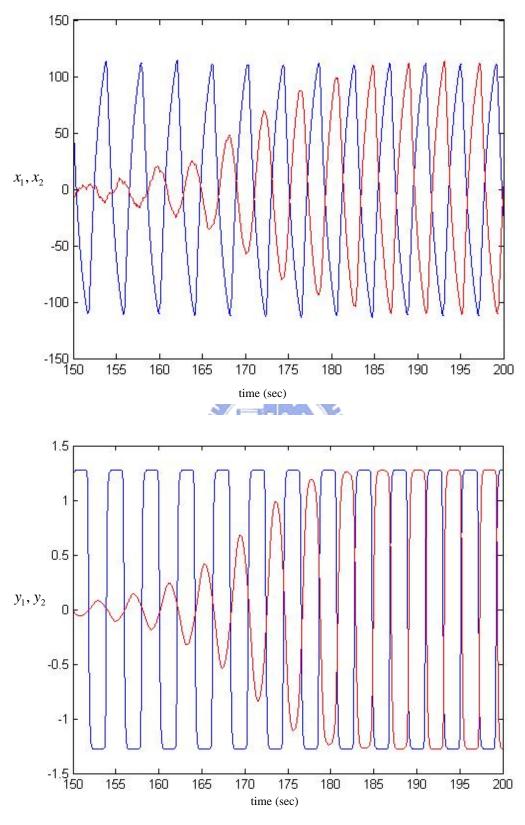


Fig.3.3 (b) The time histories of x_1 (blue) and x_2 (red), y_1 (blue) and y_2 (red) of double Ikeda systems with initial conditions in Table 3.2.

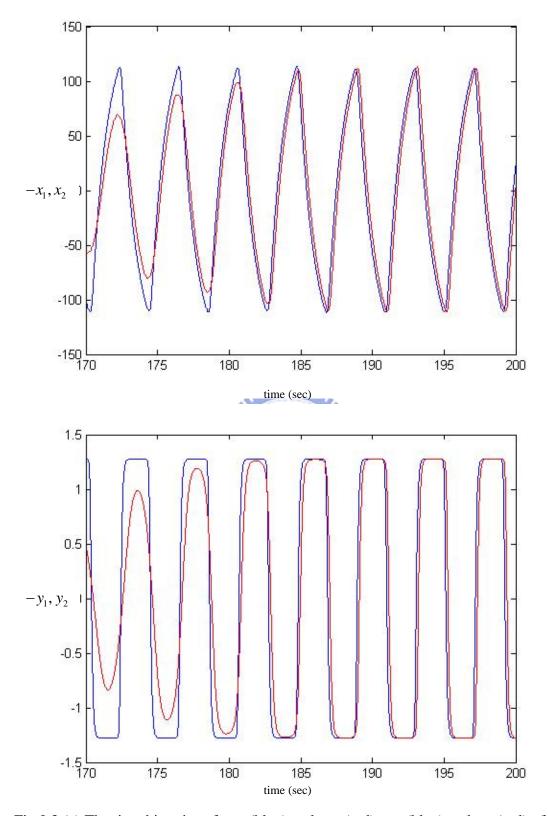


Fig.3.3 (c) The time histories of $-x_1$ (blue) and x_2 (red), $-y_1$ (blue) and y_2 (red) of double Ikeda systems with initial conditions in Table 3.2.

Chaos and Chaotization of a Double Ikeda System by Chaotic Delay Time

4.1 Preliminaries

In this Chapter, the chaotization of a double Ikeda delay system by replacing the original constant delay time by a function of chaotic state variable of a second double Ikeda chaotic system. It is found that chaos exists for many cases.

4.2 Chaos of a double Ikeda system....

According to Chapter 3, we consider two double Ikeda systems which consist of two coupled Ikeda equations:

$$\begin{cases} \dot{x}_1 = -a_1 x_1 - b_1 \sin x_{1\tau_1} - c \sin y_1 \\ \dot{y}_1 = -a_2 y_1 - b_2 \sin y_{1\tau_2} \end{cases}$$
(4.1)

and

$$\begin{cases} \dot{x}_2 = -a_1 x_2 - b_1 \sin x_{2\tau_1} - d \sin y_2 \\ \dot{y}_2 = -a_2 y_2 - b_2 \sin y_{2\tau_2} \end{cases}$$
(4.2)

The scheme is to replace the constant delay time τ_1 or τ_2 in system (4.1) by a function of a chaotic state of system (4.2).

4.3 Chaotization scheme of a double Ikeda system by chaotic delay time

Creating chaos is called chaotization at times. The above replacements enhance the existing chaos of the originally system effectively. The results are demonstrated by numerical results.

In order to induce chaotic phenomena of the double Ikeda system (4.1), $p+qx_2$ and $p+qy_2$ replace the delay times in system (4.1) respectively, where p, q are constant.

4.4 Numerical Simulations for Chaos and Chaotization by time delay driven by a chaotic signal

In following simulations, we replace the delay time in system (4.1) by x_2 where x_2 is a state variable in system (4.2). In our numerical simulations, $a_1 = 1$, $a_2 = 15$, $b_1 = 20$, $b_2 = 20$, c = 157, $\tau_1 = \tau_2 = 2s$ of system (4.1) and (4.2) are fixed. The initial states of system (4.1) and (4.2) are $x_1(0) = -1$, $y_1(0) = -1$, $x_2(0) = -1.1$, $y_2(0) = -1.1$. The numerical simulations are carried out by MATLAB.

Case 1: The parameters $a_1 = 1$, $a_2 = 15$, $b_1 = 20$, $b_2 = 20$, c = 157, d = 157 of system (4.1) and (4.2) are fixed. The delay time τ_1 of system (4.1) is replaced by $200 + x_2$, where x_2 is the state variable of system (4.2). The phase portraits are shown in Fig. 4.1. The phase portrait with the original constant delay time is shown in Fig. 4.1(b)

Case 2: The parameters $a_1 = 1$, $a_2 = 15$, $b_1 = 20$, $b_2 = 20$, c = 157, d = 157 of system (4.1) and (4.2) are fixed. The delay time τ_2 of system (4.1) is replaced by $200 + x_2$, where x_2 is the state variable of system (4.2). The phase portraits are shown in Fig. 4.2.

Case $3\sim4$: The delay times are changed by $200+0.5x_2$, the sequences which are similar to Case $1\sim2$. The phase portraits are shown in Fig. 4.3 ~4.4 .

Case $5\sim6$: The delay times are changed by $200+0.612x_2$, the sequences which are similar to Case $1\sim2$. The phase portraits are shown in Fig. $4.5\sim4.6$.

In Case 2 and 4, it is not found that chaos exists. The time history are not conform the chaos. The time histories are shown in Fig. 4.7~4.8.

For the chaotization, we replace the delay time in system (4.1) by a function of chaotic states of system (4.2). The system (4.1) is periodic and the system (4.2) is chaotic.

Case 7: The parameters $a_1 = 1$, $a_2 = 15$, $b_1 = 20$, $b_2 = 20$, c = 151 of system (4.1) is fixed. And the parameters d = 157 of system (4.2), the others are the same .The delay time τ_1 of system (4.1) is replaced by $200 + 0.612x_2$, where x_2 is the state variable of system (4.2). The phase portraits are shown in Fig. 4.9. The phase portrait that the state is periodic is shown in Fig. 4.9(b)

Case 8: The parameters $a_1 = 1$, $a_2 = 15$, $b_1 = 20$, $b_2 = 20$, c = 151 of system (4.1) is fixed. And the parameters d = 157 of system (4.2), the others are the same .The delay time τ_2 of system (4.1) is replaced by $200 + 0.612x_2$, where x_2 is the state variable of system (4.2). The phase portraits are shown in Fig. 4.10.

Case $9\sim10$: The delay times are changed by $200+0.1y_2$, the sequences which are similar to Case $7\sim8$. The phase portraits are shown in Fig. 4.11 ~4.12 .

In Case 10, no chaotization exists. The time histories are shown in Fig. 4.13.

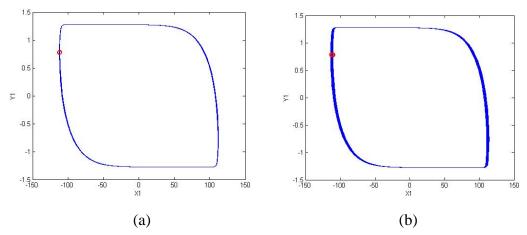


Fig. 4.1 (a) The phase portrait, the delay time τ_1 of system (4.1) is replaced by $200 + x_2$.

Fig. 4.1 (b) The phase portrait, the original constant delay time is fixed.

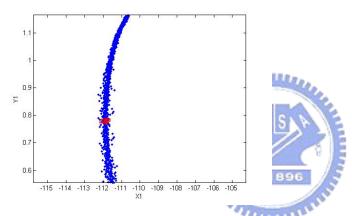


Fig. 4.1 (c) It is the magnified diagram of Fig.4.1 (a).

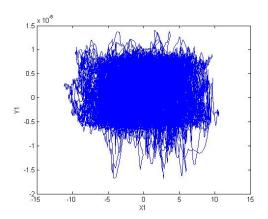


Fig. 4.2 The phase portrait, the delay time $\ \tau_2$ of system (4.1) is replaced by $200+x_2$.

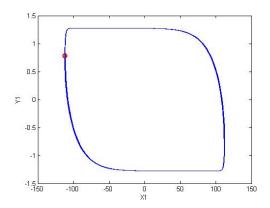


Fig. 4.3 (a) The phase portrait, the delay time τ_1 of system (4.1) is replaced by $200+0.5x_2$.

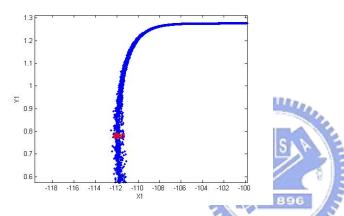


Fig. 4.3 (b) It is the magnified diagram of Fig. 4.3 (a).

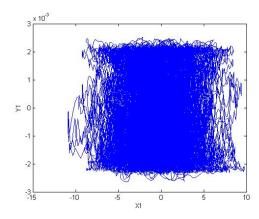


Fig. 4.4 The phase portrait, the delay time τ_2 of system (4.1) is replaced by $200+0.5x_2$.

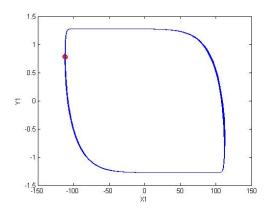


Fig. 4.5 (a) The phase portrait, the delay time τ_1 of system (4.1) is replaced by $200+0.612x_2$.

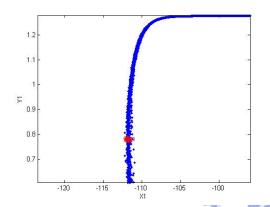


Fig. 4.5 (b) It is the magnified diagram of Fig.5 (a).

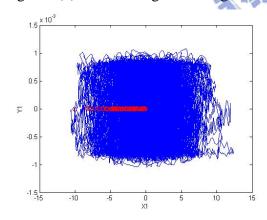


Fig. 4.6 The phase portrait, the delay time τ_2 of system (4.1) is replaced by $200+0.612x_2$.

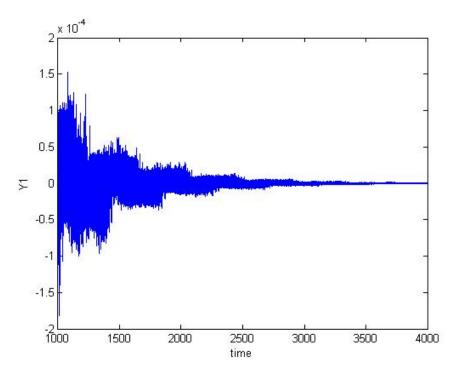


Fig. 4.7 The time history when the delay time τ_2 of system (4.1) is replaced by $200 + x_2$.

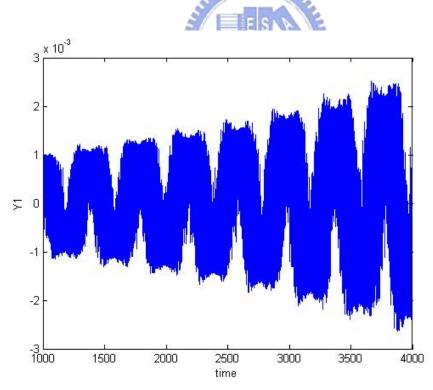


Fig. 4.8 The time history when the delay time τ_2 of system (4.1) is replaced by $200+0.5x_2$.

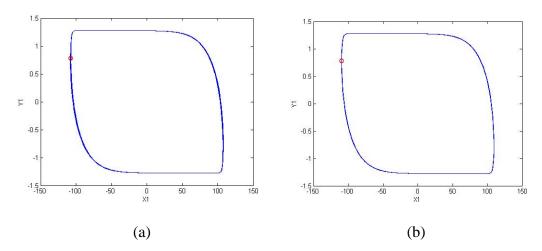


Fig. 4.9 (a) The phase portrait, the delay time τ_1 of system (4.1) is replaced by $200+0.612x_2$.

Fig. 4.9 (b) The phase portrait, the state is periodic.

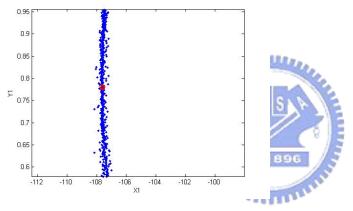


Fig. 4.9 (c) It is the magnified diagram of $\overline{\text{Fig.4.9}}$ (a).

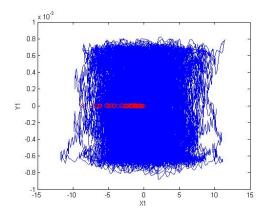


Fig. 4.10 The phase portrait, the delay time τ_2 of system (4.1) is replaced by $200+0.612x_2$.

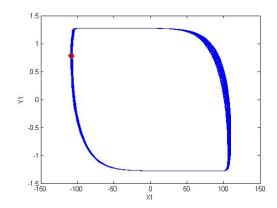


Fig. 4.11 (a) The phase portrait, the delay time τ_1 of system (4.1) is replaced by $200+0.1y_2$.

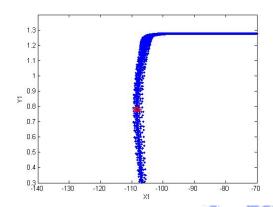


Fig. 4.11 (b) It is the magnified diagram of Fig.4.11 (a).

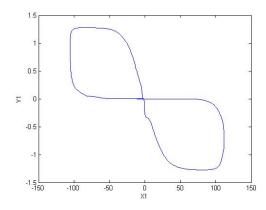


Fig. 4.12 The phase portrait, the delay time τ_2 of system (4.1) is replaced by $200+0.1y_2$.

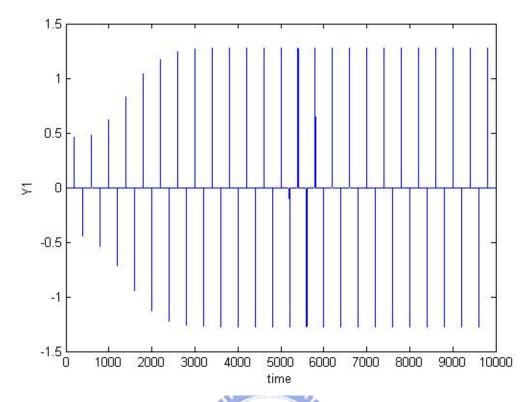


Fig. 4.13 The time history, the delay time τ_2 of system (4.1) is replaced by $200+0.1y_2$.

Chapter 5

Robust Lag Chaos Synchronization, Lag Chaos Quasi-Synchronization and Chaos Control of Double Ikeda System by Uncoupled Parameter Excited Method

5.1 Preliminaries

In this chapter, lag synchronization, lag quasi-synchronization and chaos control of two uncoupled double Ikeda systems, are achieved by replacing the corresponding parameters of two systems by the chaotic state variables of a third chaotic system. It is named as parameter excited method for synchronization. Numerical simulations are illustrated by phase portraits, Poincaré maps and time histories.

5.2 Lag Chaos Synchronization, Lag Chaos Quasi-Synchronization and Chaos Control with Parameters Replaced by Different Chaotic Signals

In this section, consider three identical double Ikeda systems:

$$\begin{cases} \dot{x}_1 = -a_1 x_1 - b_1 \sin x_{1\tau_1} - c \sin y_1 \\ \dot{y}_1 = -a_2 y_1 - b_2 \sin y_{1\tau_2} \end{cases}$$
 (5.1)

and

$$\begin{cases} \dot{x}_2 = -a_1 x_2 - b_1 \sin x_{2\tau_1} - c \sin y_2 \\ \dot{y}_2 = -a_2 y_2 - b_2 \sin y_{2\tau_2} \end{cases}$$
 (5.2)

and

$$\begin{cases} \dot{x}_3 = -a_1 x_3 - b_1 \sin x_{3\tau_1} - c \sin y_3 \\ \dot{y}_3 = -a_2 y_3 - b_2 \sin y_{3\tau_2} \end{cases}$$
 (5.3)

The scheme is to replace the parameters b_1 , b_2 in system (5.1) and (5.2) by two different chaotic states of system (5.3), respectively. One replaces a parameter in system (5.1) by a chaotic state of system (5.3), and replaces the corresponding parameter in system (5.2) by the same chaotic state of system (5.3) with a delay time τ_3 .

In following simulations, the parameter in system (5.1) is replaced by x_3 where x_3 is a state variable in system (5.3). The parameter in system (5.2) is replaced by $x_3(t-\tau_3)=x_{3_{\tau_3}}$ where τ_3 is a constant. In our numerical simulations, $a_1=1$, $a_2=15$, $b_1=20$, $b_2=20$, c=157, $\tau_1=\tau_2=2s$ of system (5.1), (5.2) and (5.3) are fixed. The initial states of system (5.1), (5.2) and (5.3) are $x_1(0)=-1$, $y_1(0)=-1$, $x_2(0)=-1.1$, $y_2(0)=-1.1$, $x_3(0)=-1$. Chaotic phase portraits are shown in Fig. 5.1. The numerical simulations are carried out by MATLAB.

The parameter b_1 of system (5.1) is replaced by x_3 , and the parameter b_1 of system (5.2) is replaced by $x_{3_{r_3}}$. τ_3 =10, 50, 100 are used in simulation. The phase portraits, Poincaré maps and time histories are shown in Figs. 5.2~5.4. Fig.5.3 (d) is the magnified diagram of Fig.5.3 (b). Fig.5.3 (e) is the magnified diagram of Fig.5.4 (d) is the magnified diagram of Fig.5.4 (e) is the magnified diagram of Fig.5.4 (c).

In these cases, we can find the lag quasi-synchronization between x_1 and x_2 , and lag synchronization between y_1 and y_2 .

Then the parameter b_2 of system (5.1) is replaced by y_3 , and the parameter b_2 of system (5.2) is replaced by $y_{3_{r_3}}$, where τ_3 =100. The phase portraits are shown in Fig. 5.5. In this case, by parameter excited method we successfully control x_1 , x_2 to constants, and y_1 , y_2 to zero.

Parameter excited method for lag synchronization by replacements of a_1 , a_2 and c has not been found effective.

5.3 Robustness of Lag Chaos Synchronization

In the section, noises are added to systems (5.1) and (5.2) as follows:

$$\begin{cases} \dot{x}_1 = -a_1 x_1 - b_1 \sin x_{1\tau_1} - c \sin y_1 + k N_1 \\ \dot{y}_1 = -a_2 y_1 - b_2 \sin y_{1\tau_2} \end{cases}$$
(5.4)

and

$$\begin{cases} \dot{x}_2 = -a_1 x_2 - b_1 \sin x_{2\tau_1} - c \sin y_2 + k N_2 \\ \dot{y}_2 = -a_2 y_2 - b_2 \sin y_{2\tau_2} \end{cases}$$
 (5.5)

where k is strength constant, and N_1, N_2 are different kinds of noise.

Three kind of noise are used. There are Gaussian noise, Rayleigh noise, and Rician noise.

The probability density function of n-dimensional Gaussian noise is

$$f(x) = ((2\pi)^n \det K)^{-\frac{1}{2}} \exp(-(x-\mu)^T K^{-1}(x-\mu)/2)$$
(5.6)

where x is a length-n vector, K is the n-by-n covariance matrix, μ is the mean value vector, and the superscript T indicates matrix transpose. The Simulink Communications toolbox provides the Gaussian Noise Generator block. The initial

seed, the mean value and the variance in the simulation must be specified. We take the initial seed 41, the mean value 1 and the variance 1 in the simulation.

The Rayleigh probability density function is given by

$$f(x) = \begin{cases} \frac{x}{\sigma^2} e^{-\frac{x^2}{2\sigma^2}} & x \ge 0\\ 0 & x < 0 \end{cases}$$
 (5.7)

where σ^2 is known as the fading envelope of the Rayleigh distribution. The Simulink Communications toolbox provides the Rayleigh Noise Generator block. The initial seed and the sigma parameter in the simulation must be specified. We specify the initial seed 47 and the sigma parameter 1 in the simulation.

The Rician probability density function is given by

$$f(x) = \begin{cases} \frac{x}{\sigma^2} I_0(\frac{mx}{\sigma^2}) e^{\frac{x^2 + m^2}{2\sigma^2}} & x \ge 0\\ 0 & x < 0 \end{cases}$$
 (5.8)

where σ is the standard deviation of the Gaussian distribution that underlies the Rician distribution noise, $m^2 = m_I^2 + m_Q^2$, where m_I and m_Q are the mean values of two independent Gaussian components, and I_0 is the modified 0th-order Bessel function of the first kind given by

$$I_0(y) = \frac{1}{2\pi} \int_{-\pi}^{\pi} e^{y\cos t} dt$$
 (5.9)

Note that m and σ are not the mean value and standard deviation for the Rician noise. The Simulink Communications toolbox provides the Rician Noise Generator block. The initial seed, Rician K-factor and the sigma parameter must be specified in the simulation. We specify the initial seed 59, Rician K-factor 2 and the sigma parameter 1 in the simulation.

External terms kN_1 , kN_2 are added to systems of (5.4) and (5.5), where N_1 ,

 N_2 are the Gaussian noise and Rayleigh noise, respectively. In addition, the parameter b_1 of system (5.4) is replaced by x_3 , and the parameter b_1 of system (5.5) by $x_{3_{r_3}}$ where τ_3 =10. k=1, k=50 are used. For k=1, as shown in Fig. 5.6, robustness of lag synchronization of y_1 , y_2 and of lag quasi-synchronization of x_1 , x_2 is satisfactory. For k=50, as shown in Fig. 5.7, robustness of lag synchronization of y_1 , y_2 is still satisfactory, but robustness of the lag quasi-synchronization of x_1 , x_2 is lost.

Lastly, external terms kN_1 , kN_2 are added to the systems of (5.4) and (5.5), where N_1 , N_2 are the Gaussian noise and Rician noise, respectively. The parameter b_1 of system (5.4) is replaced by x_3 , and the parameter b_1 of system (5.5) by x_3 , where t_3 =10. k=1, k=30 are used. For k=1, as shown in Fig. 5.8, robustness of lag synchronization of t_1 , t_2 is satisfactory. For k=30, as shown in Fig. 5.9, robustness of lag synchronization of t_1 , t_2 is still satisfactory, but robustness of the lag quasi-synchronization of t_1 , t_2 is lost.

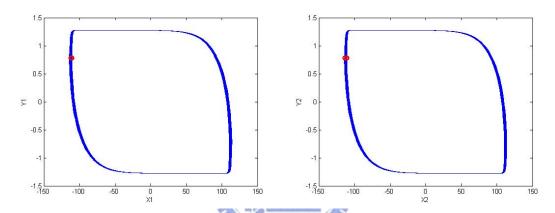
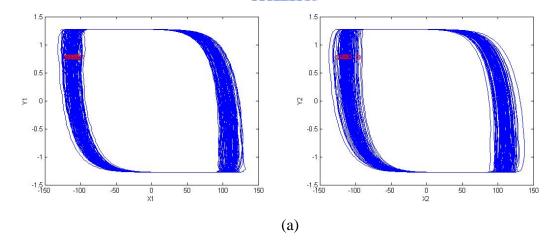


Fig. 5.1 The phase portraits for systems (5.1) and (5.2).



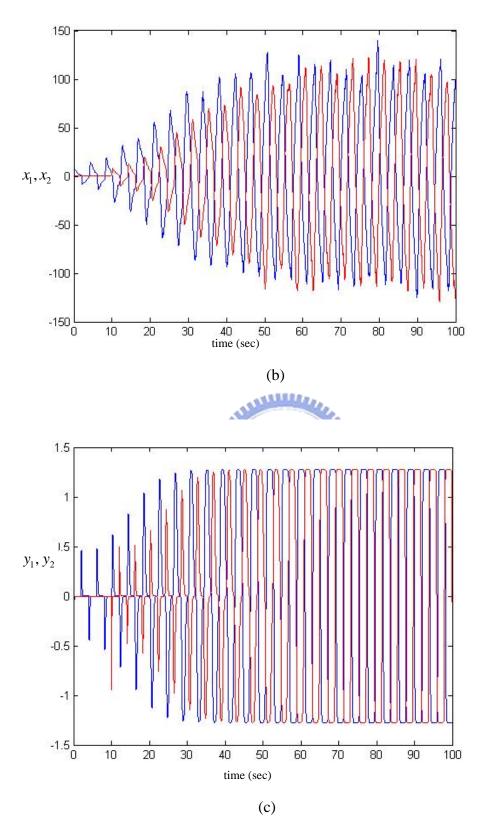
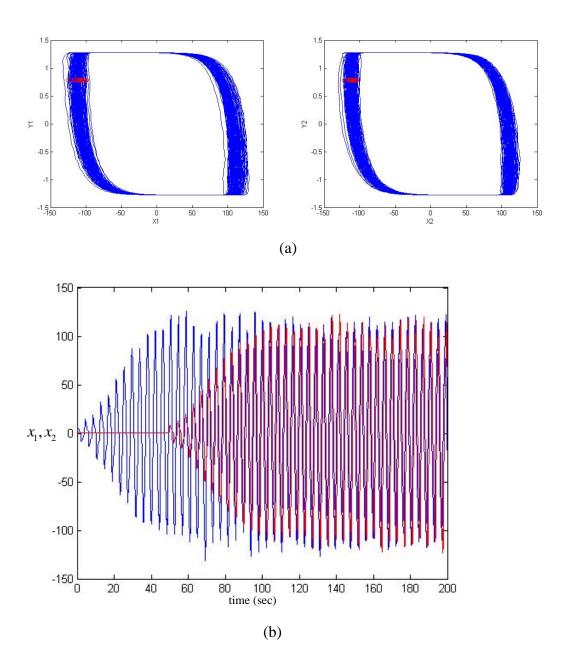
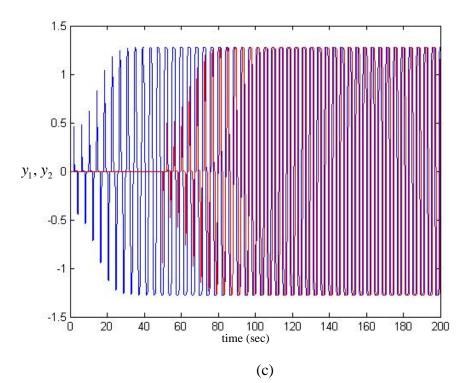
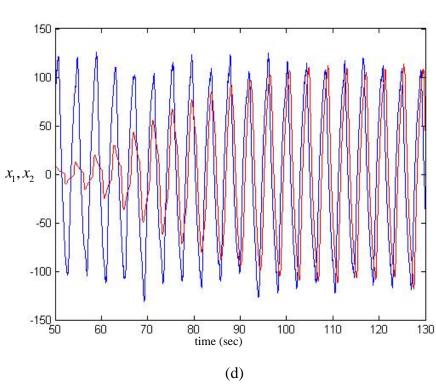


Fig. 5.2 (a) The phase portraits, $\tau_3=10$. (b) The time histories of x_1 (blue) and x_2 (red), $\tau_3=10$. (c) The time histories of y_1 (blue) and y_2 (red), $\tau_3=10$.







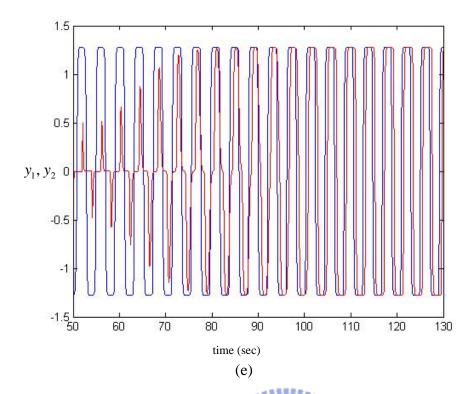
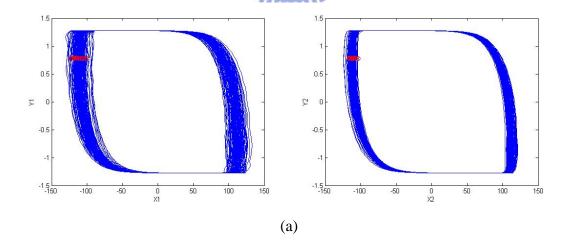
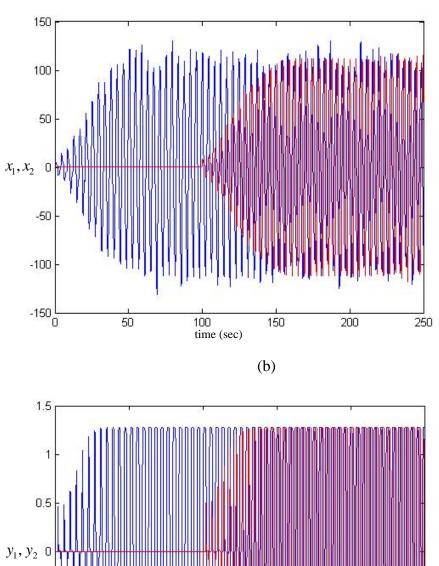
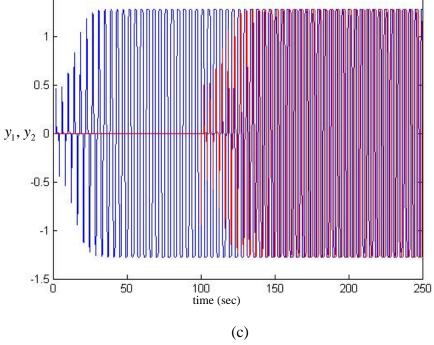


Fig. 5.3 (a) The phase portraits, τ_3 =50. (b) The time histories of x_1 (blue) and x_2 (red), τ_3 =50. (c) The time histories of y_1 (blue) and y_2 (red), τ_3 =50. (d) Magnified diagram of (b). (e) Magnified diagram of (c).







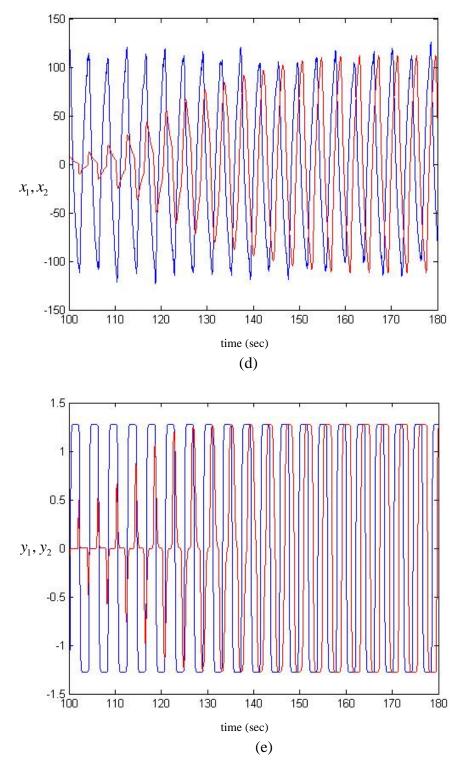


Fig. 5.4 (a) The phase portraits, τ_3 =100. (b) The time histories of x_1 (blue) and x_2 (red), τ_3 =100. (c) The time histories of y_1 (blue) and y_2 (red), τ_3 =100. (d) Magnified diagram of (b). (e) Magnified diagram of (c).

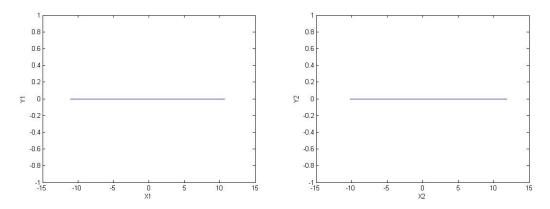
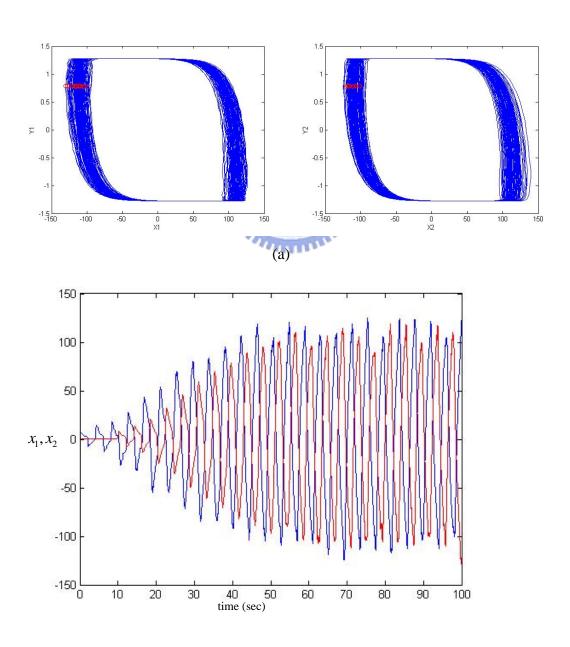


Fig. 5.5 The phase portraits, $\tau_3 = 100$.



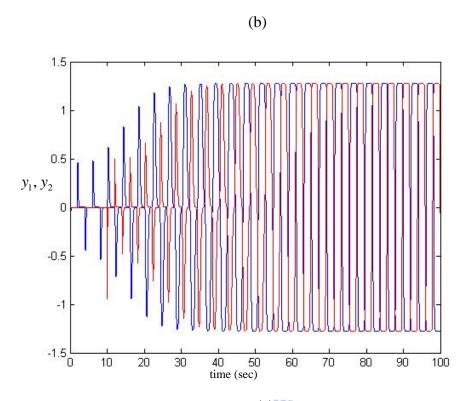
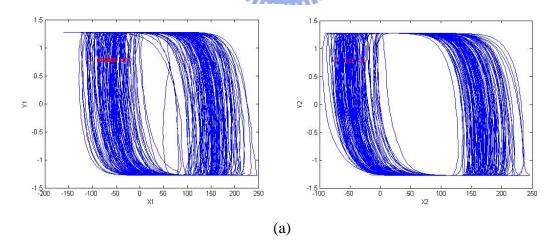


Fig. 5.6 (a) The phase portraits, $\tau_3 = 10$ and k = 1. (b) The time histories of x_1 (blue) and x_2 (red), $\tau_3 = 10$ and k = 1. (c) The time histories of y_1 (blue) and y_2 (red), $\tau_3 = 10$ and k = 1.



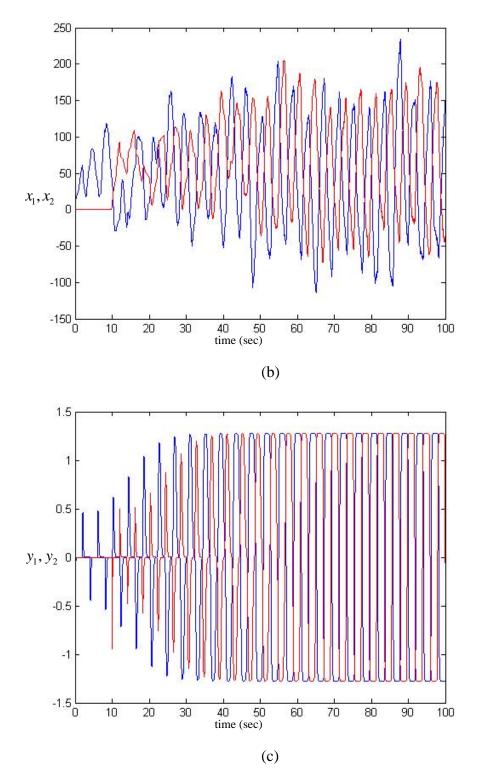
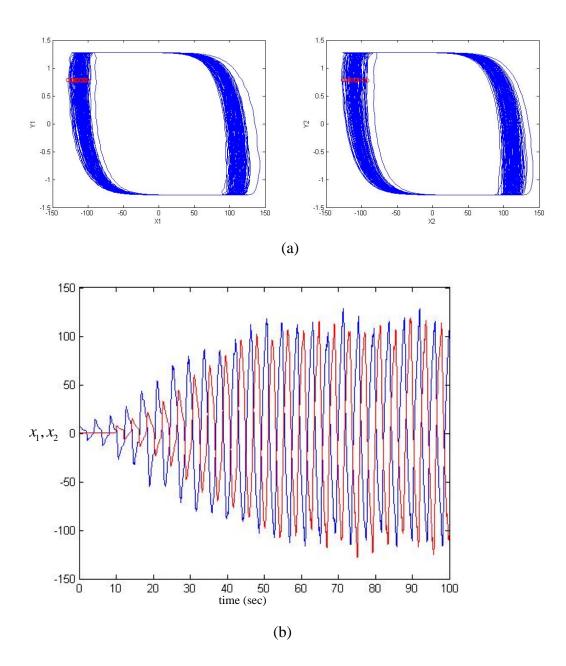


Fig. 5.7 (a) The phase portraits, τ_3 =10 and k=50. (b) The time histories of x_1 (blue) and x_2 (red), τ_3 =10 and k=50. (c) The time histories of y_1 (blue) and y_2 (red), τ_3 =10 and k=50.



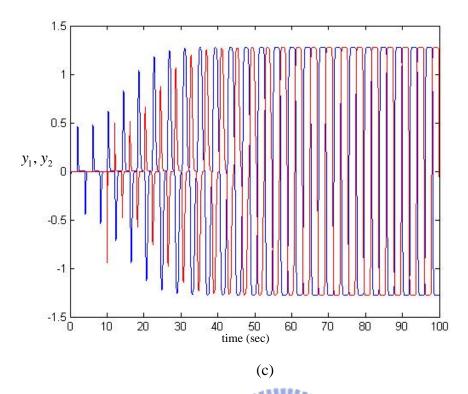
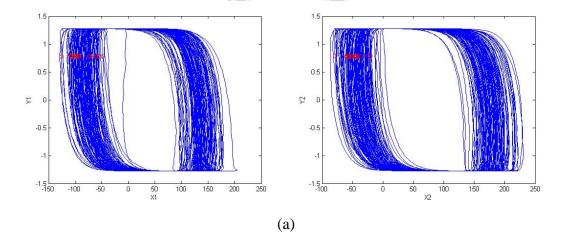


Fig. 5.8 (a) The phase portraits, $\tau_3 = 10$ and k = 1. (b) The time histories of x_1 (blue) and x_2 (red), $\tau_3 = 10$ and k = 1. (c) The time histories of y_1 (blue) and y_2 (red), $\tau_3 = 10$ and k = 1.



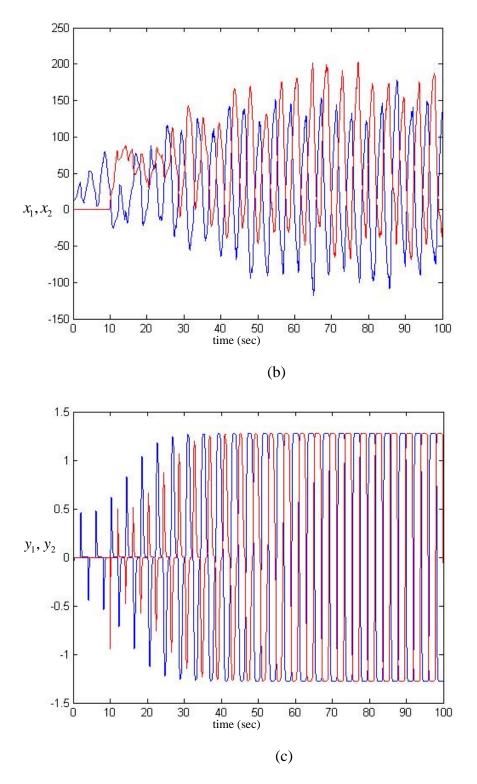


Fig. 5.9 (a) The phase portraits, τ_3 =10 and k=30. (b) The time histories of x_1 (blue) and x_2 (red), τ_3 =10 and k=30. (c) The time histories of y_1 (blue) and y_2 (red), τ_3 =10 and k=30.

Chapter 6

Conclusions

In this thesis, we have studied the chaos in the integral and fractional order double Ikeda systems by phase portraits, Poincaré maps and bifurcation diagrams. In Chapter 2, the chaos in integral and fractional order double Ikeda systems with total order of derivatives from 2 to 0.2 are studied by phase portraits, Poincaré maps and bifurcation diagrams. It is found that chaos exists in all cases.

In Chapter 3, lag or anticipated synchronization and the lag or anticipated anti-synchronization of two double Ikeda systems with different initial conditions are discovered. Cases 1~8 are the lag or anticipated synchronizations. Cases 9~16 are the lag or anticipated anti-synchronizations.

In Chapter 4, the chaotic behaviors of double Ikeda systems are obtained by replacing their delay time by a function of chaotic state variables of a second chaotic system. It is found that chaos exists for Case 1, 3, 5, 6. The chaotization of a double Ikeda system is studied by using a function of state variable of a second identical system to replace the delay time of the first system. It is found that in Case 7, 8, 9, chaotization exists.

In Chapter 5, robust lag chaos synchronization, lag quasi-synchronization and chaos control of two uncoupled double Ikeda system, are achieved by replacing the corresponding parameters of two systems by different chaotic state variables of a third chaotic system. Robustness of synchronization is studied by addition of various noises. The results are satisfactory.

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Appendix

Table 1. FRACTIONAL OPERATORS WITH APPROXIMATELY 2 db ERROR FROM $\mathcal{O}=10^{-2}$ TO 10^2 rad/sec

$$\frac{1}{s^{0.1}} \approx \frac{220.4s^4 + 5004s^3 + 503s^2 + 234.5s + 0.484}{s^5 + 359.8s^4 + 5742s^3 + 4247s^2 + 147.7s + 0.2099}$$

$$\frac{1}{s^{0.2}} \approx \frac{60.95s^4 + 816.9s^3 + 582.8s^2 + 23.24s + 0.04934}{s^5 + 134s^4 + 956.5s^3 + 383.5s^2 + 8.953s + 0.01821}$$

$$\frac{1}{s^{0.3}} \approx \frac{23.76s^4 + 224.9s^3 + 129.1s^2 + 4.733s + 0.01052}{s^5 + 64.51s^4 + 252.2s^3 + 63.61s^2 + 1.104s + 0.002267}$$

$$\frac{1}{s^{0.4}} \approx \frac{25s^4 + 558.5s^3 + 664.2s^2 + 44.15s + 0.1562}{s^5 + 125.6s^4 + 840.6s^3 + 317.2s^2 + 7.428s + 0.02343}$$

$$\frac{1}{s^{0.5}} \approx \frac{15.97s^4 + 593.2s^3 + 1080s^2 + 135.4s + 1}{s^5 + 134.3s^4 + 1072s^3 + 543.4s^2 + 20.1s + 0.1259}$$

$$\frac{1}{s^{0.6}} \approx \frac{8.579s^4 + 255.6s^3 + 405.3s^2 + 35.93s + 0.1696}{s^5 + 94.22s^4 + 472.9s^3 + 134.8s^2 + 2.639s + 0.009882}$$

$$\frac{1}{s^{0.7}} \approx \frac{4.406s^4 + 177.6s^3 + 209.6s^2 + 9.179s + 0.0145}{s^5 + 88.12s^4 + 279.2s^3 + 33.3s^2 + 1.927s + 0.0002276}$$

$$\frac{1}{s^{0.8}} \approx \frac{5.235s^3 + 1453s^2 + 5306s + 254.9}{s^4 + 658.1s^3 + 5700s^2 + 658.2s + 1}$$

$$\frac{1}{s^{0.9}} \approx \frac{1.766s^2 + 38.27s + 4.914}{s^3 + 36.15s^2 + 7.789s + 0.01}$$