

Kinetics and Mechanism of the NCN + NO₂ Reaction Studied by Experiment and TheoryTsong-Ju Yang,[†] Niann S. Wang,^{*,‡} L. C. Lee,[‡] Z. F. Xu,[§] and M. C. Lin^{*,†,§}

Department of Applied Chemistry and Center for Interdisciplinary Molecular Science, Chiao Tung University, Taiwan 30010, Taiwan, Department of Electrical and Computer Engineering, San Diego State University, San Diego, California 92182, and Department of Chemistry, Emory University, Atlanta, Georgia 30322

Received: June 16, 2008; Revised Manuscript Received: August 19, 2008

The rate constant for the NCN + NO₂ reaction has been measured by a laser photolysis/laser-induced fluorescence technique in the temperature range of 260–296 K at pressures between 100 and 500 Torr with He and N₂ as buffer gases. The NCN radical was produced from the photolysis of NCN₃ at 193 nm and monitored by laser-induced fluorescence with a dye laser at 329.01 nm. The rate constant was found to increase with pressure but decrease with temperature, indicating that the reaction occurs via a long-lived intermediate stabilized by collisions with buffer gas. The reaction mechanism and rate constant are also theoretically predicted for the temperature range of 200–2000 K and the He and N₂ pressure range of 10⁻⁴ Torr to 1000 atm based on dual-channel Rice–Ramsperger–Kassel–Marcus (RRKM) theory with the potential energy surface evaluated at the G2M//B3LYP/6-311+G(d) level. In the low-temperature range (<700 K), the most favorable reaction is the barrierless association channel that leads to the intermediate complex (NCN–NO₂). At high temperature, the direct O-abstraction reaction with a barrier of 9.8 kcal/mol becomes the dominant channel. The rate constant calculated by RRKM theory agrees reasonably well with experimental data.

1. Introduction

The cyanonitrene (NCN) radical has been observed by optical emission and presumed as an intermediate in active-nitrogen hydrocarbon flames long ago by Jennings and Linnett.¹ This radical was detected in a microwave discharge of a CF₄/N₂/He mixture by laser-induced fluorescence (LIF) by Smith et al.,² who further suggested that this radical could play a significant role in the combustion of nitramine propellants. The LIF technique is a sensitive method to detect this radical in combustion media, and it is applied to kinetic measurement. The importance of how this radical plays in hydrocarbon combustion is well recognized only after Moskaleva et al.^{3,4} revealed by a computational study that it can be produced by the spin-conserved reaction



This reaction suggests that NCN is the primary product leading to prompt NO formation in the hydrocarbon flame zone.^{3,4} This suggestion was shortly confirmed by Smith,⁵ who observed the presence of NCN in a methane/air flame. The production of NCN from the above reaction has recently been confirmed by Vasudevan et al.,⁶ who detected NCN by optical absorption in heating mixtures of diketene/N₂ and ethane/N₂ behind reflected shock waves. Their experimental data conclude that H + NCN is the dominant (and possibly the only) path for the CH + N₂ reaction. Recently, Harding et al.⁷ revisited the kinetics of CH + N₂ with multireference ab initio calculations. They found that the rates of this reaction at varied temperatures are higher than those calculated by Moskaleva and Lin.⁴ The higher reaction rate, resulting from the predicted lower heat of formation of NCN, agrees better with recent experimental data.^{6,7}

This result, in turn, proves that NCN is indeed an important intermediate in hydrocarbon/air combustion media.

Because NCN is a dominant prompt NO formation intermediate, knowledge of its reaction mechanism and kinetics is needed for understanding NO formation in the combustion process. For instance, El bakali et al.⁸ and Lamoureux et al.⁹ have demonstrated that the NO fractions observed in natural gas (a methane/ethane/propane mixture) flames agree better with the prediction if NCN is included in kinetic modeling. The kinetic data of NCN were almost not known prior to the new prompt NO formation mechanism put forth by Lin and co-workers^{3,4} in 2000, and the database continues to build up since then. Extensive theoretical calculations for reactions with various combustion species, such as C, H, CH_x (x = 1–3), CN, etc., have been reported by Moskaleva¹⁰ in her Ph.D. dissertation. Baren and Hershberger¹¹ measured the reaction rates of NCN with NO, O₂, C₂H₄, and NO₂ in a He buffer gas by the photolysis/LIF technique. The rate constant for the reaction with NO was measured more extensively over the temperature range of 298–573 K at varied He pressures, but for the other reactions, only the upper limits were estimated at 298 K at a low He pressure (<3 Torr). The reaction mechanisms of NCN with NO and NS were investigated by Chen and Ho.¹² Zhu and Lin have recently investigated the oxidation rates of NCN by O₂¹³ and O(³P)¹⁴ by ab initio calculations with molecular orbital and transition state (TS) theory.

Recently, Huang et al.¹⁵ have extensively measured the rate constants for the NCN + NO reaction in the presence of He (40–600 Torr) and N₂ (30–528 Torr) buffer gases in the temperature range of 254–353 K using the photolysis/LIF technique. The NCN radical was produced from the photodissociation of NCN₃ at 193 nm and monitored by LIF with a dye laser at 329.01 nm. They also calculated the rate constants by dual-channel Rice–Ramsperger–Kassel–Marcus (RRKM) theory based on the potential energy surface (PES) derived by the G2M

* To whom correspondence should be addressed.

[†] Chiao Tung University.

[‡] San Diego State University.

[§] Emory University.

method.¹⁵ The theoretical results compare reasonably well with experimental data.

The NCN radical has also been known as an interstellar species¹⁶ and a critical growth species in the synthesis of carbon nitride (C₃N₄)¹⁷ as well; thus, its reaction kinetics, in general, may also be relevant to such fields.

The study on NCN reaction kinetics is now extended to the reactant NO₂, and the results are reported in this article. Because NO₂ is the final product of nitrogen oxidation, it is expected to be present in most combustion media. The rate constant of the NCN + NO₂ reaction is certainly of interest in modeling combustion processes. However, up to now only an upper limit was given¹¹ for the NCN + NO₂ reaction at 298 K and low He pressure (<3 Torr).

In our current study, the rate constants for NCN + NO₂ were measured under varied pressures of He and N₂ (100–500 Torr) in the temperature range of 260–296 K using the experimental setup and the photolysis/LIF technique that we used for measuring the NCN + NO rate constants.¹⁵ We also calculated the reaction rate constants using the PES evaluated at the G2M//B3LYP/6-311+G(d) level of theory for the reaction of both NCN (³Σ_g⁻) and NO₂ (²A₁) in the ground state. The results of our calculations are compared with measured experimental data.

2. Experimental Section

2.1. Experimental Setup. The experimental setup for the measurement of reaction rate constants has been described previously.^{15,18} In brief, the reaction cell is a double-walled pyrex flask of volume of about 250 mL. Photolysis and LIF laser beams entered and exited the cell through two pairs of mutually perpendicular baffled side arms and intersected at the center of the cell. The LIF signal was projected through appropriate lenses and an interference filter into a photomultiplier tube (Hamamatsu R1136) mounted at a side window perpendicular to both laser beams. The LIF signal was amplified (Hamamatsu C1053, 5 MHz) and averaged with a gated integrator (Stanford Research System, SR250). The data were stored and analyzed by a computer.

The NCN radical was produced by photolysis of NCN₃ with an ArF excimer laser at 193 nm (Lambda Physik LEXTRA 50). NCN₃ was prepared¹⁹ by the reaction of NaN₃ + BrCN overnight and then degassed under vacuum before dilution in He or N₂. The absorption band for NCN₃ in the range 160–200 nm is very broad;²⁰ optical excitation at this band mainly dissociates into NCN + N₂.^{20–22} The concentration of NCN₃ was determined via absorption of a 184.9 nm Hg line with a 9.8 cm cell in the upstream of the reaction cell. The absorption cross section of 5.7 × 10⁻¹⁸ cm² at this wavelength²⁰ is used to determine the NCN₃ concentration.

The NCN radical was excited from the ground state, X³Σ_g⁻(0,0¹,0), to the excited state, A³Π_u(0,2⁰,0), at 329.01 nm by a dye laser (Continuum ND60) that was pumped by a doubled-output Nd:YAG laser (Continuum Surelite). The LIF signal was isolated by an interference filter centered at 330 nm (fwhm = 10 nm, 15% transmission). A delay/pulse generator (Stanford Research System, DG535) was used to control the firing of photolysis and LIF laser beams and to initiate the boxcar integrator as well. Typically, 30–100 laser shots at a repetition rate of 5–10 Hz were averaged to obtain a single LIF data point.

A radical precursor (NCN₃), reactant (NO₂), and buffer gas (He or N₂) were mixed in a flexible tube (30 cm long) before entering the reactor. The flow velocity (10–20 cm/s) of the gas mixture was adjusted so that the reaction zone was replenished

with fresh gas for each laser pulse, but it was in a semistatic condition during the time frame of kinetic measurement. The pressure of the system was measured with a Baratron gauge (MKS 122A). The flow rates of reactants and buffer gas were measured with a calibrated mass flowmeter (Tylan FM360). A temperature-controlled fluid flowed through the jacket of the reactor to maintain a constant reaction temperature (±2 K), which was measured with a K-type thermocouple placed 5 mm from the detection region.

The concentrations of the reactants were determined by

$$[A] = 9.66 \times 10^{18} P F_A / F_T T \quad (\text{molecules/cm}^3)$$

where P is the reaction pressure (in Torr), T is the reaction temperature (K), and F_A and F_T are the flow rates (in STP cm³/s, STP = 273.15 K, 760 Torr) of reactant A and the total reaction mixture, respectively.

The initial concentration of NCN radicals is estimated by

$$[\text{NCN}]_0 = \sigma \Phi F [\text{NCN}_3]$$

where σ is the absorption cross section (3.6 × 10⁻¹⁸ cm²)²⁰ of NCN₃ at 193 nm, Φ is the quantum yield of NCN from photolysis (assumed to be ~1), and F is the photon flux (photons/cm²) of the ArF laser.

He (99.999%) and N₂ (99.999%) were used without further purification. NO₂ was prepared by slowly reacting NO with excessive O₂ and stored in 3 atm of O₂ for 2 days before use. NO₂ was purified in freeze–pump–thaw cycles and diluted in He to a 2% mixture by a standard gas handling technique before it was used in experiments. Fourier transform infrared (FTIR) spectrometry was used to examine the impurities, mainly NO, in NO₂.

2.2. Rate Coefficient Measurement. All kinetic measurements were made under pseudo-first-order conditions, with [NO₂] more than 200 times greater than [NCN]. The initial concentrations for NCN were kept small, [NCN]₀ ≤ 1 × 10¹² molecules/cm³, to avoid possible radical–radical reactions and to ensure pseudo-first-order conditions. The rate equation is expressed as

$$-d \ln [\text{NCN}] / dt = k' = k_d + k'' [\text{NO}_2]$$

where k' (in s⁻¹) is the first-order decay coefficient, k_d is the radical loss rate coefficient mainly due to diffusion from the detection region, and k'' is the second-order reaction-rate coefficient of interest.

The rate coefficients of the NCN + NO₂ reaction were measured in the temperature range of 260–296 K at 100–500 Torr of He or N₂ pressure. The plots of the pseudo-first-order rate coefficients for k' vs [NO₂] at 296 K and He pressures of 100, 200, 300, 400, and 500 Torr are shown in Figure 1. The line slopes give the second-order bimolecular rate coefficients, k'' , which depend on the third-body pressure. The k'' values measured at temperatures of 260, 283, and 296 K with He and N₂ pressures of 100–500 Torr are summarized in Table 1. The k'' values at 283 K at varied He and N₂ pressures are further illustrated in Figure 2. As shown in the figure, the k'' values for N₂ are generally larger than those of He, indicating expectedly that N₂ is more efficient as a third body.

To examine the temperature effect, the k'' values at three temperatures at varied pressures of He are plotted in Figure 3. The log(k'') values depend very linearly on T^{-1} with varying negative slopes at different pressures; thus, they can be fitted by the Arrhenius equation, $k'' = A \exp(E_a/T)$.

The Arrhenius parameters derived from fitting the data at He and N₂ pressures of 100–500 Torr are summarized in Table 2,

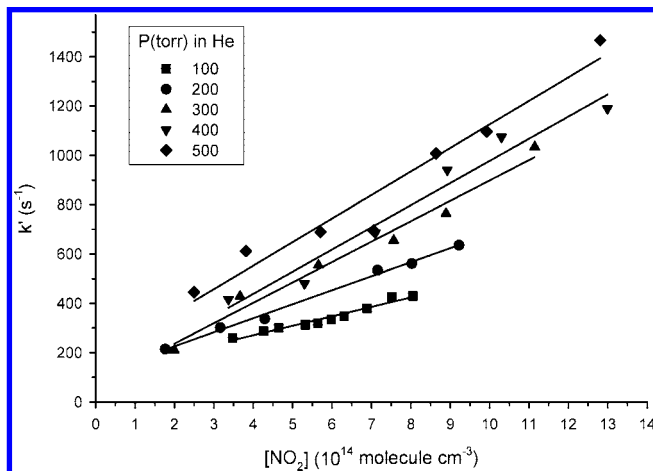


Figure 1. Plots of the pseudo-first-order decay rates of NCN, k' vs $[\text{NO}_2]$ at 296 K, in He pressures of (■) 100, (●) 200, (▲) 300, (▼) 400, and (◆) 500 Torr.

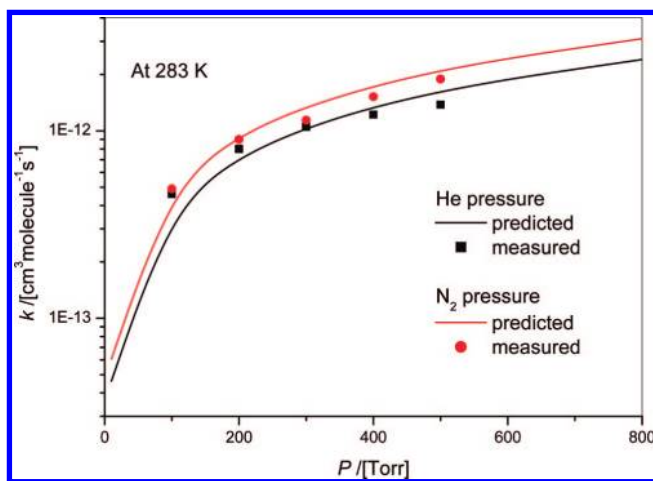


Figure 2. Plots of the reaction rate constants of NCN + NO₂, k'' vs pressure at 283 K for He (■) and N₂ (○).

TABLE 1: Experimental Conditions and Bimolecular Rate Coefficients, k'' , of the NCN + NO₂ Reaction

T (K)	P (Torr)		k'' [10^{-12} cm ³ /(molecule s)] ^a	
	He	N ₂	He	N ₂
296	100	100	0.38 ± 0.03	0.42 ± 0.06
	200	200	0.63 ± 0.02	0.73 ± 0.04
	300	300	0.78 ± 0.06	0.88 ± 0.11
	400	400	0.90 ± 0.09	1.21 ± 0.06
	500	500	0.95 ± 0.09	1.44 ± 0.12
283	100	100	0.46 ± 0.04	0.49 ± 0.03
	200	200	0.80 ± 0.02	0.90 ± 0.04
	300	300	1.05 ± 0.04	1.14 ± 0.03
	400	400	1.22 ± 0.03	1.52 ± 0.06
	500	500	1.38 ± 0.04	1.89 ± 0.10
260	100	100	0.77 ± 0.05	0.86 ± 0.07
	200	200	1.38 ± 0.07	1.55 ± 0.12
	300	300	1.88 ± 0.05	1.98 ± 0.15
	400	400	2.36 ± 0.07	2.89 ± 0.10
	500	500	2.63 ± 0.09	3.74 ± 0.14

^a Error limits are 1σ .

where A is in units of cm³/(molecule s) and both E_a and T are in K. The uncertainty shown in Table 2 represents 1 standard deviation in a linear least-squares fit. Both Figure 3 and Table 2 show that the k'' values have negative-temperature dependence, indicating that the NCN + NO₂ reaction in this low-temperature range is a typical associative one.

2.3. Discussion on Rate Measurements. The presence of NO molecules in the system, as either an impurity or a product

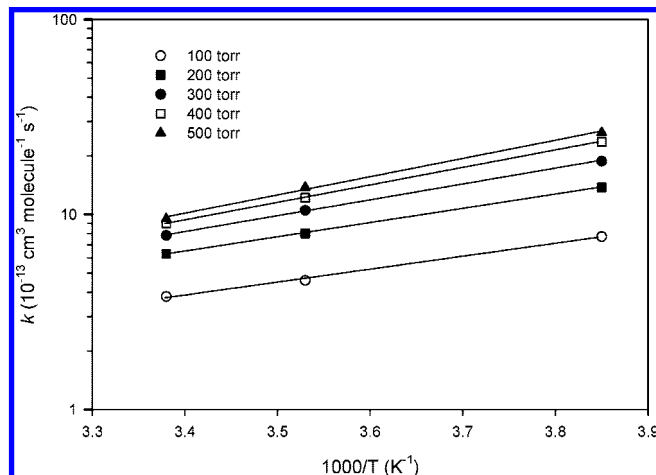


Figure 3. Arrhenius plots of k'' for He pressures of (○) 100, (■) 200, (●) 300, (□) 400, and (▲) 500 Torr.

TABLE 2: Arrhenius Parameters of the NCN + NO₂ Reaction at Various Pressures of He and N₂ [$k'' = A \exp(E_a/T)$]^a

P (Torr)	He		N ₂	
	A [10^{-15} cm ³ /(molecule s)]	E_a (K)	A [10^{-15} cm ³ /(molecule s)]	E_a (K)
100	2.12 ± 0.80	1530 ± 103	2.01 ± 1.56	1570 ± 199
200	2.10 ± 0.38	1686 ± 51	2.93 ± 1.04	1628 ± 96
300	1.37 ± 0.07	1878 ± 13	2.47 ± 0.22	1738 ± 24
400	0.83 ± 0.11	2068 ± 37	1.99 ± 1.08	1889 ± 145
500	0.65 ± 0.20	2166 ± 85	1.34 ± 0.53	2061 ± 106

^a Error limits are 1σ .

of 193-nm photolysis of NO₂ or both, might introduce errors in our measurements due to the rapid NCN + NO reaction.¹⁵ Efforts to minimize this effect were needed. As stated in the Experimental Section, NO₂ was stored under a high pressure of O₂ to ensure that all NO was converted to NO₂ before use. The impurities, especially NO, were examined by FTIR spectrometry, and the results showed no detectable NO in our NO₂ samples.

Because the ArF laser is powerful, its photodissociation of NO₂ may produce substantial photofragments, NO and O, to distort reaction rate measurements. The absorption cross section²³ of NO₂ at 193 nm is about 7×10^{-19} cm², which may totally dissociate into NO and O (i.e., $\Phi \sim 1$). This cross section is about a factor of 5 smaller than that of NCN₃.²⁰ However, [NO₂] is about 200 times [NCN₃], and the [NO] and [O] produced by photodissociation could initially be 40 times as high as [NCN], which may significantly distort the reaction rate measured. However, the laser energy (6.4 eV) is much higher than the dissociation energy of NO₂ (3.11 eV²⁴), and NO and O may carry initial kinetic energies of about 1.1 and 2.2 eV, which make the initial velocities about 8×10^3 and 1.6×10^4 m/s, respectively. The photofragments will thus fly fast away from the laser zone (less than 0.2 μs), and their concentrations are quickly diluted to far less than the initial value. Because the dilution time (approximately microseconds) is very short compared to the time frame of kinetic measurement (on the order of milliseconds), the distortion of measured rate constants by photofragments is negligible. This assertion is corroborated by a previous kinetic modeling,²⁵ in which [O] produced from photodissociation of NO₂ at 193 nm is less than 1% of [NO₂]. Under our experimental conditions, [NO]/[NO₂] was estimated to be less than 0.004. At 260 K and 400 Torr of N₂, the reaction rate constant¹⁵ for NCN + NO is 9.75×10^{-12} cm³/(mole-

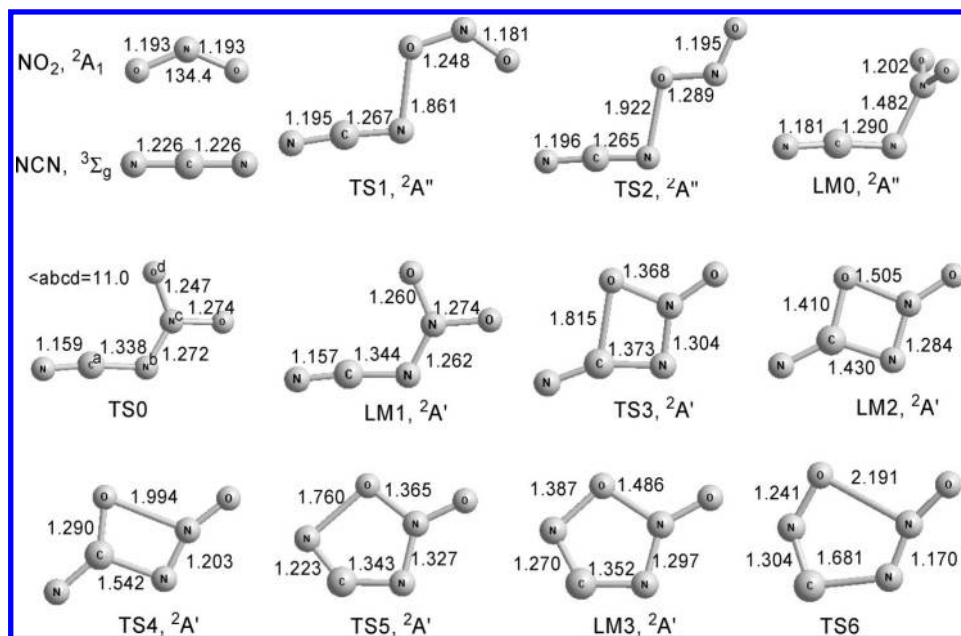


Figure 4. Geometric parameters of reactants, intermediates, and TSs optimized at the B3LYP/6-311+G(d) level of theory (in units of angstroms and degrees).

cule s), which is 3.4 times faster than that of NCN + NO₂. A 0.4% NO impurity in NO₂ would introduce an error of 1.4% in our measurements, which is within the quoted error limits. No further correction would be necessary. Another possible NO source in the system was an impurity in NO₂. Our FTIR measurements showed that they were less than those from photodissociation of NO₂ and hence would introduce a negligible error in the measured rate constants.

The bond energy of NCN–N₂ is quite low (0.3 eV).²⁰ However, this molecule is so complex that it may mainly dissociate into electronically and/or vibrationally excited species;^{20,21} thus, NCN may not move as fast as NO and O, so it can stay in the laser zone for kinetic measurement. The excited species should be quickly quenched to the ground state by a buffer gas. This assertion is corroborated by the fact that [NCN] in the ground state was observed to decay by pseudo-first-order kinetics in all of our measurements.

The NCN + NO₂ reaction has been investigated only once by Baren and Hershberger.¹¹ NCN was produced by photolysis of CH₂N₂ + C₂N₂ with a ArF excimer laser at 193 nm and detected by the LIF technique. They found that the reaction rate constant at 298 K and 3 Torr of He was very small; the upper limit of 1.5 × 10⁻¹⁴ cm³/(molecule s) was given.¹¹ We did not do the measurement at such a low pressure, but the data in Figure 2 show that the *k*' value at low pressure is indeed very low. Our data are compared with theoretical calculations as described below.

3. Computational Section

3.1. Computational Methods. The geometric parameters of the reactants, products, TSs, and molecular complexes of the NCN + NO₂ reaction system are optimized at the spin-unrestricted B3LYP/6-311+G(d) level of theory.^{26–29} All of the stationary points were identified for local minima and TSs by vibrational analysis. Intrinsic reaction-coordinate analyses^{30,31} were performed to confirm the connections of TSs with designated reactants, products, and intermediates. On the basis of the optimized geometries, higher level single-point energy calculations with the stationary points were performed by the

G2M(CC5) method,³² which calculated the base energy at the MP4/6-311G(d, p) level of theory and improved with the expanded basis set and the coupled cluster method as well as “higher level corrections (HLC)”. All electronic structure calculations were performed with the GAUSSIAN 03 program.³³

The rate constant for the barrierless association reaction that produces the NCNNO₂ complex was calculated by the VARI-FLEX program³⁴ based on the microcanonical RRKM theory.^{35–38} The component rates were evaluated at the *E*/*J*-resolved level, and the pressure dependence was treated by one-dimensional master equation calculations using the Boltzmann probability of the complex for the *J* distribution. For the barrierless association/decomposition process, the fitted Morse function was used to approximate the minimum potential energy path. The Lennard-Jones parameters ($\epsilon/k = 304$ K and $\sigma = 4.44$ Å) required for the collisional stabilization of the excited NCNNO₂ complex were computed from the NCNNO₂–He molecular interaction potential predicted at the B3LYP/6-311+G(d,p) level similar to the analogous NCN + NO reaction.¹⁵ The Lennard-Jones parameters for He and N₂ were taken from the literature.³⁹ The averaged step sizes for energy transfer per collision, $\langle\Delta E\rangle_{\text{down}}$, for He and N₂ were taken to be 150 and 250 cm⁻¹, respectively.

3.2. Reaction Potential Surface. The geometric parameters of the reactants, intermediates, and TSs optimized at the B3LYP/6-311+G(d) level of theory are shown in Figure 4. The relative potential energy diagram obtained at the G2M level based on the optimized structures is drawn in Figure 5. The intermediate complex LM0 is readily formed barrierlessly from the reactants NCN (³Σ_g⁻) and NO₂ (²A₁) as shown Figure 5, where the NCN in the ground state, ³Σ_g⁻, is designated as ³NCN. The barrierless association minimum energy path is presented by the Morse function, $V(R) = 22.7\{1 - \exp[-5.015(R - 1.595)]\}^2$, as shown in Figure 6. LM0 is in a planar symmetry of ²A'', and it is 19.4 kcal/mol below the reactants. It can isomerize into the complex LM1 (²A') of 3.1 kcal/mol energy below the reactants via a rotational TS0 with a potential barrier of 18.5 kcal/mol. NCO + NNO can be produced from LM1 via TS3, LM2, and TS4, which are four-membered ring structures. The energies of TS3,

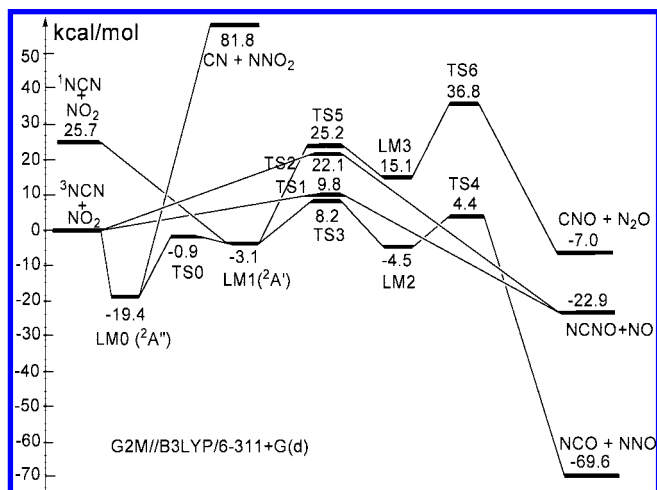


Figure 5. PES of ³NCN + NO₂ at the G2M//B3LYP/6-311+G(d) level of theory (in units of kilocalories per mole). Those of ¹NCN + NO₂ are also included.

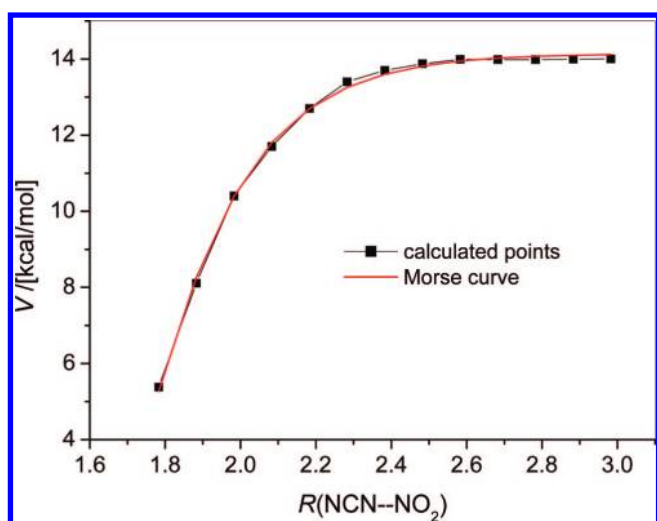


Figure 6. Morse function, $V(R)$, for the minimum energy path of the association process, ³NCN + NO₂ → LM0, at the B3LYP/6-311+G(d) level of theory.

LM2, and TS4 relative to the reactants are 8.2, -4.5, and 4.4 kcal/mol, respectively, and the formation of NCO + NNO is an exothermal process of 69.6 kcal/mol. Another reaction product, CNO + NNO, can be formed with higher barriers of five-membered ring structures from LM1 via TS5, LM3, and TS6, which have the relative energies of 25.2, 15.1, and 36.8 kcal/mol, respectively. The formation of CNO + NNO is an exothermal process of 7 kcal/mol. In addition to the association/isomerization decomposition channels, Figure 5 also shows two direct O-abstraction processes that produce the NCNO + NO products via TS1 and TS2, which are isomers with barriers of 9.8 and 22.1 kcal/mol, respectively. The energy level for the products, NCNO + NO, is lower than that of the reactants by 22.9 kcal/mol. The reaction channel for NCN (¹Δ_g) in the first excited state (designated as ¹NCN) is also shown in Figure 5.

3.3. Rate Constant Prediction. The rate constants are calculated based on the above PES using the VARIFLEX program³⁴ in the temperature range from 200 to 2000 K and in the pressure range from 10⁻⁴ Torr to 1000 atm. Overall, the most favorable reaction channels are

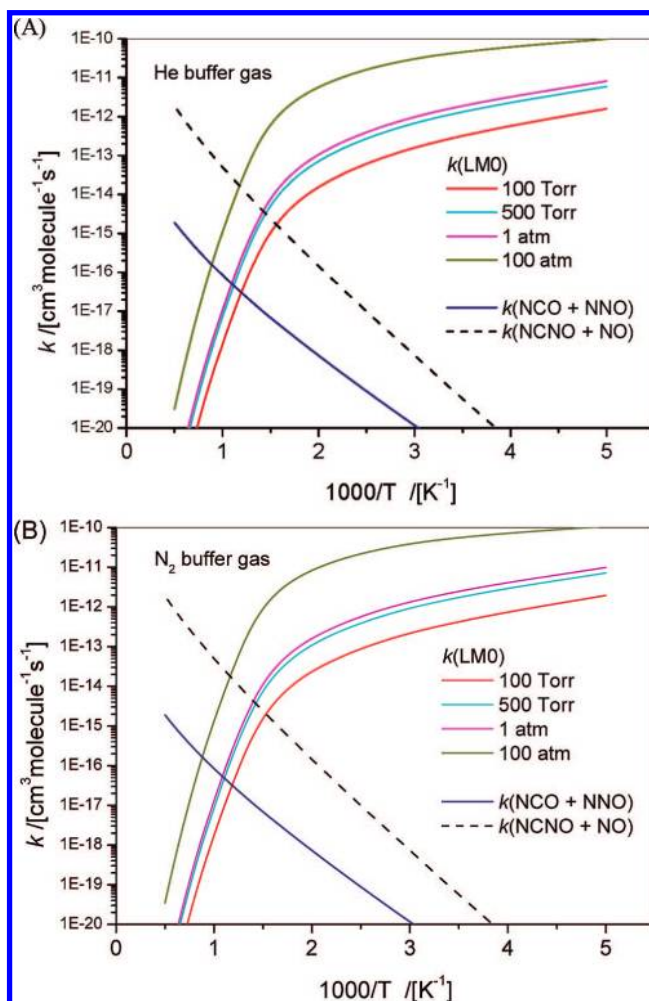


Figure 7. Calculated rate constants for the formation of LM0 and NCO + NNO via LM0 and NCNO + NO by direct abstraction: (A) He buffer gas and (B) N₂ buffer gas at varied pressures as indicated.



In channel (1), the decomposition of LM0 into the products NCO + NNO is controlled by TS3; thus, TS0 and LM1 can be ignored in the rate constant calculations. The $k(\text{LM0})$ and $k(\text{NCO} + \text{NNO})$ values calculated at the He and N₂ buffer-gas pressures of 100 Torr to 100 atm are plotted as a function of the temperature in Figure 7. The $k(\text{LM0})$ values increase with pressure but decrease with temperature. This positive-pressure and negative-temperature dependence effect is typical for an associative reaction. The $k(\text{NCO} + \text{NNO})$ values are independent of pressure but increase more than 5 orders of magnitude when the temperature increases from 300 to 2000 K. This is consistent with the calculated PES that the NCO + NNO products produced in channel (1) go through a potential barrier. As shown in Figure 7, the $k(\text{NCO} + \text{NNO})$ value is much less than those of $k(\text{LM0})$ at low temperature (<1000 K), but it increases systematically at high temperature.

The rate constant for NCNO + NO production by channel (2) is also shown in Figure 7 by the long-dashed curve. This direct O-abstraction reaction does not depend on buffer gas pressure, but it increases more than 12 orders of magnitude when the temperature increases from 200 to 2000 K. This positive-

temperature effect is caused by the fact that the NCNO + NO products produced in channel (22) go through the potential barriers of TS1 and/or TS2, as shown in Figure 5. As shown in Figure 7, the $k(\text{NCNO}+\text{NO})$ values are higher than the $k(\text{NCO}+\text{NNO})$ values by about 3 orders of magnitude over the whole temperature range of 200–2000 K. Thus, the contribution of $k(\text{NCO}+\text{NNO})$ to the total rate constant is negligible. Although $k(\text{NCNO}+\text{NO})$ is less than $k(\text{LM0})$ in the low-temperature range (<700 K), it becomes dominant at high temperature. The rate constants for the formation of LM0 (NCN–NO₂) at the high- and low-pressure limits with He and N₂ as third bodies can be given by the expressions

$$k^\infty(\text{LM0}) = 4.33 \times 10^{-11} T^{0.297} \exp(-47/T) \text{ cm}^3/(\text{molecule s}) \quad (200\text{--}2000 \text{ K})$$

$$k^0(\text{He}) = 7.58 \times 10^{20} T^{-18.22} \exp(-4354/T) \text{ cm}^6/(\text{molecule}^2 \text{ s}) \quad (200\text{--}2000 \text{ K})$$

$$k^0(\text{N}_2) = 2.09 \times 10^{22} T^{-18.63} \exp(-4532/T) \text{ cm}^6/(\text{molecule}^2 \text{ s}) \quad (200\text{--}2000 \text{ K})$$

The rate constant for the direct abstraction channel (2) can be represented by

$$k(\text{NCNO}+\text{NO}) = 1.22 \times 10^{-11} \exp(-5463/T) \text{ cm}^3/(\text{molecule s}) \quad (200\text{--}2000 \text{ K})$$

The total rate constants (k_t) calculated at indicated buffer pressures, including the high-pressure limit, are shown in Figure 8. As shown in the figure for the P and T effects, the k_t curve for every pressure shows a minimum, resulting effectively from a combination of redissociation and stabilization of the excited NCN–NO₂ intermediate and the direct abstraction reaction giving NCNO + NO. The minimum shifts to higher temperature when the pressure increases. The k_t values at low temperature increase with pressure because they are mainly contributed from $k(\text{LM0})$. On the other hand, the k_t values at high temperature do not depend on the pressure because they are mainly due to contribution from the abstraction rate constant $k(\text{NCNO}+\text{NO})$. As shown in Figure 8, at the high-pressure limit, the rate constant increases moderately with temperature in the whole range because the partition function of the loose association TS increases more rapidly with temperature than those of the reactants.

This product relation is further illustrated by the branching ratios calculated for producing LM0 and NCNO + NO at He buffer gas pressures of 10 Torr, 1 atm, and 10 atm, as shown in Figure 9. The branching ratio is almost equal to 1 for $k(\text{LM0})$ at low temperature and equal to 1 for $k(\text{NCNO}+\text{NO})$ at high temperature. The crossing points at the branching ratio of 0.5 are about 560, 700, and 780 K, when the He pressures are at 10 Torr, 1 atm, and 10 atm, respectively. This shift of the branching ratio to higher temperature can be explained again by the competition between stabilization and NCNO abstraction product formation.

4. Discussion

The measured rate constants and calculated values are compared in Figures 10 and 11 as functions of pressure and temperature, respectively. Because our measurements were performed in the low-temperature range of 260–296 K, the major product is LM0 (NCNNO₂) as alluded to above. Because of the high barriers for the formation of decomposition products NCO + N₂O via NCNNO₂ and the direct abstraction products,

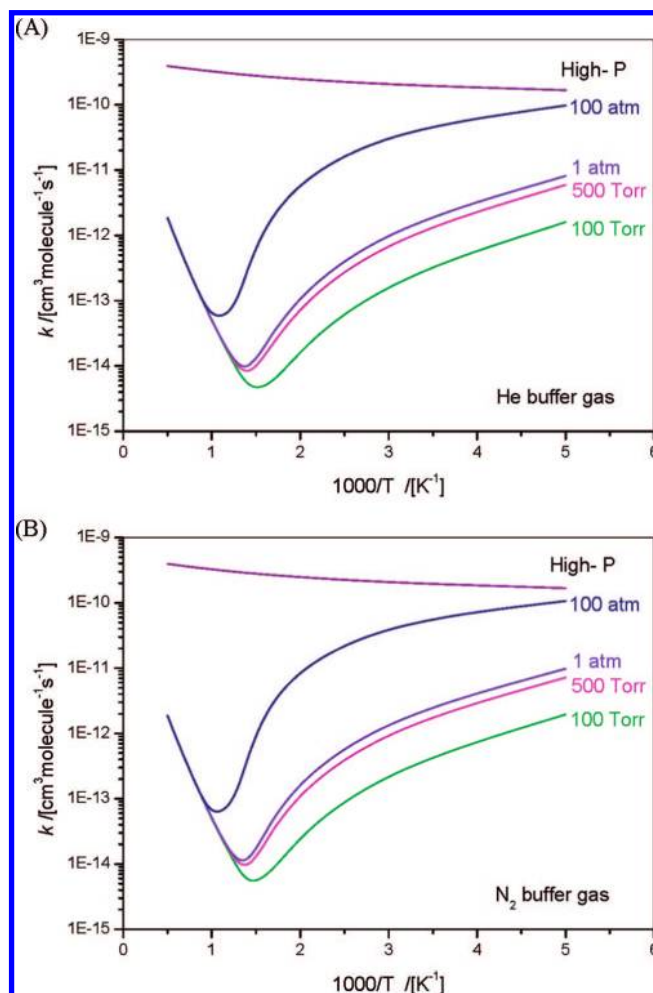


Figure 8. Calculated total rate constants of the NCN + NO₂ reaction at varied pressures as indicated: (A) He buffer gas; (B) N₂ buffer gas.

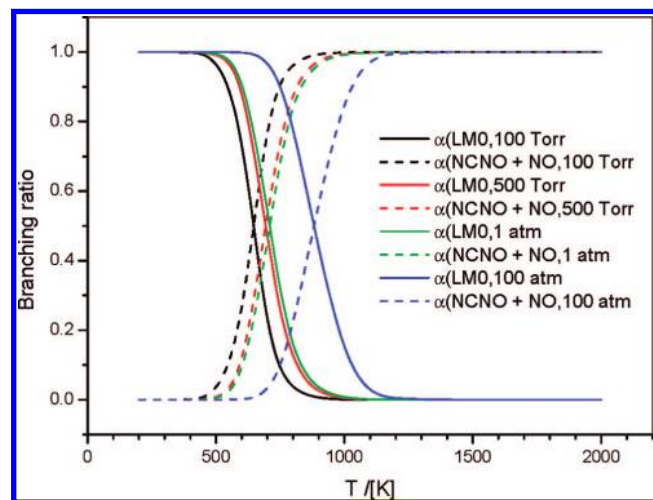


Figure 9. Calculated branching ratios for the production of LM0 and NCNO + NO at varied He pressures as indicated.

NCNO + NO, their appearance can occur only at $T > 700$ K; at such high temperature, the radical precursor NCN₃ is too unstable for kinetic studies. Both Figures 10 and 11 indicate that the observed P and T effects on k_t can be quite reasonably accounted for by the ab initio molecular orbital statistical RRKM theory calculations, similar to the NCN + NO case.¹⁵

It is of interest to compare the current rate constants with those of other reactants available in the literature. At 2000 K

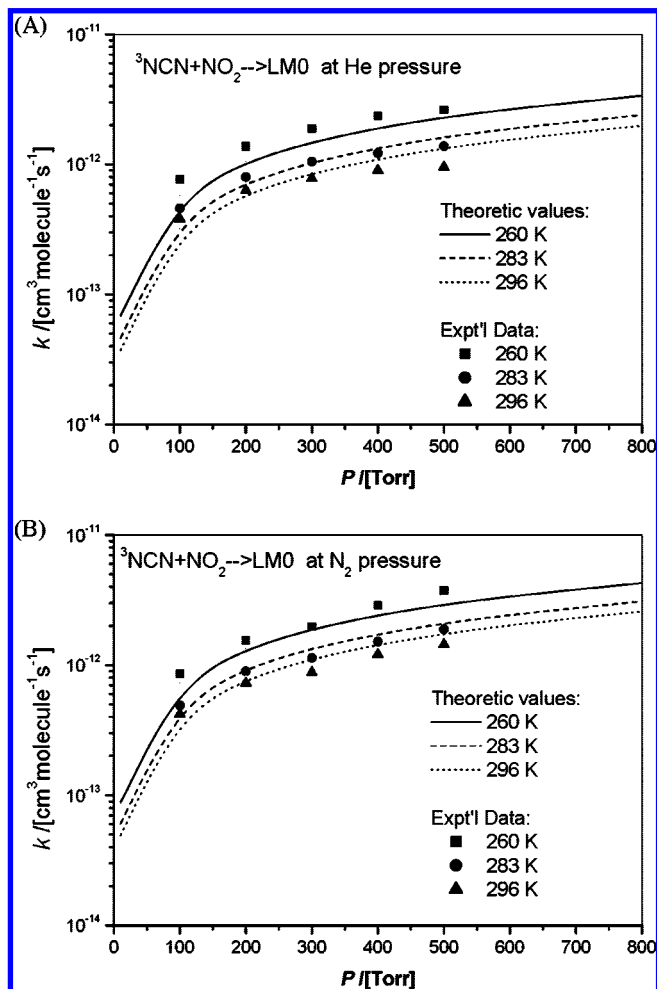


Figure 10. Comparison of the calculated rate constants with experimental data at 260, 283, and 296 K for (A) He buffer gas and (B) N₂ buffer gas.

and ~ 1 atm pressure, the total rate constant, k_t , for the NCN + NO₂ reaction is predicted to be 1.9×10^{-12} cm³/(molecule s). This is much larger than the k_t value of 7.2×10^{-16} cm³/(molecule s) predicted for the NCN + O₂ reaction¹³ under the same temperature and pressure conditions, but it is much smaller than that of the NCN + O reaction, 1.3×10^{-10} cm³/(molecule s).¹⁴ These results show that NO₂ is much more reactive than O₂ but much less than O, as one would expect.

At 296 K and 300 Torr of He pressure, the association rate constant for NCN + NO₂ is 7.8×10^{-13} cm³/(molecule s), as shown in Table 1. Interestingly, this value is smaller than that of the NCN + NO reaction, 6.7×10^{-12} cm³/(molecule s), under the same temperature and pressure conditions.¹¹ Similarly, at 296 K and 400 Torr of N₂ pressure, the association rate constant for NCN + NO₂ is 1.2×10^{-12} cm³/(molecule s), which is smaller than that of 7.4×10^{-12} cm³/(molecule s) for the NCN + NO reaction under the same temperature and pressure conditions.¹⁵ Because both reactions take place by the barrierless association mechanism, the apparent higher reactivity of NCN toward NO than NO₂ actually reflects the higher stability of the NCNNO intermediate (31 kcal/mol for *trans*-NCNNO and 22 kcal/mol for *cis*-NCNNO)¹⁵ than NCNNO₂ (19 kcal/mol as shown in Figure 5). The higher stability of NCNNO is expected to have a smaller redissociation rate, which, in turn, gives rise to a higher stabilization and overall rate constant.

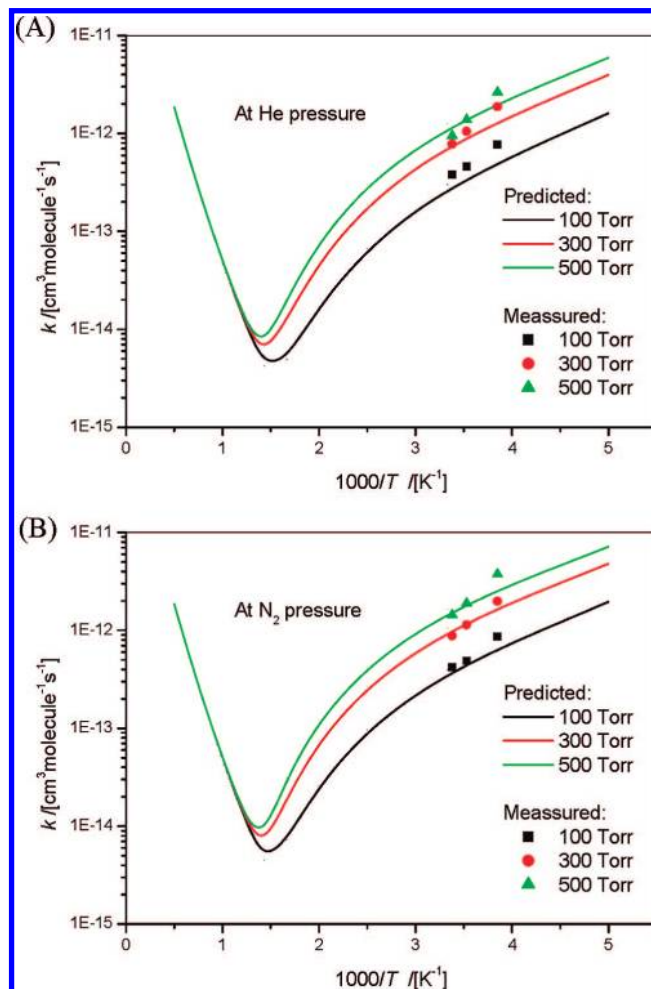


Figure 11. Comparison of calculated rate constants with experimental data at 100, 300, and 500 Torr for (A) He buffer gas and (B) N₂ buffer gas.

5. Conclusion

The kinetics and mechanism for the NCN + NO₂ reaction have been investigated experimentally and computationally. In our measured temperature range of 260–296 K, the reaction was found to be strongly and positively dependent on the pressure and negatively dependent on the temperature. The observed pressure and temperature effects can be reasonably accounted for by the results of a dual-channel RRKM calculation, with the predicted mechanism initially taking place via a moderately strongly bonded NCNNO₂ intermediate in the ²A'' state. The reaction rate constant for the NCNO + NO products by direct abstraction is calculated to be about 3 orders of magnitude higher than that for NCO + N₂O formation via NCNNO₂ in the temperature range of 200–2000 K. Although the constant for the production of NCNO + NO is smaller than of the association reaction at low temperatures (<700 K), this O-abstraction channel is predicted to be dominant at high temperatures, for example, under nitramine propellant combustion conditions.

Acknowledgment. L.C.L. and M.C.L. acknowledge the National Science Council of Taiwan for support of their visiting professorships at Taiwan National Chiao Tung University. Z.F.X. is grateful for partial support from the U.S. Office of Naval Research under Grant N00014-02-1-0133. T.-J.Y. and N.S.W. thank the National Science Council of Taiwan (Contract NSC 95-2113-M-009-028) for support.

References and Notes

- (1) Jennings, K. R.; Linnett, J. W. *Trans. Faraday Soc.* **1960**, *56*, 1737.
- (2) Smith, G. P.; Copeland, R. A.; Crosley, D. R. *J. Chem. Phys.* **1989**, *91*, 1987.
- (3) Moskaleva, L. V.; Xia, W. S.; Lin, M. C. *Chem. Phys. Lett.* **2000**, *331*, 269.
- (4) Moskaleva, L. V.; Lin, M. C. *Proc. Combust. Inst.* **2000**, *28*, 2393.
- (5) Smith, G. P. *Chem. Phys. Lett.* **2003**, *367*, 541.
- (6) Vasudevan, V.; Hanson, R. K.; Bowman, C. T.; Golden, D. M.; Davidson, D. F. *J. Phys. Chem. A* **2007**, *111*, 11818.
- (7) Harding, L. B.; Klippenstein, S. J.; Miller, J. A. *J. Phys. Chem. A* **2008**, *112*, 522.
- (8) El bakali, A.; Pillier, L.; Desgroux, P.; Lefort, B.; Gasnot, L.; Pauwels, J. F.; da Costa, I. *Fuel* **2006**, *85*, 896.
- (9) Lamoureux, N.; El bakali, A.; Gasnot, L.; Pauwels, J. F.; Desgroux, P. *Combust. Flame* **2008**, *153*, 186.
- (10) Moskaleva, L. V. Ph.D. Dissertation, Emory University, Atlanta, GA, 2001.
- (11) Baren, R. E.; Hershberger, J. F. *J. Phys. Chem. A* **2002**, *106*, 11093.
- (12) Chen, H. T.; Ho, J. J. *J. Phys. Chem. A* **2005**, *109*, 2564.
- (13) Zhu, R. S.; Lin, M. C. *Int. J. Chem. Kinet.* **2005**, *37*, 593.
- (14) Zhu, R. S.; Lin, M. C. *J. Phys. Chem. A* **2007**, *111*, 6766.
- (15) Huang, C. L.; Tseng, S. Y.; Wang, T. Y.; Wang, N. S.; Xu, Z. F.; Lin, M. C. *J. Chem. Phys.* **2005**, *122*, 184321.
- (16) Odell, C. R.; Miller, C. O.; Cochran, A. L.; Cochran, W. D.; Opal, C. B.; Barker, R. S. *Astrophys. J.* **1991**, *368*, 616.
- (17) Bernard, D. J.; Linnen, C.; Harker, A.; Michels, H. H.; Addison, J. B.; Ondercin, R. *J. Phys. Chem. B* **1998**, *102*, 6010.
- (18) You, Y. Y.; Wang, N. S. *J. Chin. Chem. Soc.* **1993**, *40*, 337.
- (19) Milligan, D. E.; Jacox, M. E.; Bass, A. M. *J. Chem. Phys.* **1965**, *43*, 3149.
- (20) Okabe, H.; Mele, A. *J. Chem. Phys.* **1969**, *51*, 2100.
- (21) Kroto, H. W. *J. Chem. Phys.* **1966**, *44*, 831.
- (22) Pontrelli, G. J.; Anastassiou, A. G. *J. Chem. Phys.* **1965**, *42*, 3735.
- (23) Nakayama, T.; Kitamura, M. Y.; Watanabe, K. *J. Chem. Phys.* **1959**, *30*, 1180.
- (24) Herzberg, G. *Electronic Spectra and Electronic Structure of Polyatomic Molecules*, Van Nostrand: New York, 1966; Vol. III.
- (25) Park, J.; Lin, M. C. *J. Phys. Chem. A* **1997**, *101*, 2643.
- (26) Becke, A. D. *J. Chem. Phys.* **1993**, *98*, 5648.
- (27) Becke, A. D. *J. Chem. Phys.* **1992**, *96*, 2155.
- (28) Becke, A. D. *J. Chem. Phys.* **1992**, *97*, 9173.
- (29) Lee, C.; Yang, W.; Parr, R. G. *Phys. Rev.* **1988**, *37B*, 785.
- (30) Gonzalez, C.; Schlegel, H. B. *J. Chem. Phys.* **1989**, *90*, 2154.
- (31) Gonzalez, C.; Schlegel, H. B. *J. Chem. Phys.* **1990**, *94*, 5523.
- (32) Mebel, A. M.; Morokuma, K.; Lin, M. C. *J. Chem. Phys.* **1995**, *103*, 7414.
- (33) Frisch, M. J.; Trucks, G. W.; Schlegel, H. B.; Scuseri, G. E.; Robb, M. A.; Cheeseman, J. R.; Zakrzewski, V. G.; Montgomery, J. A., Jr.; Stratmann, R. E.; Burant, J. C.; Dapprich, S.; Millam, J. M.; Daniels, A. D.; Kudin, K. N.; Strain, M. C.; Farkas, O.; Tomasi, J.; Barone, V.; Cossi, M.; Cammi, R.; Mennucci, B.; Pomelli, C.; Adamo, C.; Clifford, S.; Ochterski, J.; Petersson, G. A.; Ayala, P. Y.; Cui, Q.; Morokuma, K.; Malick, D. K.; Rabuck, A. D.; Raghavachari, K.; Foresman, J. B.; Cioslowski, J.; Ortiz, J. V.; Baboul, A. G.; Stefanov, B. B.; Liu, G.; Liashenko, A.; Piskorz, P.; Komaromi, I.; Gomperts, R.; Martin, R. L.; Fox, D. J.; Keith, T.; AlLaham, M. A.; Peng, C. Y.; Nanayakkara, A.; Gonzalez, C.; Challacombe, M.; Gill, P. M. W.; Johnson, B. G.; Chen, W.; Wong, M. W.; Andres, J. L.; HeadGordon, M.; Replogle, E. S.; Pople, J. A. *GAUSSIAN 03*, revision A.9; Gaussian, Inc.: Pittsburgh, PA, 1998.
- (34) Klippenstein, S. J.; Wagner, A. F.; Dunbar, R. C.; Wardlaw, D. M.; Robertson, S. H. *VARIFLEX*, version 1.00; 1999.
- (35) Wardlaw, D. M.; Marcus, R. A. *Chem. Phys. Lett.* **1984**, *110*, 230.
- (36) Wardlaw, D. M.; Marcus, R. A. *J. Chem. Phys.* **1985**, *83*, 3462.
- (37) Klippenstein, S. J. *J. Chem. Phys.* **1992**, *96*, 367.
- (38) Klippenstein, S. J.; Marcus, R. A. *J. Chem. Phys.* **1987**, *87*, 3410.
- (39) Mourits, F. M.; Rummens, F. H. A. *Can. J. Chem.* **1977**, *55*, 3007.

JP805311U